$\mathcal{L} = \frac{1}{4g^2} G_{\mu\nu} G_{\mu\nu} + \sum_{j} \overline{g}_j (i \partial^{\mu} D_{\mu} + m_j) g_j$ $\alpha_{\rm c}(Q^2)$ t decays (N110) Heavy Quarkonia (NLO) 0.3 o e'e jets & shapevines NMLO- e.w. precision fits (N2LO) 7 pp -> jets (NLO) where $G_{\mu\nu}^{\alpha} \equiv \partial_{\mu} A_{\nu}^{\alpha} - \partial_{\nu} A_{\mu}^{\alpha} + i f_{\delta e}^{\alpha} A_{\mu}^{\delta} A_{\nu}^{c}$ ▼ 00 -> 0 (NMLO) 0.2 and Du = du + it An 0.1 $QCD \alpha_s(M_z) = 0.1181 \pm 0.0011$ Q[GeV] 1000 10 100

PDG $\alpha_s(m_z)$ discussion

J. Huston, K. Rabbertz, (G. Zanderighi) Les Houches 2023

PDG averaging for $\alpha_s(m_z)$

- Every two years, the QCD section in the Particle Data Book is updated; part of that update is a review of the world average of $\alpha_s(m_z)$, revising it to include the impact of new measurements and calculations
- The last revision was in 2021 (minor update in 2022)
- The selection of results to include in the α_s averaging are restricted by the following considerations:
 - published in a peer-reviewed paper at the time of the report (or is based on a summary of results that have been published in a peer-reviewed journal, such as the FLAG report)
 - based on the most complete perturbative predictions of at least NNLO accuracy, accompanied by reliable estimates of all experimental and theoretical uncertainties



Figure 9.2: Summary of determinations of $\alpha_s(M_z^2)$ from the seven sub-fields discussed in the text. The yellow (light shaded) bands and dotted lines indicate the pre-average values of each sub-field. The dashed line and blue (dark shaded) band represent the final world average value of $\alpha_s(M_z^2)$. The "*" symbol within the "hadron colliders" sub-field indicates a determination including a simultaneous fit of PDFs.

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Non-perturbative effects

- Some inclusive quantities such as the e⁺e⁻ cross sections to hadrons have small non-perturbative corrections (~Λ⁴/Q⁴), while others such as event-shape distributions, can have corrections that go as Λ/Q
- Analyses of the τ hadronic decay width and spectral functions are performed with N³LO predictions, but low Q (m_τ) results in non-negligible non-perturbative corrections, whose treatment differs among the different calculations
- Collider measurements access the highest values of Q where nonperturbative effects are expected to be less important
- Both collider and DIS/DY data go into global PDF fits, which themselves are dependent on non-perturbative forms



Figure 9.2: Summary of determinations of $\alpha_s(M_Z^2)$ from the seven sub-fields text. The yellow (light shaded) bands and dotted lines indicate the pre-average sub-field. The dashed line and blue (dark shaded) band represent the final world $\alpha_s(M_Z^2)$. The "*" symbol within the "hadron colliders" sub-field indicates a determ a simultaneous fit of PDFs.

Significant advance

Fits of α_s using power corrections in the three-jet region

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ABSTRACT: In this work we study the impact of recent findings regarding non-perturbative corrections in the three-jet region to e^+e^- hadronic observables, by performing a simultaneous fit of the strong coupling constant α_s and the non-perturbative parameter α_0 . We extend the calculation of these power corrections, already known for thrust and Cparameter, to other e^+e^- hadronic observables. We find that for some observables the non-perturbative corrections are reasonably well behaved in the two-jet limit, while for others they have a more problematic behaviour. If one limits the fit to the three-jet region and to the well-behaved observables, one finds in general very good results, with the extracted value of α_s agreeing well with the world average. This is the case in particular for the thrust and *C*-parameter for which notably small values of α_s have been reported when non-perturbative corrections have been computed using analytic methods. Furthermore, the more problematic variables are also well described provided one stays far enough from the two-jet limit, while in this same region they cannot be described using the traditional implementation of power-corrections based on two-jet kinematics.

KEYWORDS: Perturbative QCD, QCD Phenomenology, electron-positron scattering





 0.1175 ± 0.0010

- Results within sub-fields 1-6 were pre-averaged (using unweighted average)
- FLAG19 result itself is an average and is taken $\alpha_s(M_Z^2) = 0.1182 \pm 0.0008$.



Figure 9.5: Lattice determinations that enter the FLAG2019 average. The yellow (light shaded) band and dotted line indicates the average value for this sub-field. The dashed line and blue (dark shaded) band represent the final world average value of $\alpha_s(M_s^2)$.^a

FLAG21 result (too late to be used in previous combination)

$$\alpha_{\overline{\mathrm{MS}}}^{(5)}(M_Z) = 0.1184(8)$$



Figure 9.2: Summary of determinations of $\alpha_s(M_Z^2)$ from the seven sub-fields discussed in the text. The yellow (light shaded) bands and dotted lines indicate the pre-average values of each sub-field. The dashed line and blue (dark shaded) band represent the final world average value of $\alpha_s(M_Z^2)$. The "*" symbol within the "hadron colliders" sub-field indicates a determination including a simultaneous fit of PDFs.

 A non-lattice result was determined from sub-fields 1-6 using a χ²-averaging method

 $\alpha_s(M_Z^2) = 0.1176 \pm 0.0010$,

(without lattice)

 FLAG result itself is an average and is taken as is

 $\alpha_s(M_Z^2) = 0.1182 \pm 0.0008, \quad (\text{lattice})$

- Combine two numbers in unweighted average, and take uncertainty as an average of the two uncertainties (conservative) $\alpha_s(M_Z^2) = 0.1179 \pm 0.0009$
- A weighted average of all 7 categories would give

 $\alpha_s(M_Z^2) = 0.1180 \pm 0.0006.$



Figure 9.2: Summary of determinations of $\alpha_s(M_Z^2)$ from the seven sub-fields discussed in the text. The yellow (light shaded) bands and dotted lines indicate the pre-average values of each sub-field. The dashed line and blue (dark shaded) band represent the final world average value of $\alpha_s(M_Z^2)$. The "*" symbol within the "hadron colliders" sub-field indicates a determination including a simultaneous fit of PDFs.

Collider measurements of α_s

As the number of NNLO calculations has increased, there have been a growing number of determinations of $\alpha_s(m_z)$ at that order (or higher) from the LHC experiments that have nominal uncertainties that rival the full world average uncertainty

ش Z р_т

- * event shapes
- It would be nice to understand those uncertainties better, especially if PDF uncertainties are taken into account



N³LL+N³LO



New LHC results



Exp.	√s / TeV	Lumi / fb ⁻¹	Theory	Obs.	α _s (M _Z)	Δα _s exp	Δα _s oth	Δα _s scl	Ref.
CMS	13	33.5	NNLO	Jet pT	0.1166	14 (NP)	7	4	JHEP12 (2022) 035
ATLAS	13	139	NNLO	TEEC	0.1175	6	12	+32 -11	2301.09 351
ATLAS	13	139	NNLO	ATEEC	0.1185	9	11	+22 -2	2301.09 351
CMS	13	36.3	NNLO	2D m _{jj}	0.1201	12 (NP)	9	8	SMP- 21-008
CMS	13	36.3	NNLO	3D m _{jj}	0.1201	10 (NP)	10	5	SMP- 21-008
ATLAS	8	20.2	N4LLa+ N3LO	Z pT	0.1183	4	6	4	CONF- 2023- 015



Figure 2: Transverse-momentum distribution of Z bosons predicted with DYTurbo [31] at different values of $\alpha_s(m_Z)$, using the MSHT20 PDF set [32].

The statistical analysis for the determination of $\alpha_s(m_Z)$ is performed with the xFitter framework [60]. The value of $\alpha_s(m_Z)$ is determined by minimising a χ^2 function which includes both the experimental uncertainties and the theoretical uncertainties arising from PDF variations:

$$\chi^{2}(\beta_{\exp},\beta_{th}) = \sum_{i=1}^{N_{data}} \frac{\left(\sigma_{i}^{\exp} + \sum_{j} \Gamma_{ij}^{\exp} \beta_{j,\exp} - \sigma_{i}^{th} - \sum_{k} \Gamma_{ik}^{th} \beta_{k,th}\right)^{2}}{\Delta_{i}^{2}} + \sum_{j} \beta_{j,\exp}^{2} + \sum_{k} \beta_{k,th}^{2}.$$
(1)



FIG. 7 (color online). The $\Delta \sigma / \Delta P_{\rm T}$ cross section versus $P_{\rm T}$. Cross-section values are plotted at the bin center. The horizontal bars represent the bin extent and the vertical bars are the crosssection uncertainties. The solid (black) crosses are the data and all uncertainties except the integrated luminosity uncertainty are combined and plotted. The solid (red) histogram is the RESBOS calculation. The dash-dotted (blue) bars of the $P_{\rm T} > 25$ GeV/*c* region are the FEWZ2 calculation. For the calculations, only numerical uncertainties are included but they are too small to be visible. The inset is the $P_{\rm T} < 25$ GeV/*c* region with a linear ordinate scale.



Figure 3: Determination of $\alpha_s(m_Z)$ at various different orders in the QCD perturbative expansion, using the MSHT20 PDF set. The filled area represents missing higher order uncertainties estimated through scale variations, the vertical error bars include experimental and PDF uncertainties.

Table 1: Summary of the uncertainties for the determination of $\alpha_s(m_Z)$.

Experimental uncertainty	+0.00044	-0.00044
PDF uncertainty	+0.00051	-0.00051
Scale variations uncertainties	+0.00042	-0.00042
Matching to fixed order	0	-0.00008
Non-perturbative model	+0.00012	-0.00020
Flavour model	+0.00021	-0.00029
QED ISR	+0.00014	-0.00014
N4LL approximation	+0.00004	-0.00004
Total	+0.00084	-0.00088

Table 2: Summary of N³LL fits with NNLO PDFs.

PDF set	$\alpha_{\rm s}(m_Z)$	PDF uncertainty	$g [{\rm GeV}^2]$	$q [\text{GeV}^4]$	χ^2 /dof
MSHT20 [32]	0.11839	0.00040	0.44	-0.07	96.0/69
NNPDF40 [78]	0.11779	0.00024	0.50	-0.08	116.0/69
CT18A [79]	0.11982	0.00050	0.36	-0.03	97.7 /69
HERAPDF20 [63]	0.11890	0.00027	0.40	-0.04	132.3/69

Running $\alpha_{\rm s}$



Looking forward

- We have been considering lattice QCD determinations of α_s independently of experimental/phenomenological determinations
- In the future, it may be useful to group lattice QCD determinations with experimental determinations of as that have systematics of similar origin
 - · del Debbio, Ramos; arXiv:2101.04672



Global PDF fits

- There is a wide variety of data in modern global PDF analyses, over 3500 data points for CT18
- The data includes DIS, DY (including precision W/Z), jet production, top production
- All predictions at NNLO, all depending on α_s

Experimental data in CT18 PDF analysis



α_s and gluon (Lagrange multiplier studies)

- Also, all of the experiments in the global fit do not speak with a unified voice, further weakening the discrimination power
- We end up with a fairly parabolic χ² dependence of α_s(m_z), but it's clear that different experiments have different preferences
- At 68% CL, α_s(m_z)=0.1166+/-0.0018 (for CT18 at NNLO)

