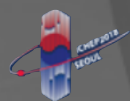


Lattice QCD for High-Energy Physics

Phiala Shanahan, MIT

Image Credit: 2018 EIC User's Group Meeting



ICHEP 2018 SE^oUL

XXXIX INTERNATIONAL CONFERENCE ON *high Energy* PHYSICS

JULY 4 - 11, 2018 COEX, SEOUL



Massachusetts
Institute of
Technology

Lattice QCD for HEP

We are entering the

- Precision era of lattice QCD for simple systems
- Beginning of reliable lattice QCD results for nuclear matrix elements

I will highlight some new results in these areas since ICHEP2016

Lattice QCD can provide input for

Decay constants, form factors, mixing parameters



Hadronic vacuum polarisation and light-by-light scattering



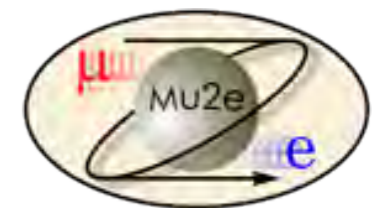
Neutrino-nucleus interactions



Dark matter-nucleon and DM-nucleus interactions



Muon-nucleus cross-sections



Parton distribution functions



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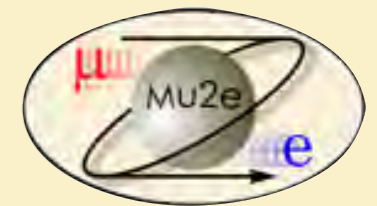
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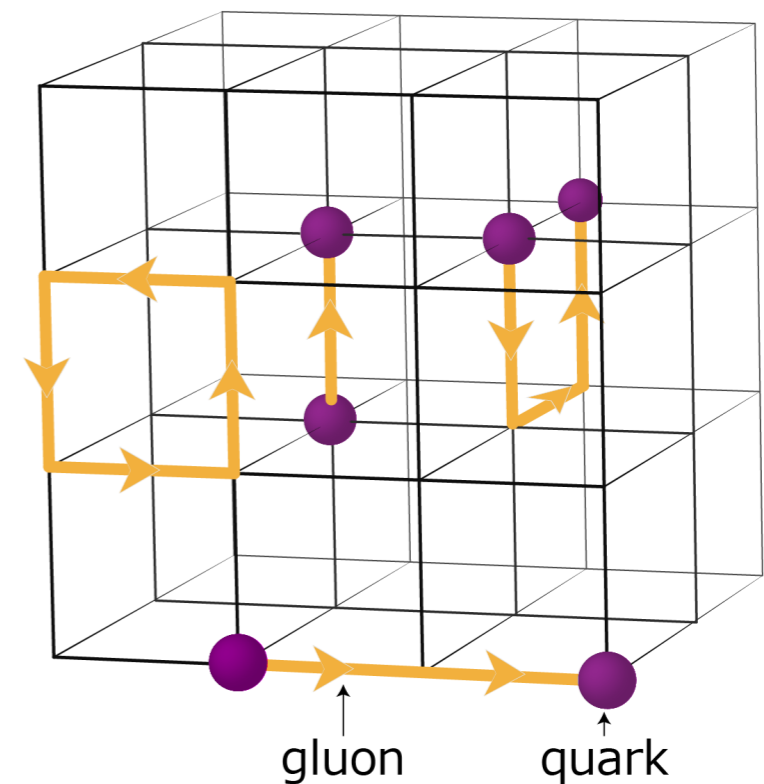
Parton distribution functions



Lattice QCD

Numerical first-principles approach to non-perturbative QCD

- Discretise QCD onto 4D space-time lattice
- Approximate QCD path integral using Monte-Carlo methods and importance sampling
- Run on supercomputers and dedicated clusters
- Take limit of vanishing discretisation, infinite volume, physical quark masses



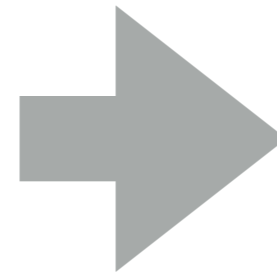
Lattice QCD

Numerical first-principles approach to non-perturbative QCD

INPUT

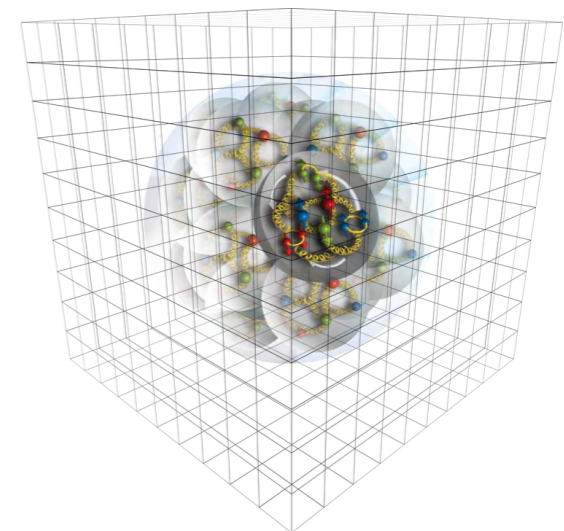
Lattice QCD action has same free parameters as QCD: quark masses, α_S

- Fix quark masses by matching to measured hadron masses, e.g., π, K, D_s, B_s for u, d, s, c, b
- One experimental input to fix lattice spacing in GeV (and also α_S), e.g., $2S-1S$ splitting in Y , or f_π or Ω mass



OUTPUT

Calculations of all other quantities are QCD predictions



Highlights since ICHEP2016

- Quark masses and α_s
- Flavour physics
- $g - 2$
- Parton distribution functions
- Neutrino physics
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Quark masses and α_s

The quark masses and α_s are the fundamental parameters of QCD
Their precise values are important for precision tests of the Standard Model

e.g., Next-generation of high-luminosity colliders will measure **Higgs partial widths to sub-percent precision** to look for deviations from Standard-Model expectations

➔ Need **Standard Model calculations at same sub-percent precision**; largest uncertainties are currently in m_c , m_b , & α_s [LHCHXSWG-DRAFT-INT-2016-008]

➔ Lepage, Mackenzie, Peskin, [arXiv:1404.0319]

- Precision goals for m_c , m_b , & α_s needed by high-luminosity ILC
- Outlined timeline for lattice QCD progress

Continued progress towards precision goals since ICHEP2016

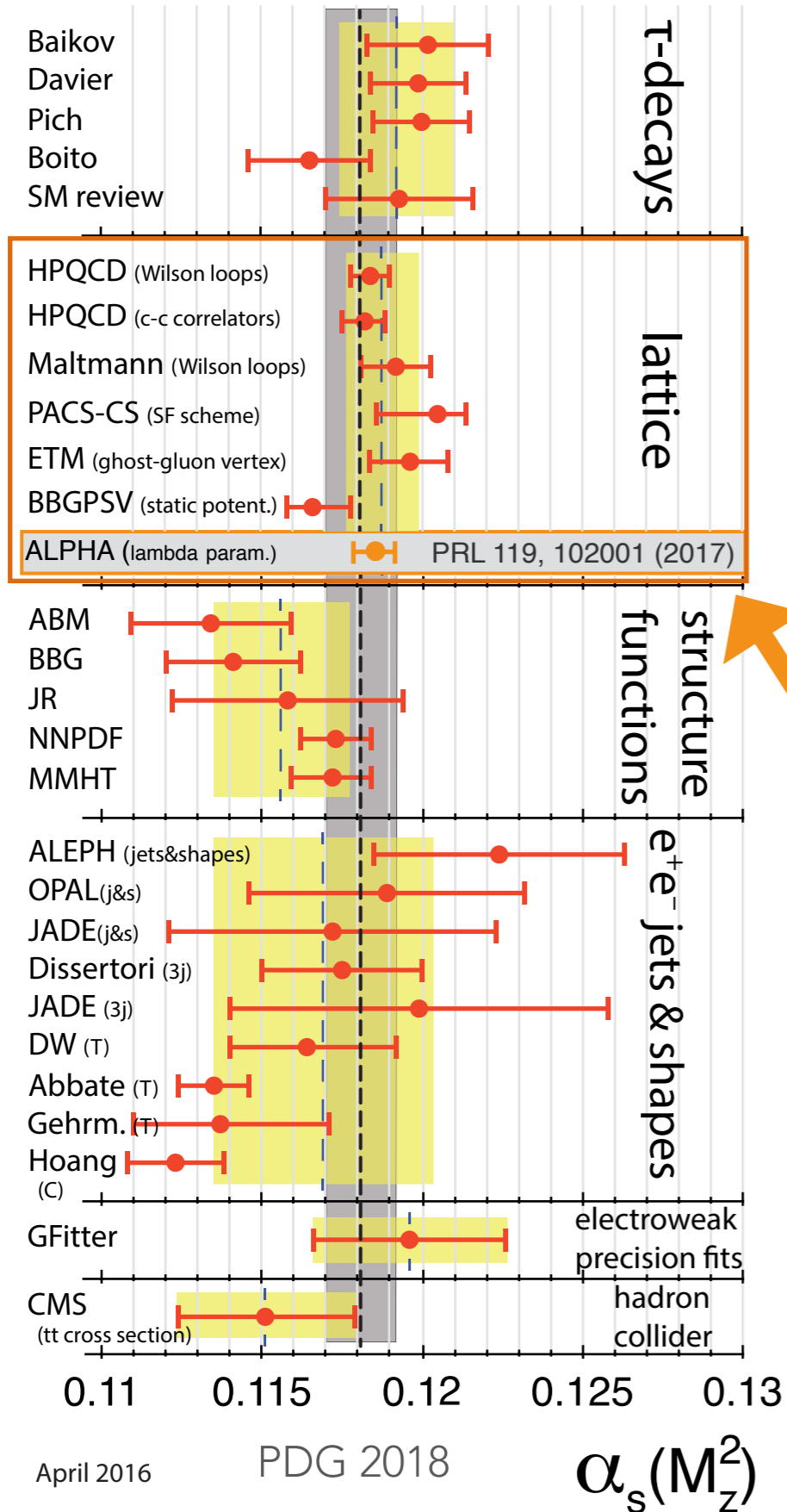
Next goals:

- Correlated determinations of m_c , m_b , and α_s
- Dynamical QED

α_s update

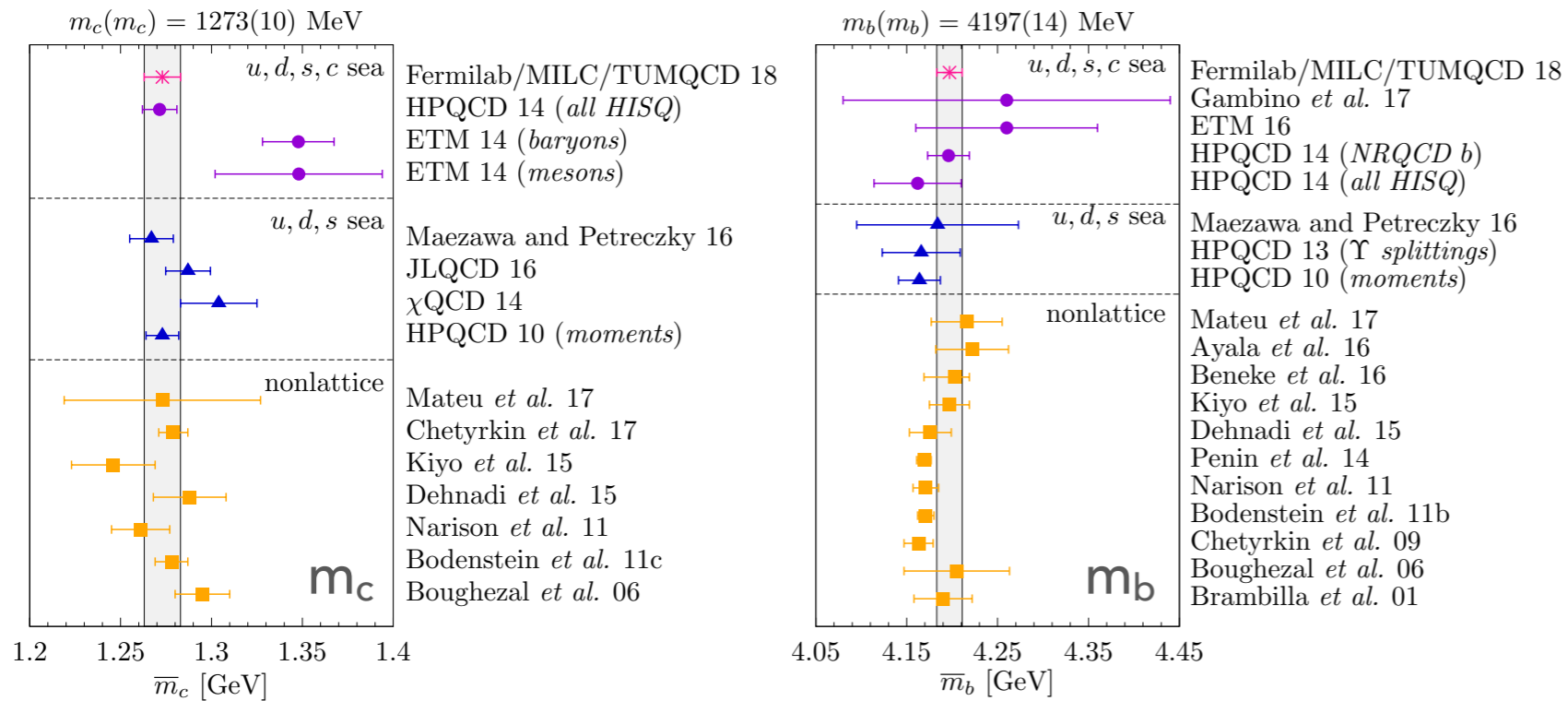
Best lattice QCD uncertainties $\sim 0.6-0.7\%$,
 approaching ILC target: 0.6% .
 Twice as precise as non-lattice world average

- Several independent lattice QCD methods available to obtain α_s
- Results consistent, despite significantly different sources of systematic uncertainty

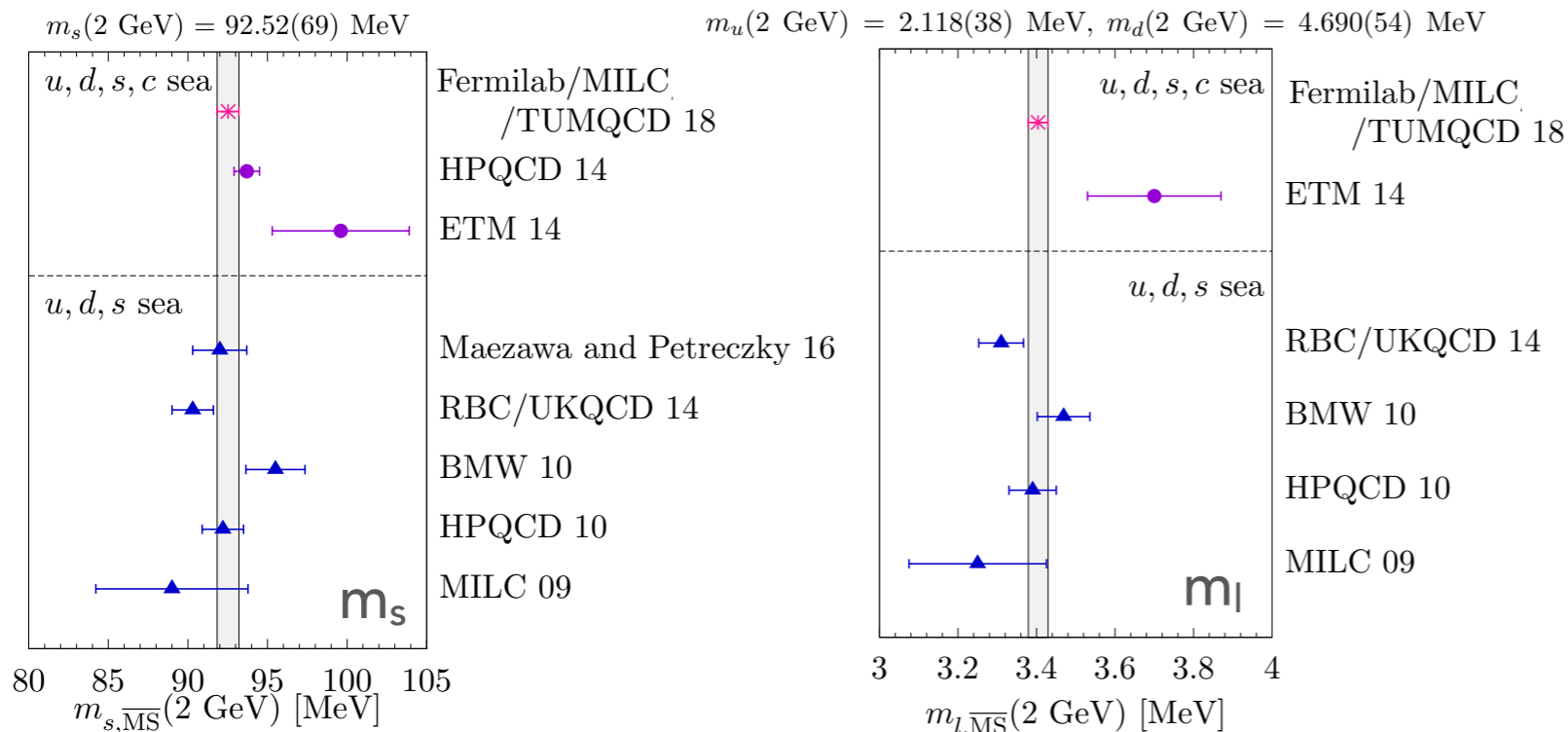


2017 Highlight:
 New lattice QCD determination based on finite size scaling (rather than Wilson Loops and quarkonia) consistent and precise: $\alpha_s = 0.1185(8)(3)$ [PRL119, 102001]

Quark masses update



MRS mass definition [arXiv:1712.04983]



2018 Highlight:

Significant improvement in heavy quark mass determinations using new method based on heavy-quark effective theory [arXiv:1802.04248]

Precision

ILC goals

$$\delta m_b \sim 0.3$$

$$0.3$$

$$\delta m_c \sim 0.8$$

$$0.7$$

Will improve further with inclusion of finer lattice spacings

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Flavour physics from lattice QCD


Lattice QCD can provide precise hadronic matrix elements for flavour physics:

$$[\text{experiment}] = [\text{known}] \times [\text{CKM}] \times [\text{hadronic matrix element}]$$

This leads to effort in two main directions:




- 1** Lattice QCD + experiment
 CKM quark mixing matrix elements
 Simple tree-level processes in lattice QCD allow determination of all elements and phases other than V_{tb}

$$\left(\begin{array}{ccc} V_{ud} & V_{us} & V_{ub} \\ \pi \rightarrow \mu\nu & K \rightarrow \pi l\nu \\ & K \rightarrow \mu\nu & B \rightarrow \pi l\nu \\ V_{cd} & V_{cs} & V_{cb} \\ D \rightarrow \pi l\nu & D \rightarrow K l\nu & B \rightarrow D, D^* l\nu \\ D \rightarrow l\nu & D_s \rightarrow l\nu & \\ V_{td} & V_{ts} & V_{tb} \\ B^0 - \bar{B}^0 \text{ mixing} & B_s - \bar{B}_s & \end{array} \right)$$

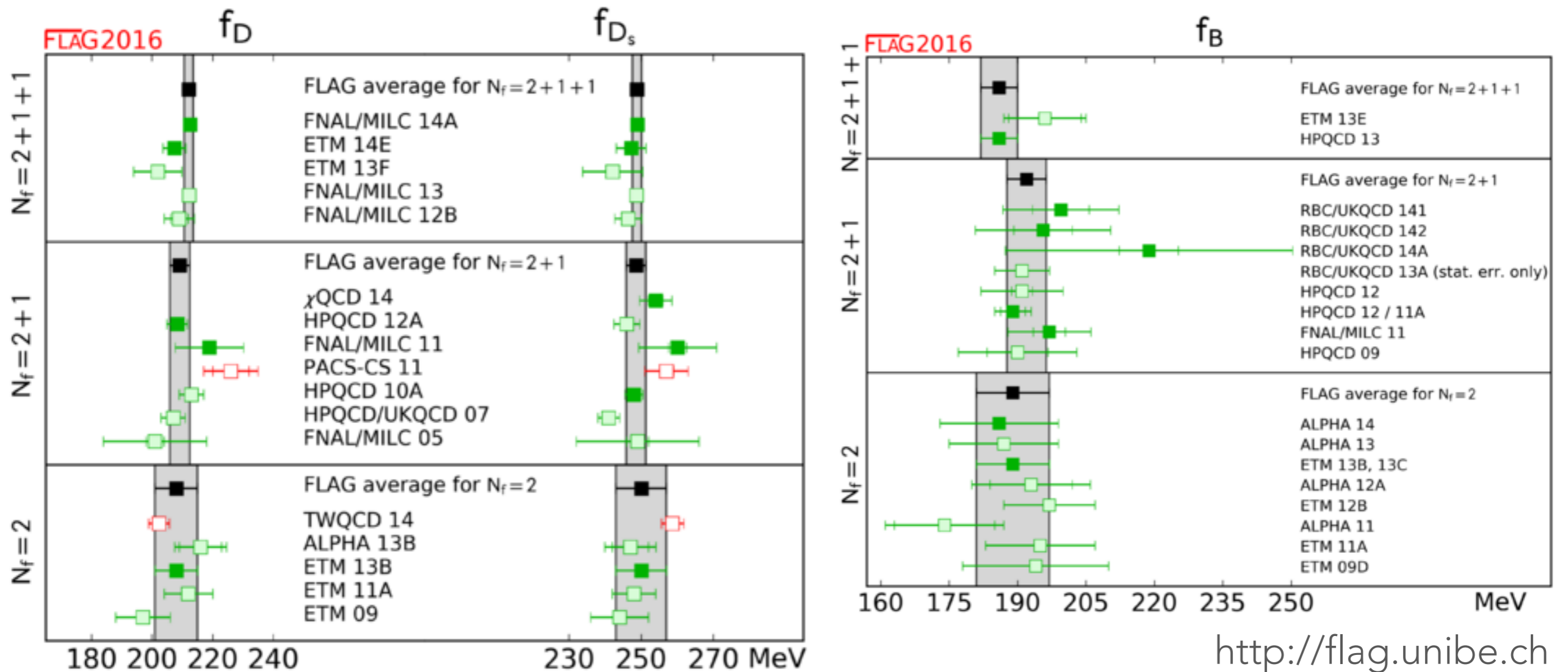
- 2** Lattice QCD + CKM
 Standard Model expectations for rare decays & mixing
 Target processes sensitive to beyond-SM physics

Many lattice calculations of what are now “simple” flavour physics quantities

FLAG: Flavour Lattice Averaging Group
Similar effort to the PDG, for Lattice QCD

- Members from most major lattice QCD collaborations
- Evaluates and grades different aspects of each calculation   
- Provides averages as the “Lattice QCD community consensus” value for a given quantity
- Includes lattice dictionary and summaries for non-experts
- Summary report every ~2 years: [arXiv:1105.3453] [1304.5422] [1607.00299]
Dec 2017 update at <http://flag.unibe.ch>
- **New version planned for 2019:** coverage expanding to include simple baryon quantities e.g., g_A

- Decay constants, form factors, kaon mixing, LECs...
- Colour coded for quality of calculation (# lattice spacings, volumes,...)



New results for $f_{B/D}$ from Fermilab lattice & MILC [arXiv:1712.09262] not yet in FLAG

$B \rightarrow D \ell \nu$ and $B \rightarrow D^* \ell \nu$

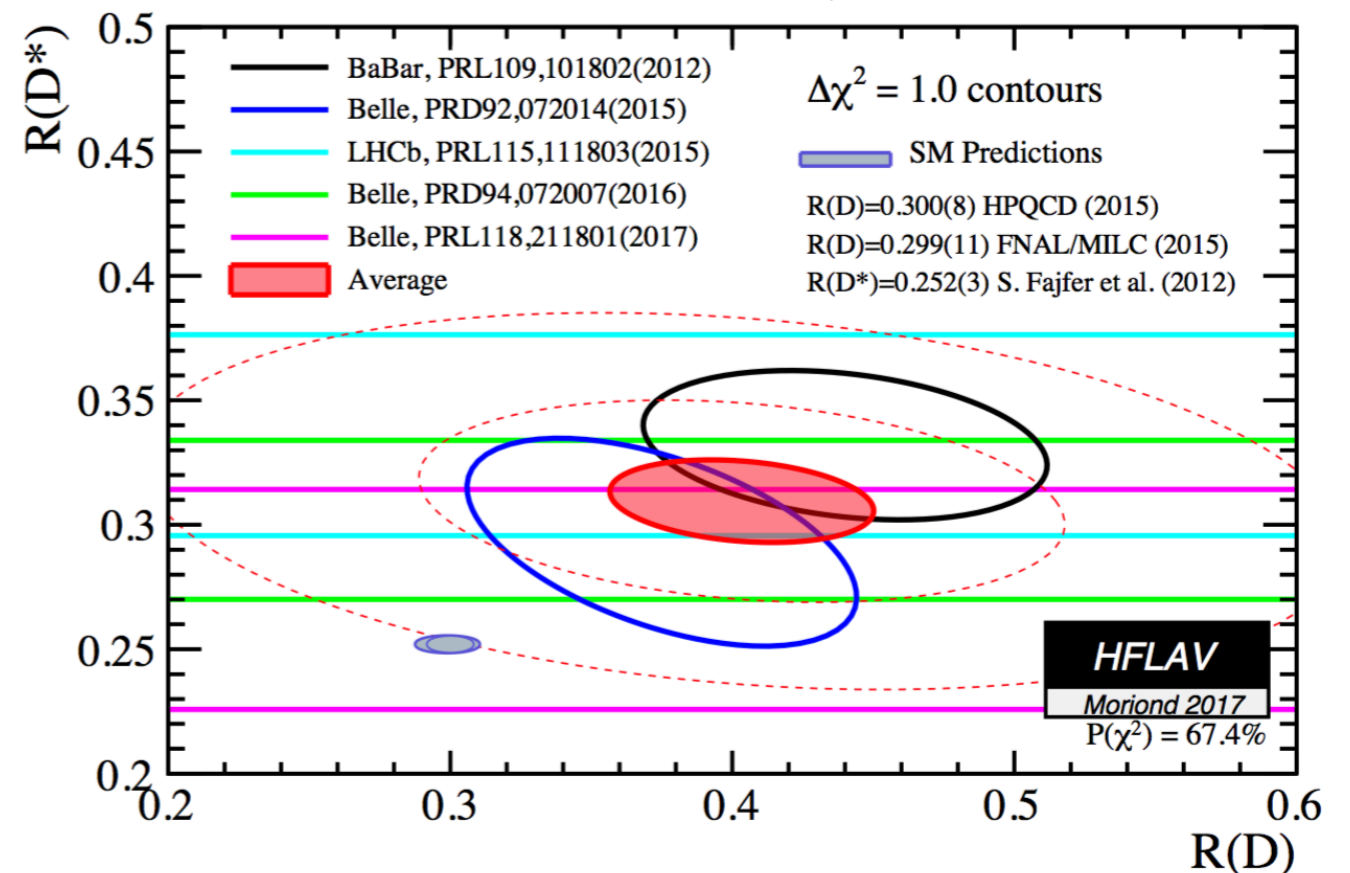
New lattice QCD results for $B \rightarrow D \ell \nu$ reveal $\sim 4\sigma$ tension with experimental measurements [HPQCD arXiv:1505.03925, MILC arXiv:1503.07237]

Belle 2 will provide further constraints

Standard Model prediction for $R(D^*)$ uses experimental data plus HQET: uncertainties may be underestimated [arXiv:1203.2654]

Lattice determinations of $\bar{B} \rightarrow D^* \ell \bar{\nu}$ form factors at nonzero recoil (in progress) will provide independent constraints on $R(D^*)$ [arXiv:1710.09817]

$$R(D^{(*)}) = \Gamma(B \rightarrow D^{(*)} \tau \nu) / \Gamma(B \rightarrow D^{(*)} \ell \nu)$$



Heavy Flavour Averaging Group [arXiv:1612.07233]

$$K^+ \rightarrow \pi^+ \nu \bar{\nu}$$

$K^+ \rightarrow \pi^+ \bar{\nu} \nu$ decay is a rare process sensitive to many new physics scenarios e.g., Custodial Randall-Sundrum, MSSM, Z, Z' models, Littlest Higgs with T-parity, LFU violation models

- NA62 Experiment aims to test Standard Model at 10% precision with data to end 2018; first dataset 2016, one event observed



- Short distance contribution to decay amplitude: perturbation theory + semileptonic kaon decay form factors
Long-distance contribution $O(5\%)$ from phenomenological estimates.

2018 Lattice QCD highlight:

Exploratory calculation demonstrated feasibility of decay amplitude calculation, in particular long-distance component [Bai et al., arXiv:1806.11520]

Expectation that a fully controlled calculation will be possible within four years

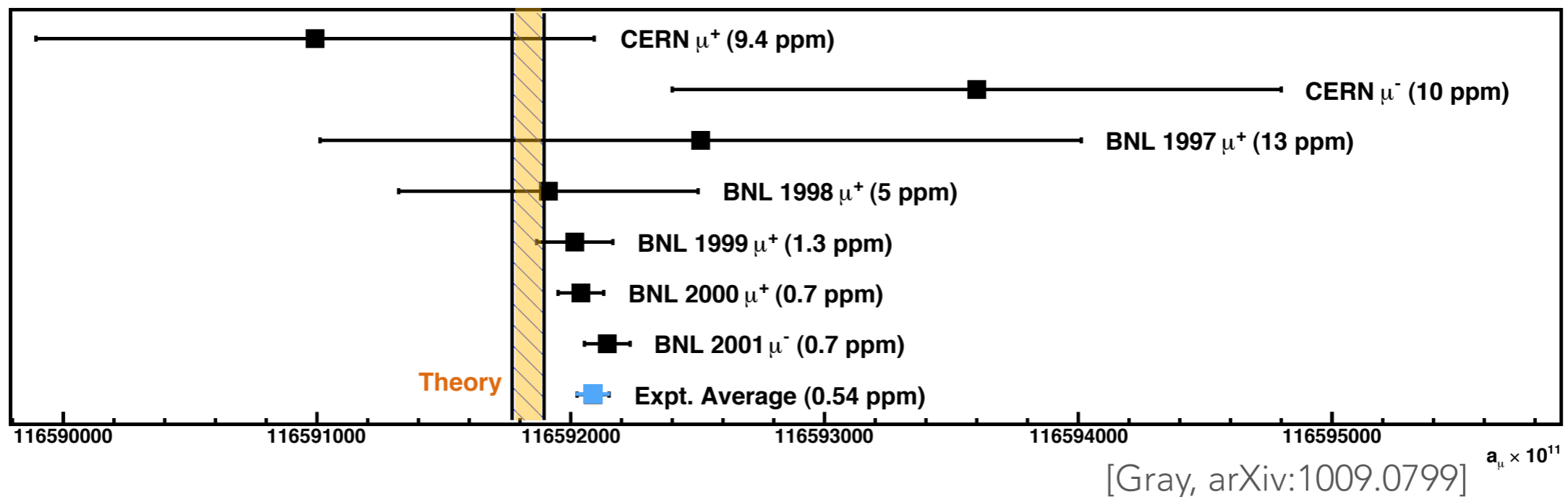
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Muon $g-2$: status and challenges

Long standing discrepancy between measured value and SM estimate for muon anomalous magnetic moment ($\sim 3\sigma$)



Sign of new physics?

New experiments aiming at 4-fold uncertainty reduction (E989 @ Fermilab, E34 @ JPARC)

➔ if no shift in central values, tension will be $\sim 7\sigma$ with projected theory improvements by 2020

Commensurate control of theory needed: Muon $g-2$ Theory Initiative formed

<https://indico.fnal.gov/event/13795/>

Standard Model muon g-2

Measured value

$$a_{\mu}^{\text{E821}} = (116\,592\,089 \pm 63) \times 10^{-11}$$

Breakdown of SM contributions
(2 evaluations of HVP)

	VALUE ($\times 10^{-11}$) UNITS
QED ($\gamma + \ell$)	$116\,584\,718.951 \pm 0.009 \pm 0.019 \pm 0.007 \pm 0.077_{\alpha}$
HVP(lo) [20]	$6\,923 \pm 42$
HVP(lo) [21]	$6\,949 \pm 43$
HVP(ho) [21]	-98.4 ± 0.7
HLbL	105 ± 26
EW	154 ± 1
Total SM [20]	$116\,591\,802 \pm 42_{\text{H-LO}} \pm 26_{\text{H-HO}} \pm 2_{\text{other}} (\pm 49_{\text{tot}})$
Total SM [21]	$116\,591\,828 \pm 43_{\text{H-LO}} \pm 26_{\text{H-HO}} \pm 2_{\text{other}} (\pm 50_{\text{tot}})$

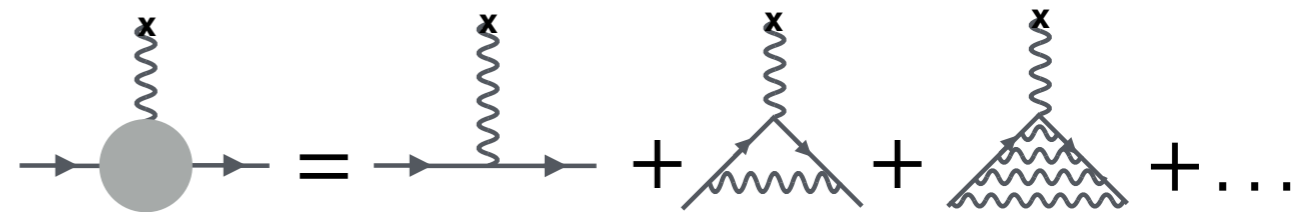
[T. Blum et al., arXiv:1311.2198]

Deviation

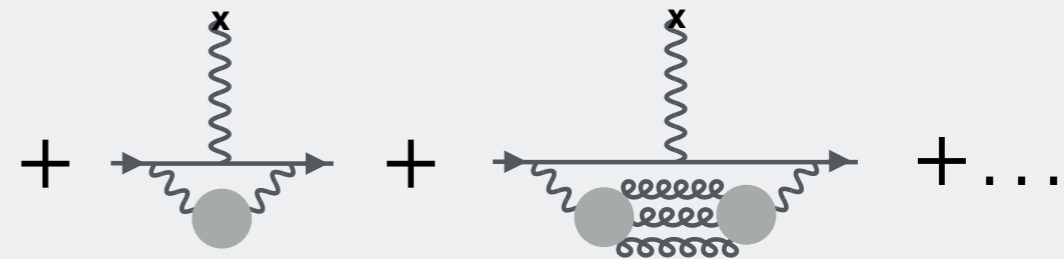
$$\begin{aligned} \Delta a_{\mu}(\text{E821} - \text{SM}) &= (287 \pm 80) \times 10^{-11} \text{ [20]} \\ &= (261 \pm 78) \times 10^{-11} \text{ [21]} \end{aligned}$$

Dominant uncertainties from hadronic corrections—calculable in lattice QCD

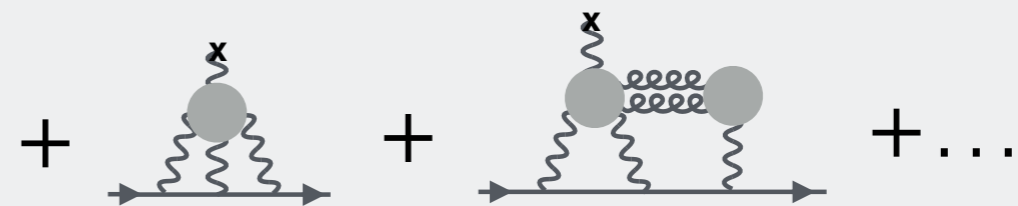
QED (5 loop) [Aoyama et al. 2012]



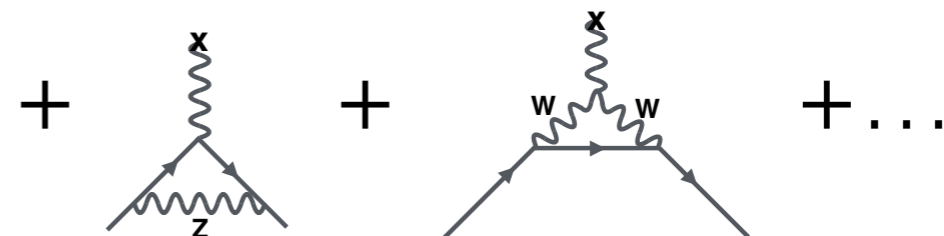
Hadronic vacuum polarisation



Hadronic light-by-light



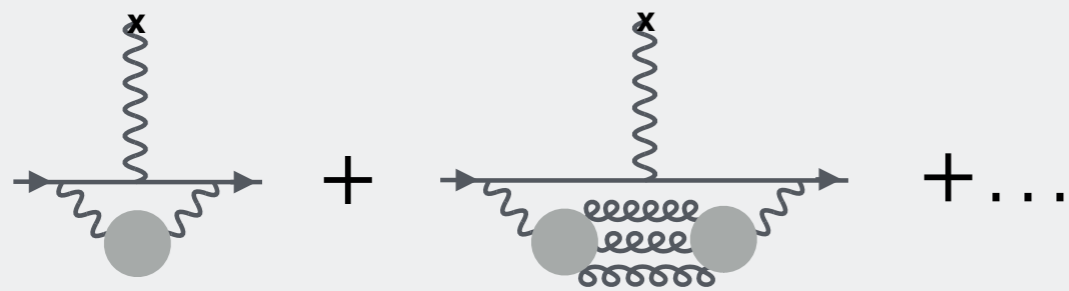
Electroweak (2 loop) [Czarnecki et al. 2006]



Standard Model muon $g-2$

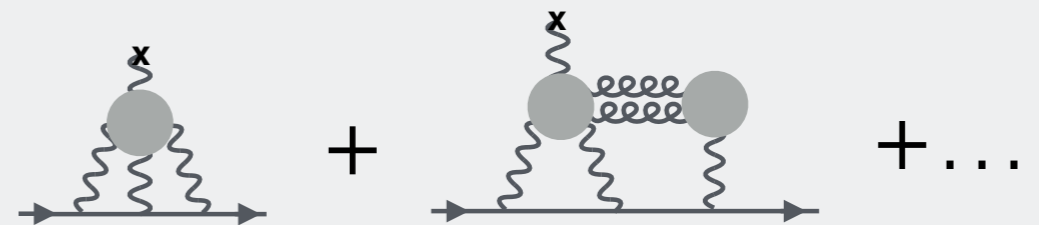
Dominant uncertainties in SM determination from hadronic corrections
Both calculable (in principle) from lattice QCD

Hadronic vacuum polarisation



Dispersion relation plus
experimental data on “R-ratio”
 $\sigma(e+e^- \rightarrow \text{hadrons})$

Hadronic light-by-light



Estimated from models
including large- N_c , chiPT,
vector meson dominance, etc.

2018: First lattice QCD calculation
with QED and isospin breaking
[T. Blum et al., arXiv:1801.07224]

Since ICHEP2016: disconnected terms
and lattice volume better controlled
[PRL118(2016)022005, PRD96(2017)034515]

Hadronic Vacuum Polarisation

Combining lattice QCD and dispersion relations yields best current determination of HVP contribution

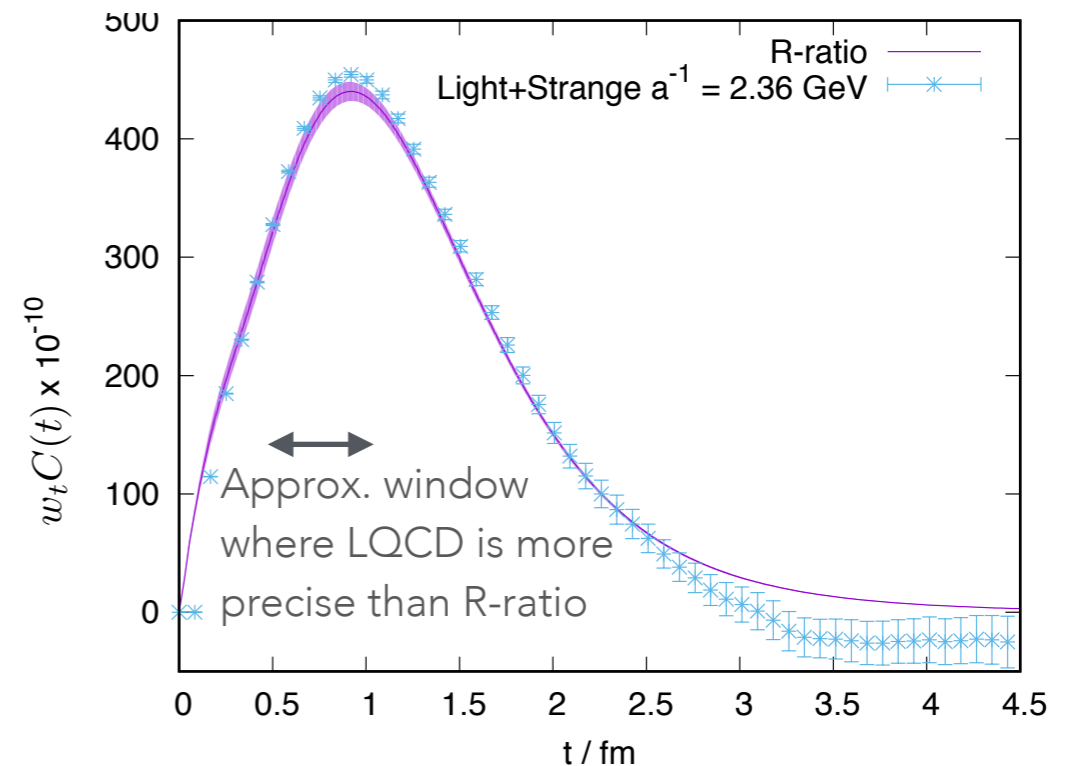
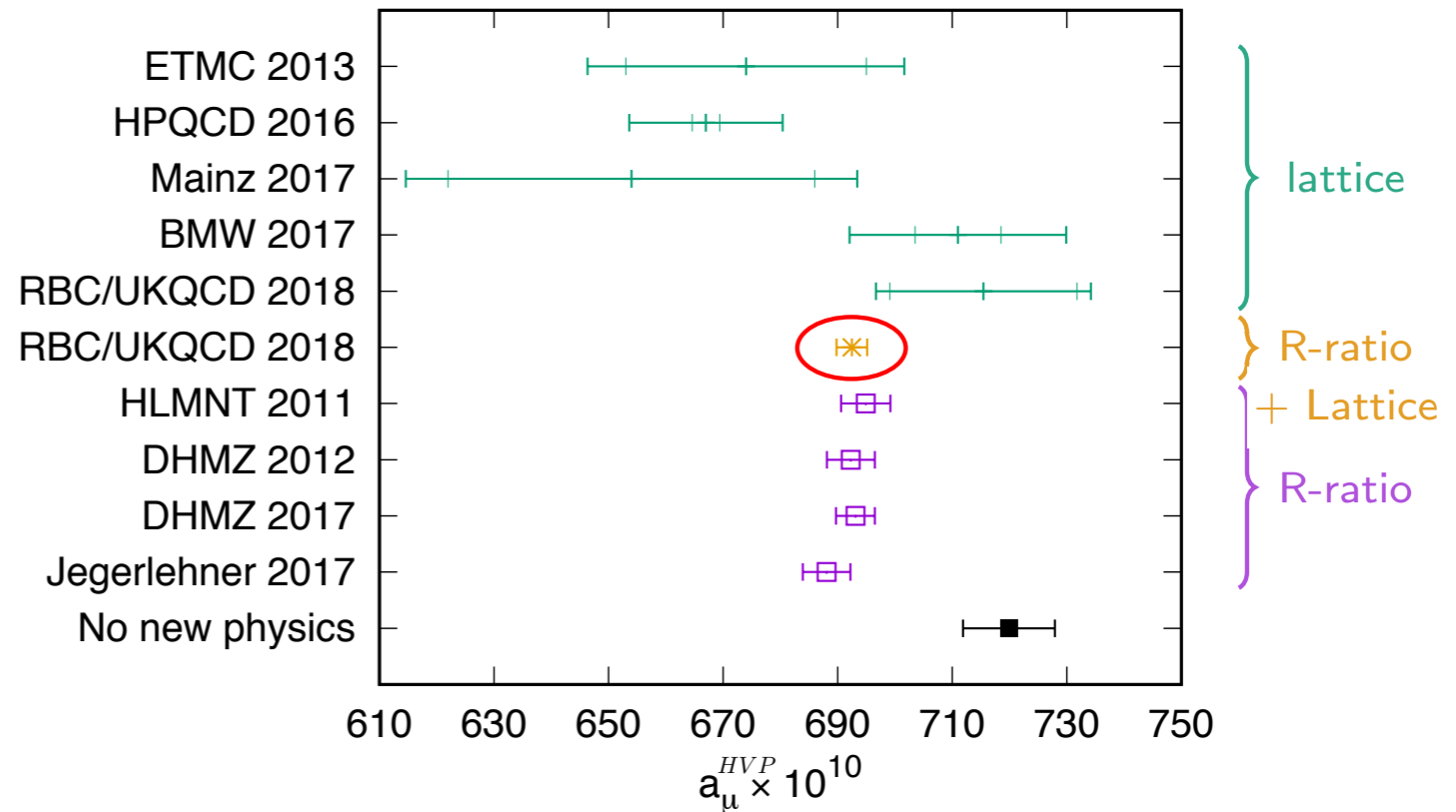
[T. Blum et al., arXiv:1801.07224]

Use "R-ratio" experimental data on $\sigma(e+e-\rightarrow\text{hadrons})$ at short and long distances t , lattice QCD in intermediate t region

$$a_{\mu}^{HVP} = \sum_t w_t C(t)$$

$$C(t) = \frac{1}{3} \sum_{\vec{x}} \sum_{j=0,1,2} \langle J_j(\vec{x}, t) J_j(0) \rangle$$

Flavour breakdown: light~90%, strange~8% and charm~2%



Highlights since ICHEP2016

- Quark masses and α_s
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Parton distribution functions

Parton distribution functions $f(x, \mu^2)$

Number densities of partons of type f with momentum fraction x at scale μ^2 in a given hadron

PDFs quantify fundamental aspects of hadron structure

Nucleon PDFs are needed for e.g., searches for new physics at the LHC through top-quark and Higgs-boson coupling measurements

Lattice QCD can provide

- Moments of PDFs with controlled uncertainties: $\int_0^1 x^n f(x, \mu^2) = \langle x^n \rangle_f(\mu^2)$
 - ➔ Inclusion in global PDF fits can reduce uncertainties
see workshop slides <http://www.physics.ox.ac.uk/confs/PDFlattice2017>
and community white paper [Prog.Part.Nucl.Phys.100 (2018) 107]
- First calculations of x -dependence of nucleon PDFs

Moments of PDFs

Lattice QCD can cleanly access low moments of PDFs ($n \lesssim 3$)

[work to move beyond: Chambers et al., arXiv:1703.01153, Davoudi & Savage, arXiv:1204.4146]

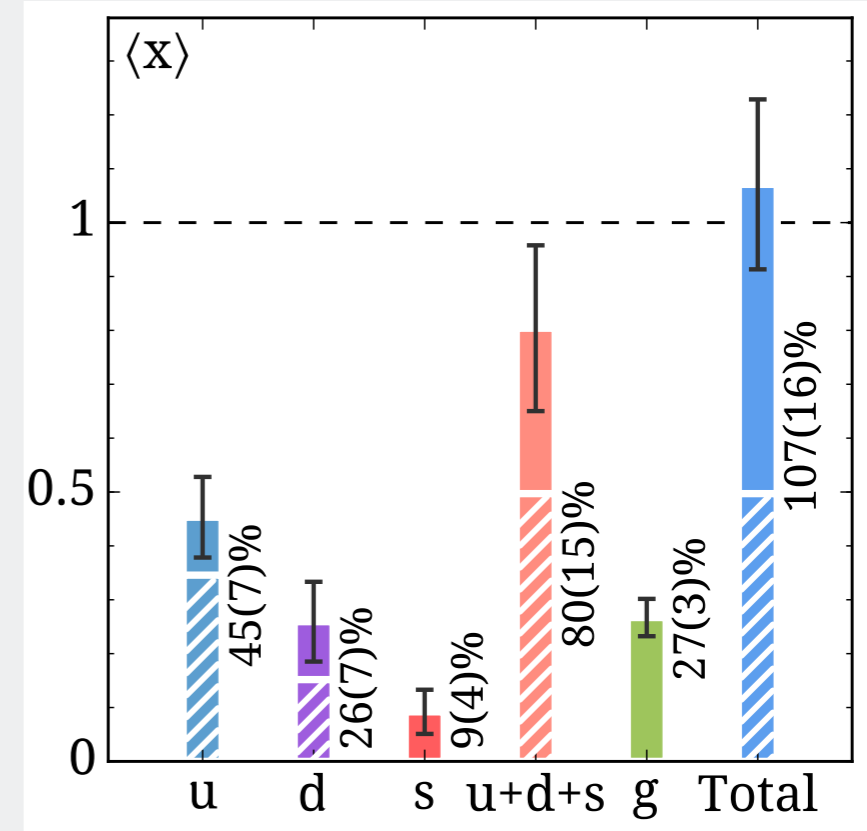
$$\int_0^1 x^n f(x, \mu^2) = \langle x^n \rangle_f(\mu^2)$$

State-of-the-art calculations have:

- Fully-controlled systematic uncertainties competitive with or better than experiment for some quantities
- Separate contributions from
 - Strangeness and light flavours
 - Charge symmetry violation
 - Gluons

2017 Highlight: All terms of nucleon momentum decomposition calculated with controlled uncertainties

$\overline{\text{MS}}$ -scheme at 2 GeV

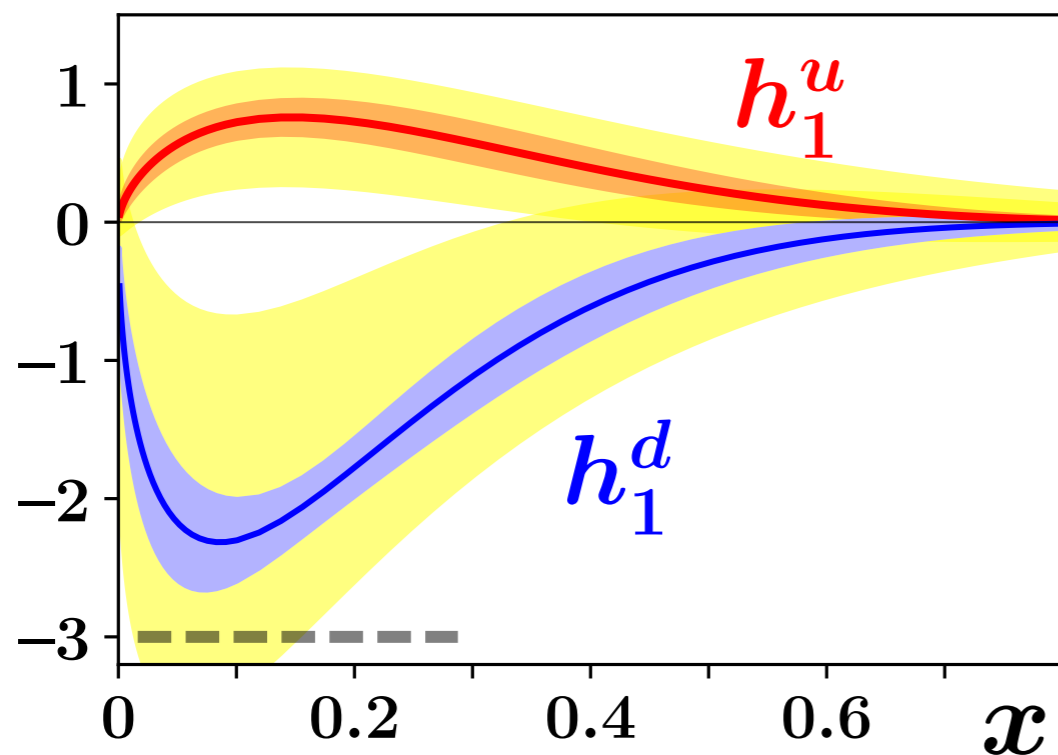


[C. Alexandrou et al., arXiv:1706.02973]

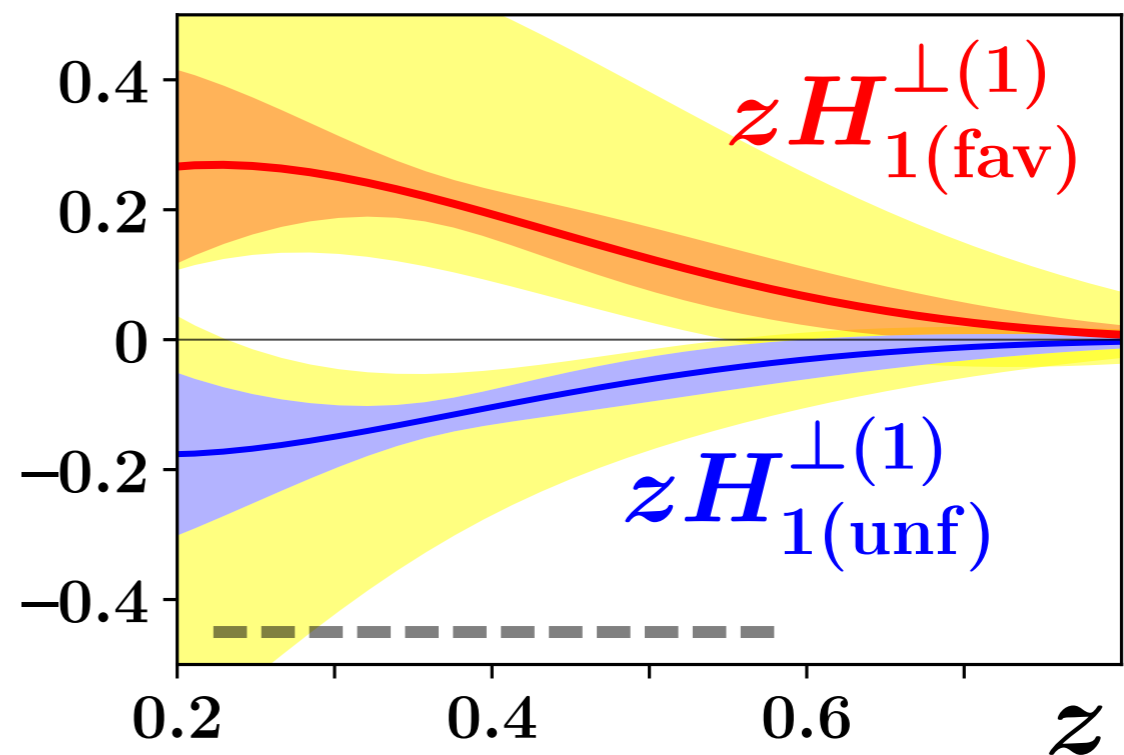
Constraints on global PDF fits

- Including lattice QCD results for moments in global PDF fits can yield significant improvements

Transversity PDFs



Collins fragmentation functions

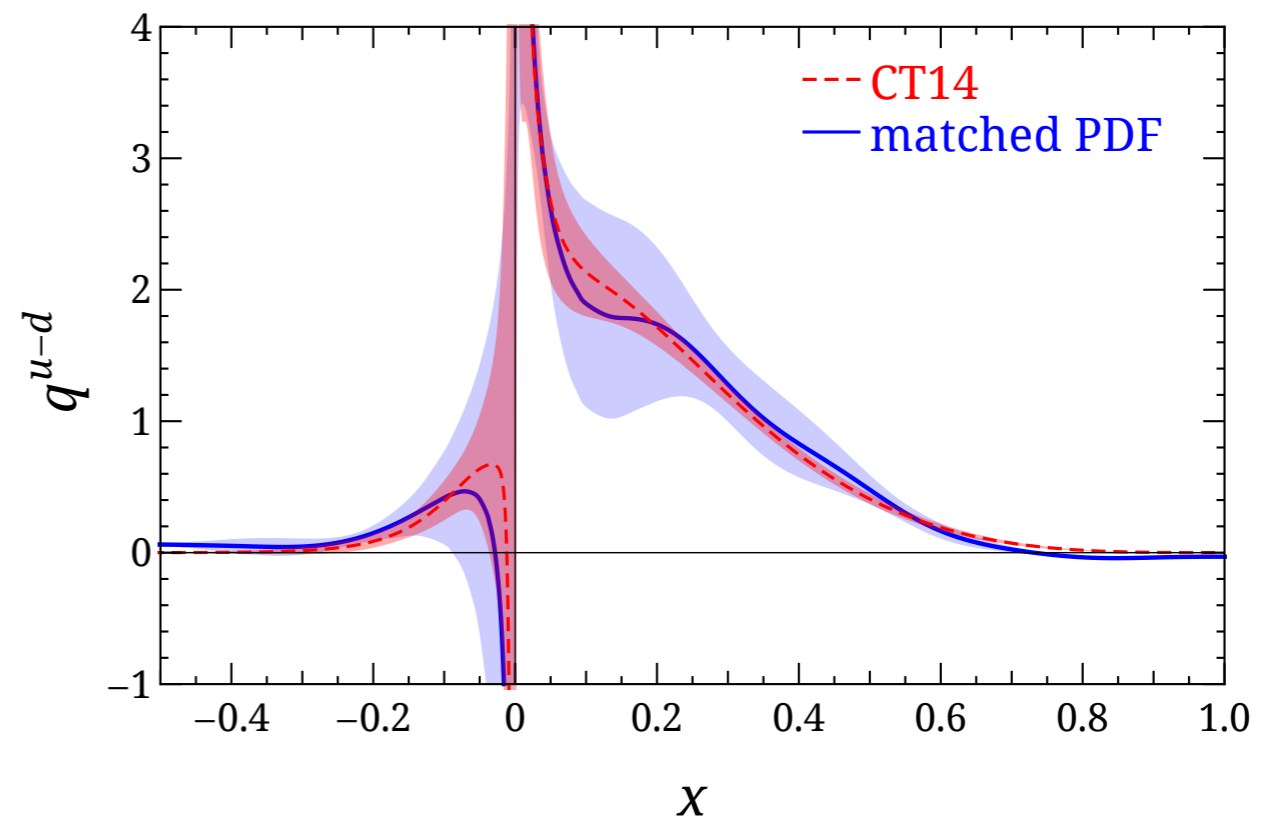
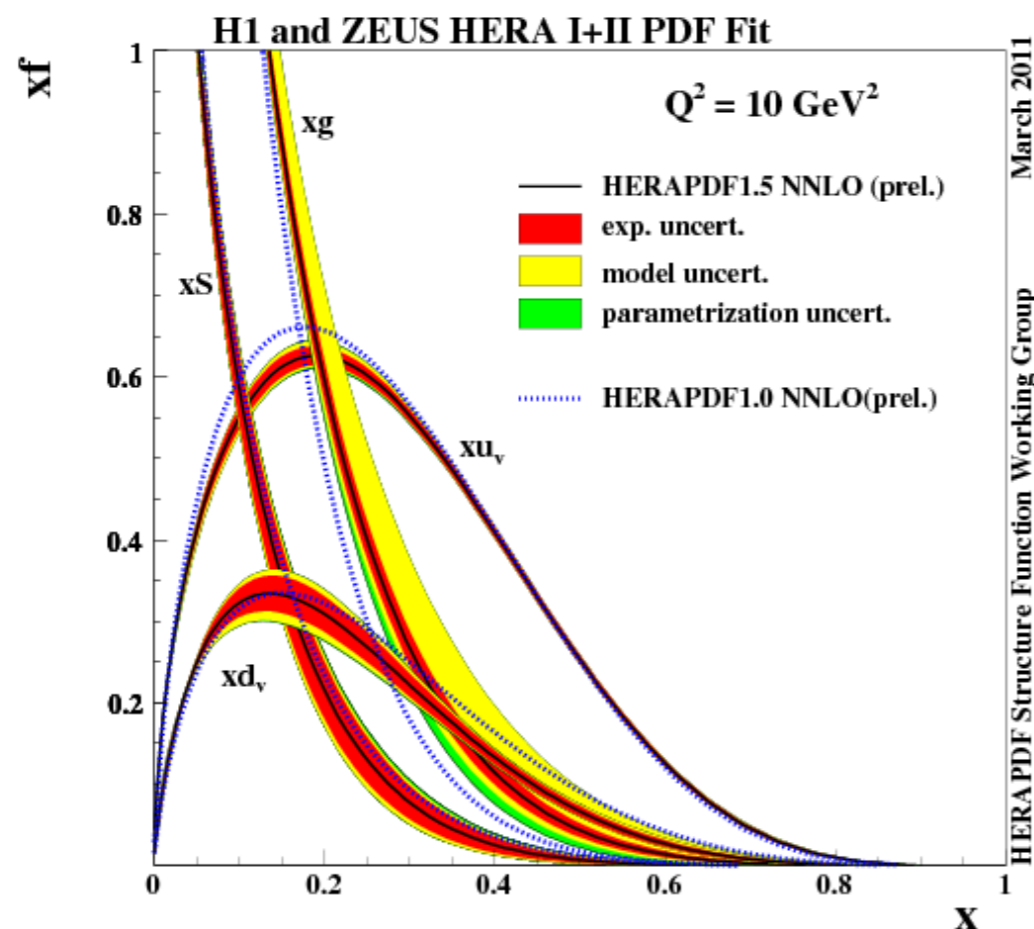


Yellow: SIDIS data only: direct constraints in region indicated by dashes
Blue/Red: SIDIS + lattice QCD for tensor charge (zeroth moment)

[H-W. Lin et al., arXiv:1710.09858]

x-dependence of PDFs

- First calculations of x-dependence of nucleon PDFs undertaken
Quasi and pseudo-PDF calculations use non-local Euclidean correlators and perturbative QCD matching in high momentum limit [X. Ji, arXiv:1305.1539]
- Extremely rapid progress, but many systematics to be controlled
- Flavour separation is relatively straightforward



[J-W Chen et al., arXiv:1803.04393]

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Long-baseline neutrino experiments



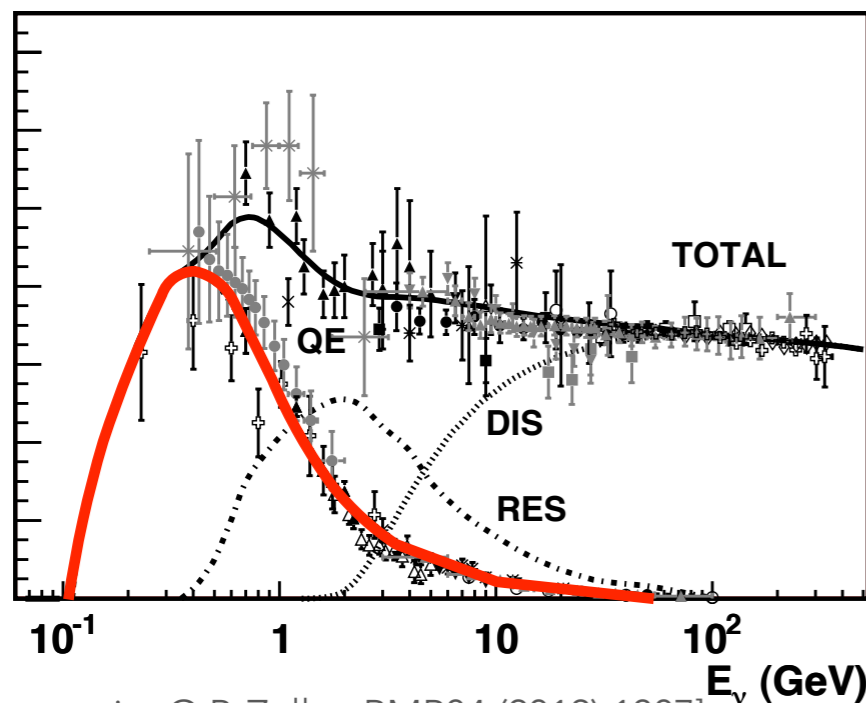
Seek to determine neutrino mass hierarchy, mixing parameters, CP violating phase

To differentiate between mixing & CP parameter scenarios



Need neutrino energy reconstruction from final state to better than 100 MeV

ν charged-current cross-section



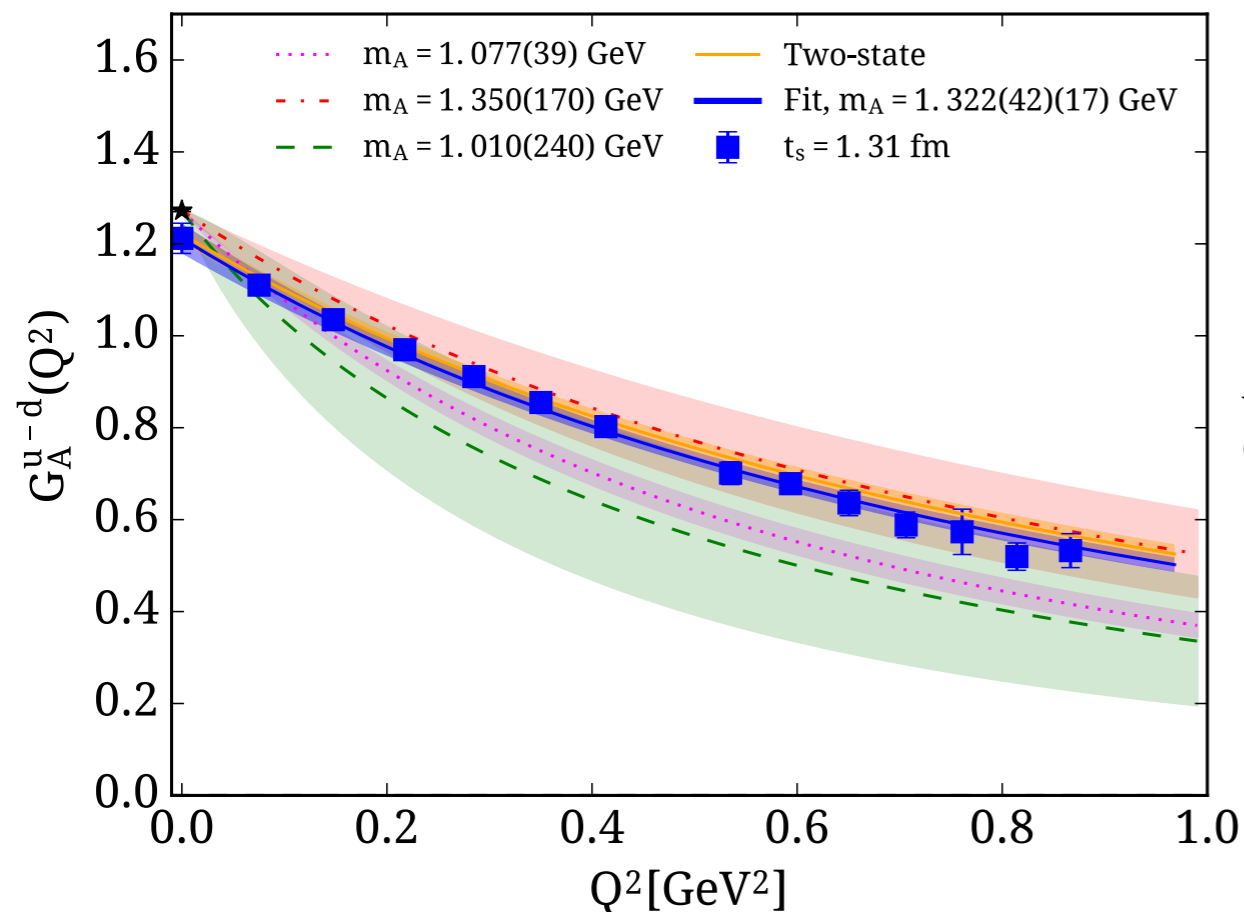
[J.A. Formaggio, G.P. Zeller, RMP84 (2012) 1307]

Lattice QCD: direct non-perturbative QCD predictions for nucleon and nuclear matrix elements

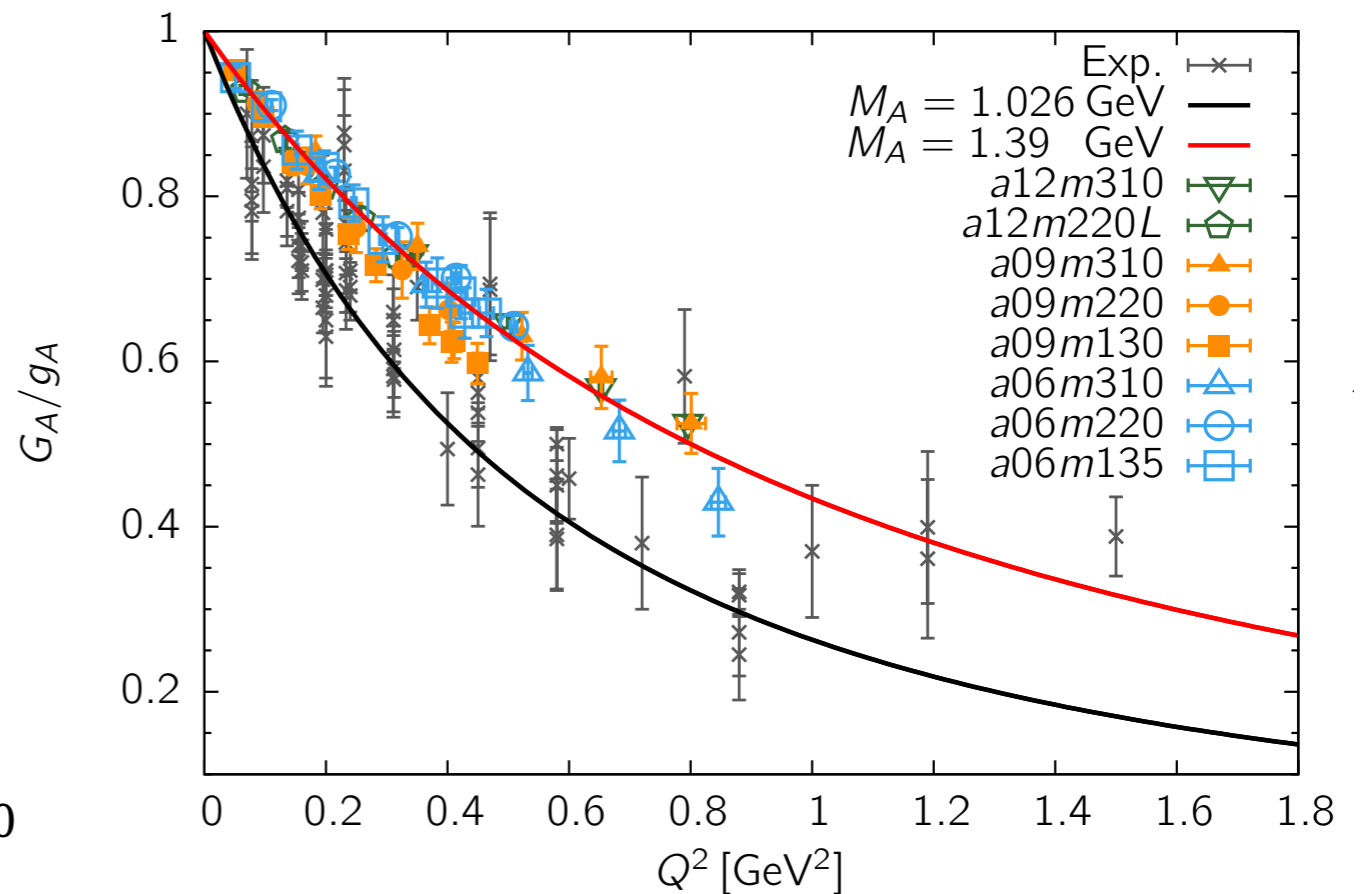
e.g., axial and pseudo-scalar form factors important in **quasi-elastic** region (Lattice QCD can also provide useful input in resonance and DIS regions, see appendix)

Nucleon axial form factors

- Nucleon properties are historically difficult calculations
- Recent calculations of nucleon form factors including axial in agreement with experiment with fully-controlled uncertainties
- Q^2 -dependence well-determined in LQCD: competitive with experiment



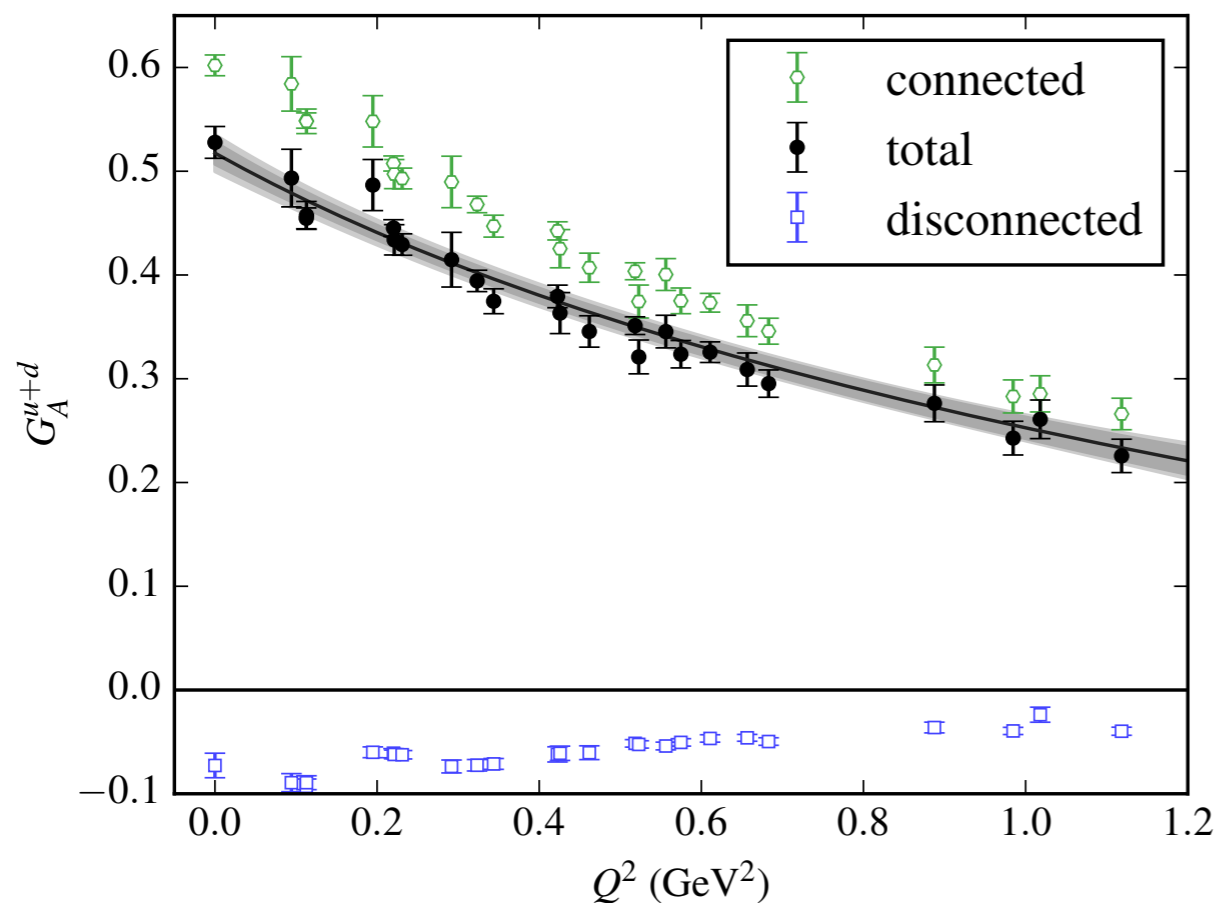
[Alexandrou et al., PRD96 (2017), 054507]



[Gupta et al., PRD96 (2017), 114503]

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[Green et al., PRD 95, 114502 (2017)]

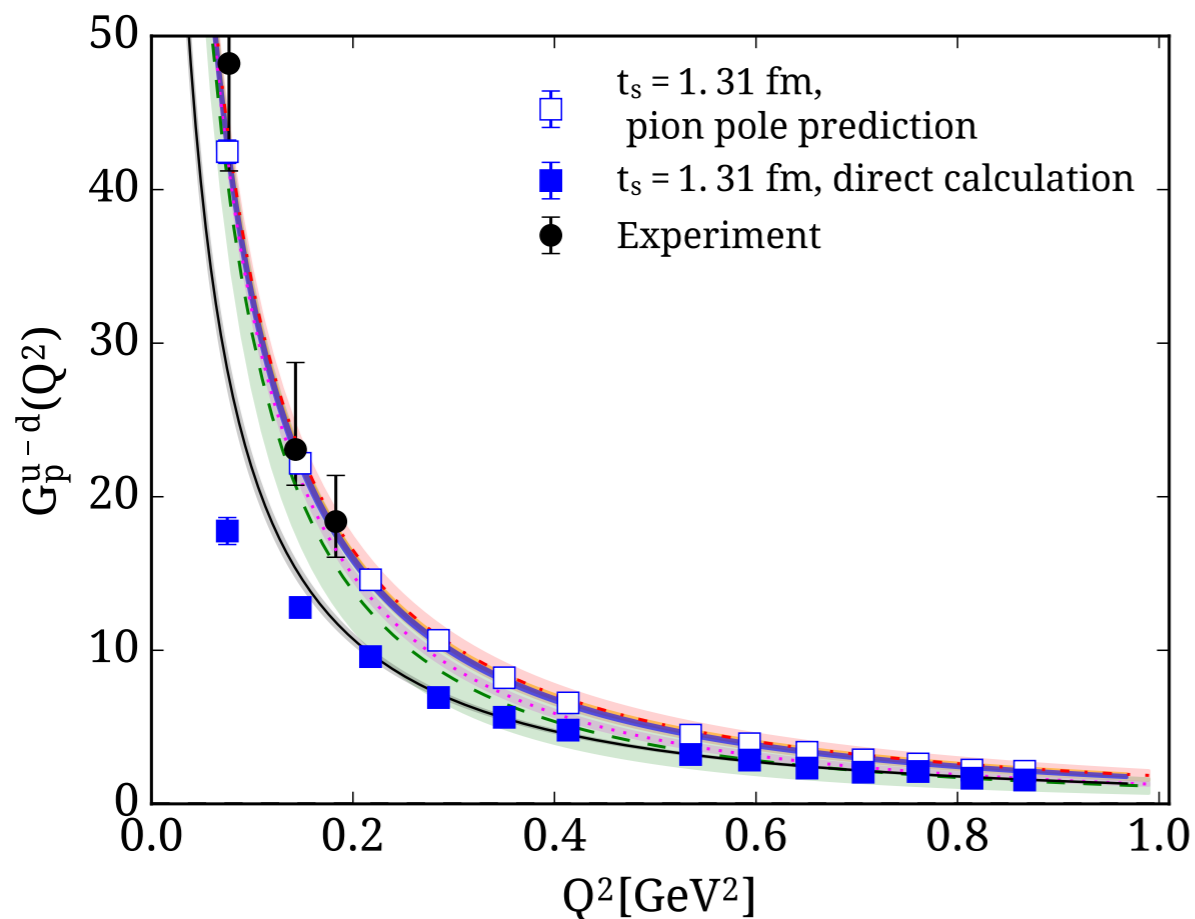
Strange quark contributions are determined separately and can be isolated

[Also Gupta et al., EPJ Web Conf. 175 (2018) 06029, Alexandrou et al., PRD96 (2017), 054507]

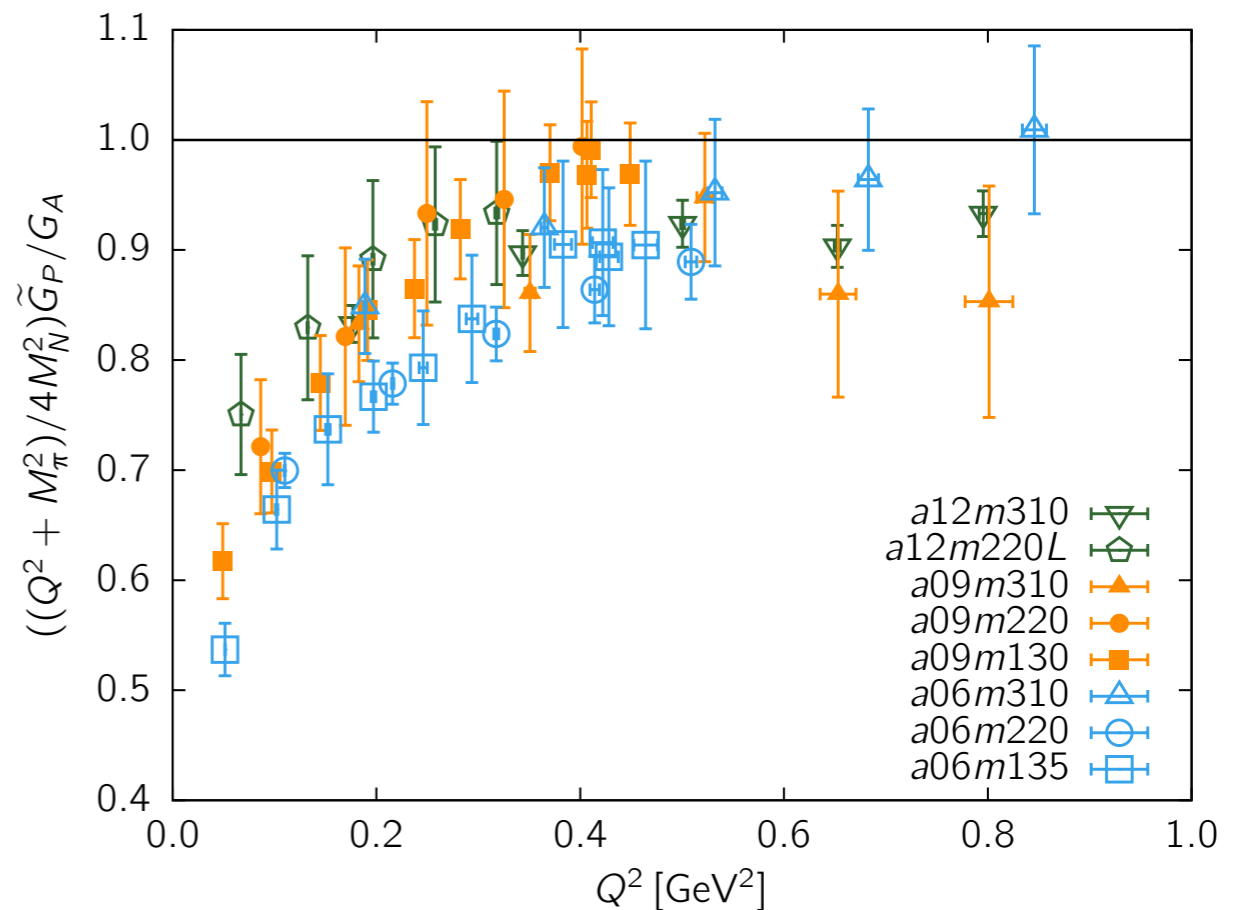
Nucleon pseudo-scalar form factors

- First calculations with controlled uncertainties
- Clear deviations from pion-pole dominants ansatz at low Q^2

$$\tilde{G}_P(Q^2) = G_A(Q^2) \left[\frac{4M_N^2}{Q^2 + M_\pi^2} \right]$$



[Alexandrou et al., PRD96 (2017), 054507]



[Gupta et al., PRD96 (2017), 114503]

Highlights since ICHEP2016

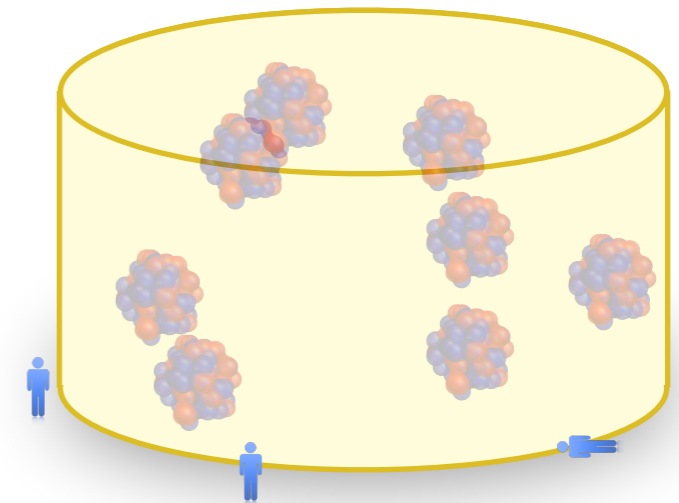
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Dark matter direct detection

Look for scattering of WIMP dark matter on nuclear target

Detection rate depends on

- Dark matter properties
- Probability of interaction with nucleus
i.e., *nuclear* effects are important



Low-energy limit of a generic
spin-independent interaction
is scalar



Determine **nucleon and
nuclear scalar matrix
elements** from lattice QCD

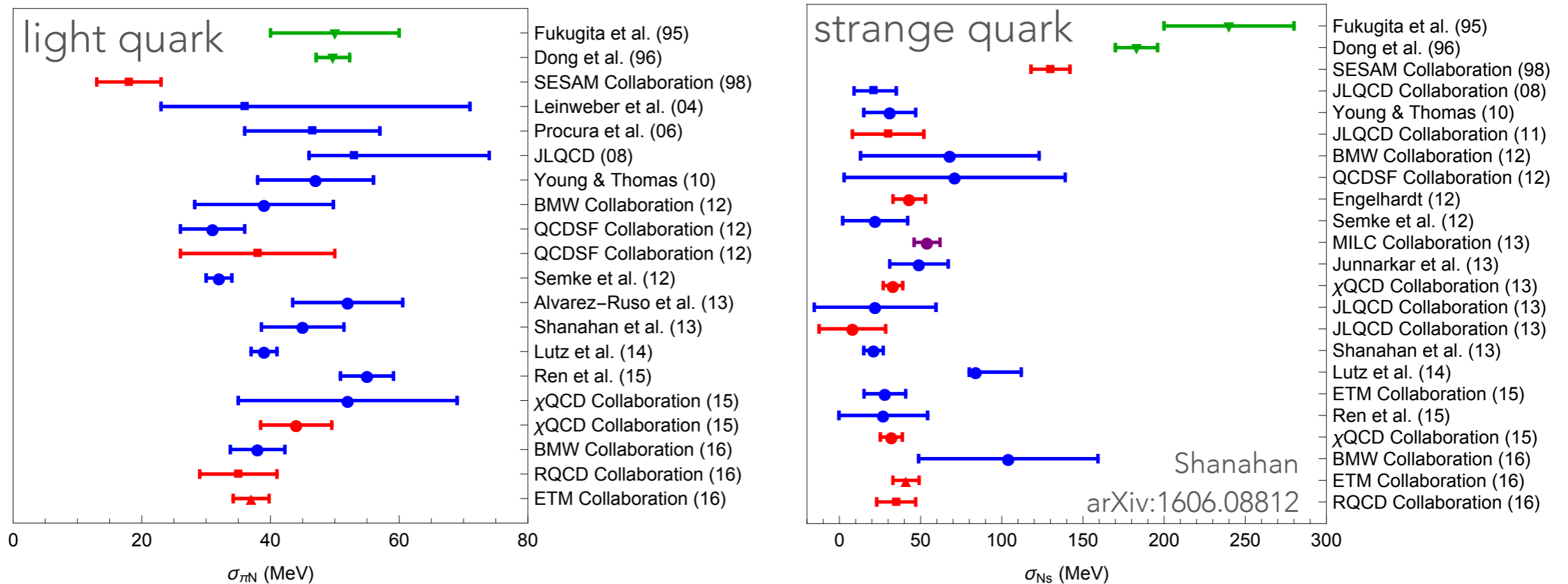
Other e.g., spin-dependent couplings can also be constrained

e.g., [Hoferichter et al., arXiv:1503.04811], [Hill et al., arXiv:1409.8290], [Fitzpatrick et al., arXiv:1203.3542]

Dark matter direct detection

Spin-independent scattering of many WIMP candidates governed by scalar matrix elements

Lattice QCD nucleon scalar matrix elements



Light quark: competitive with phenomenology

Note: tension with extraction using Roy–Steiner equations [Hoferichter arXiv:1506.04142]

Strange quark: much more precise than phenomenology

Dark matter direct detection

Direct detection experiments use nuclear targets e.g., Xenon

Determine interaction cross-section (with nucleus) for a given dark matter model

- Born approximation – interacts with a single nucleon

$$\sigma \sim |A \langle N | DM | N \rangle|^2$$

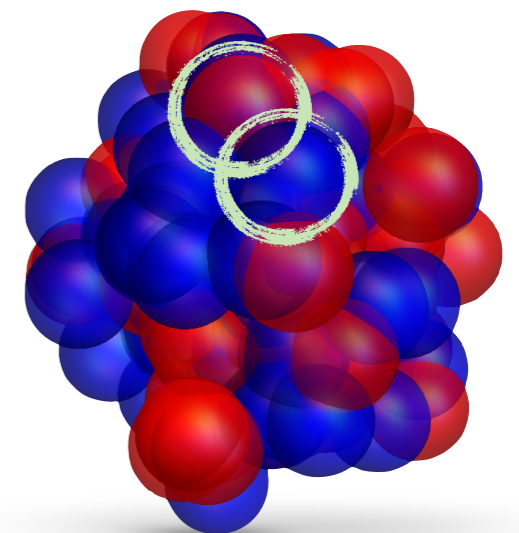
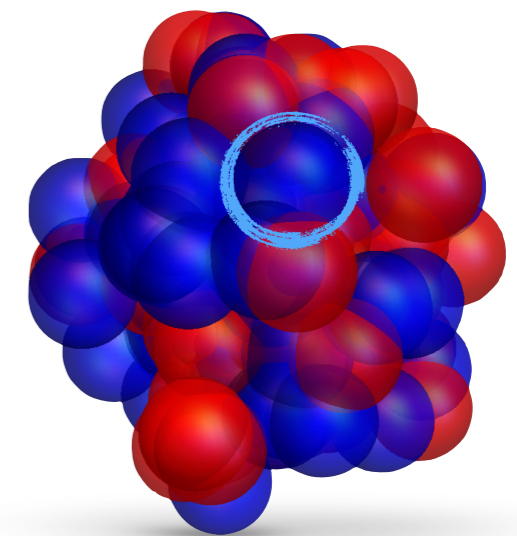
known from LQCD

- Interacts non-trivially with multiple nucleons

$$\sigma \sim |A \langle N | DM | N \rangle + \alpha \langle NN | DM | NN \rangle + \dots|^2$$

poorly known!

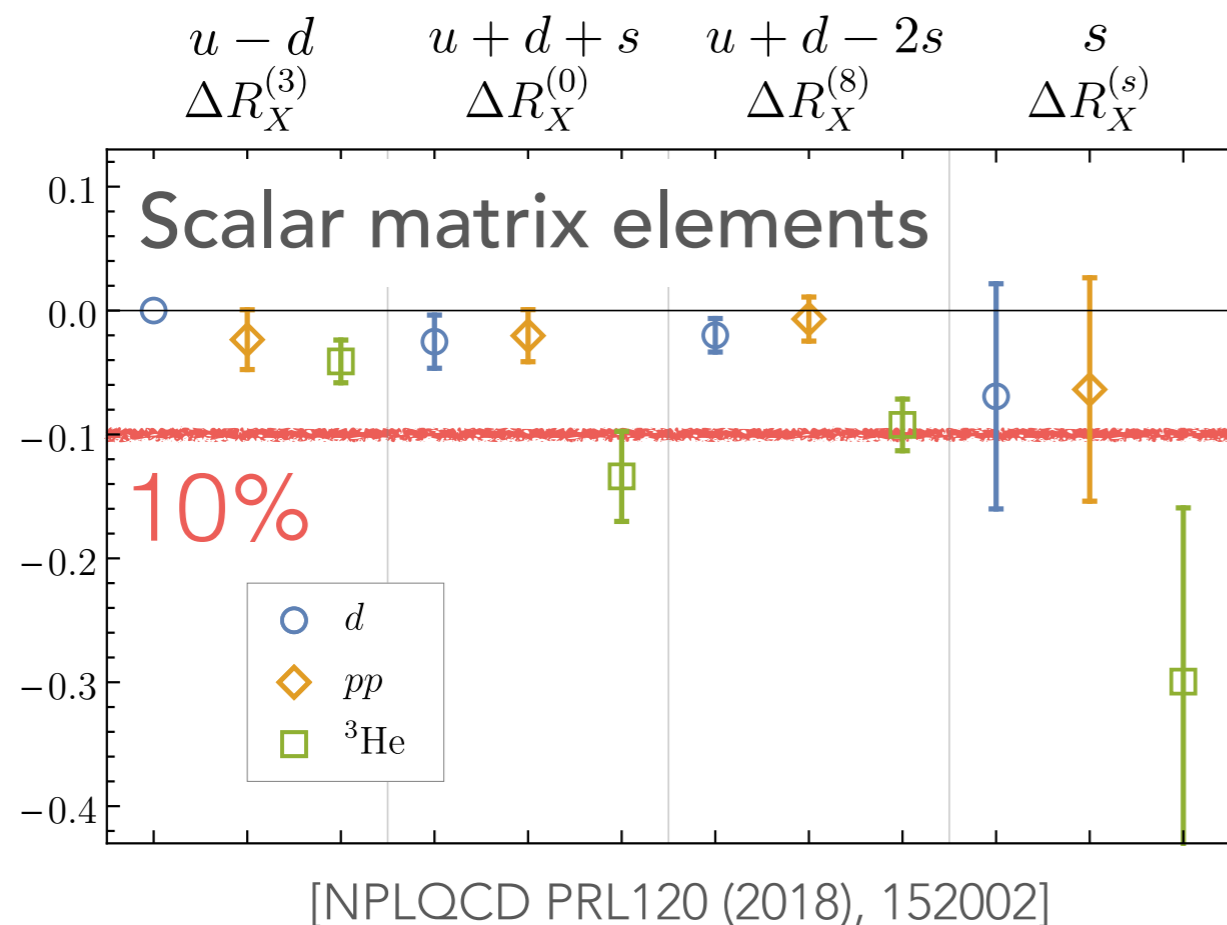
Second term may be significant!



Dark matter direct detection

Spin-independent scattering of WIMP candidates is governed by scalar matrix elements

- Lattice QCD calculation with $m_\pi \sim 800$ MeV shows 10% nuclear effects in ${}^3\text{He}$ \longrightarrow potentially very significant effects in e.g., Xenon
- Same calculation gives axial and tensor nuclear effects around $\sim 1\%$

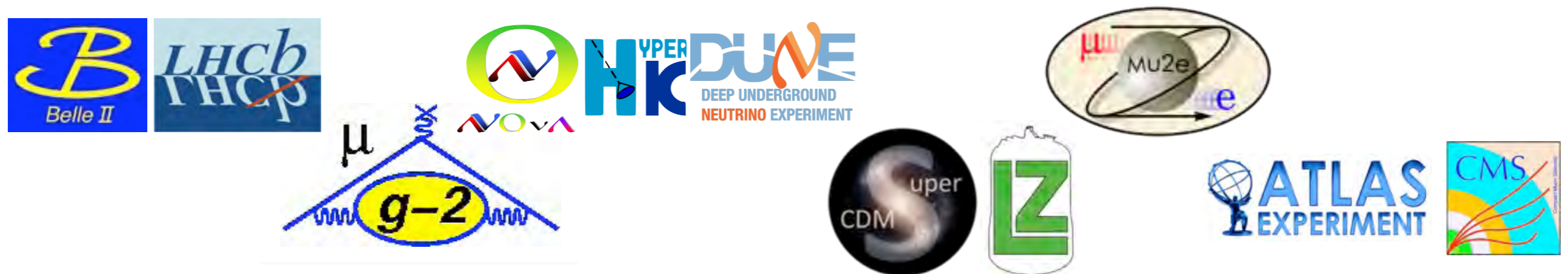


$$\frac{\text{ME}}{A(\text{Nucleon ME})} \sim 1$$

Older work: [NPLQCD PRD89 (2014) 074505]

Summary and outlook

Lattice QCD (+QED) is providing essential Standard Model input for high-energy physics experiments



- Precision lattice QCD results for simple systems including hadrons
 - ➔ FLAG lattice averaging to include hadron structure in 2019
 - ➔ Many more quantities will have fully-controlled systematic uncertainties by 2020
- Beginning of reliable lattice QCD results for nuclear matrix elements

Summary and outlook

Lattice QCD (+QED) is providing essential Standard Model input for high-energy physics experiments

... and also for many other topics not covered here

- Hadron spectroscopy
- Finite temperature and density
- Nuclear structure and reactions
- BSM physics (Technicolour theories, SUSY, ...)

New results will be presented at Lattice conference July 22-28: see conference website and proceedings



East Lansing, MI, USA

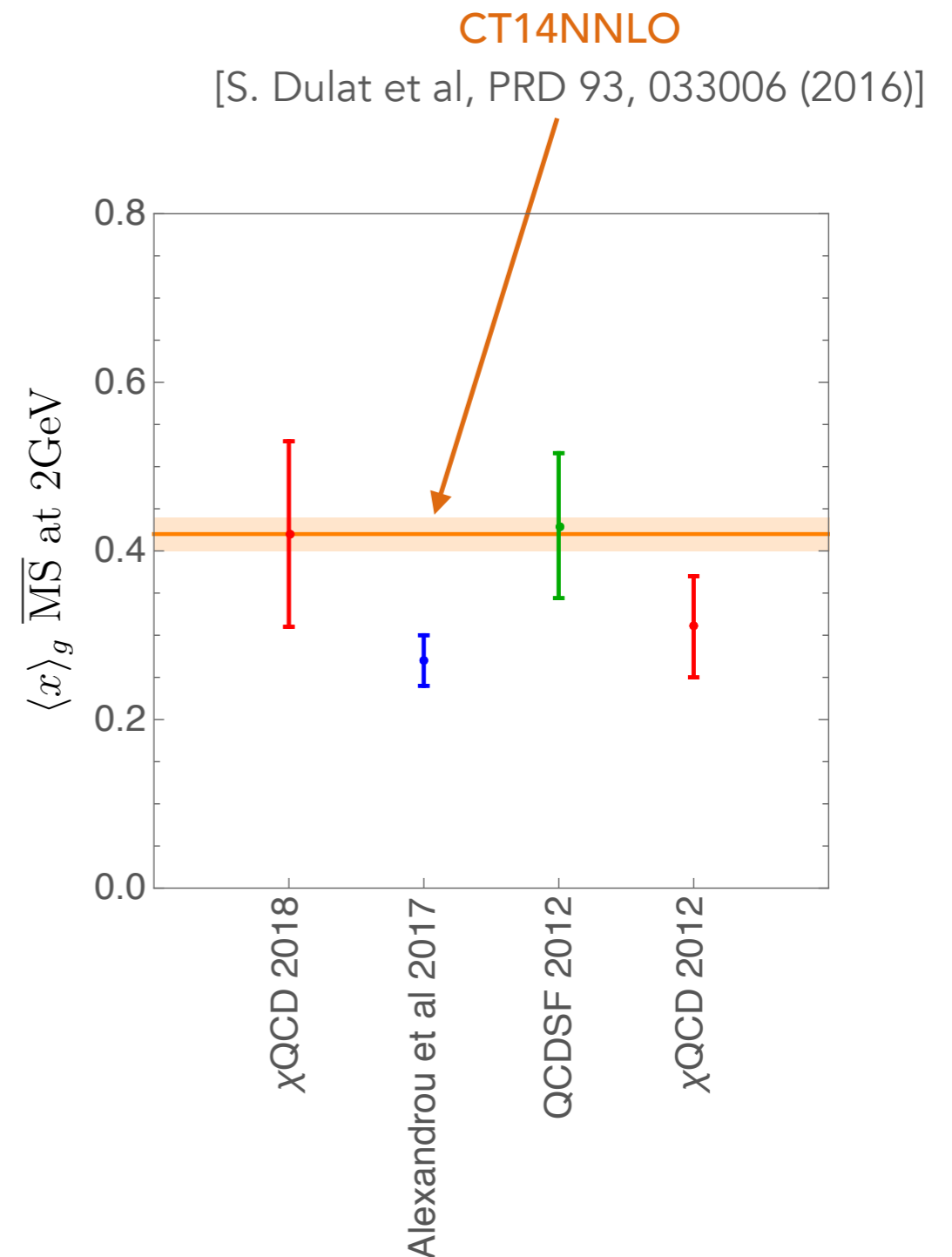
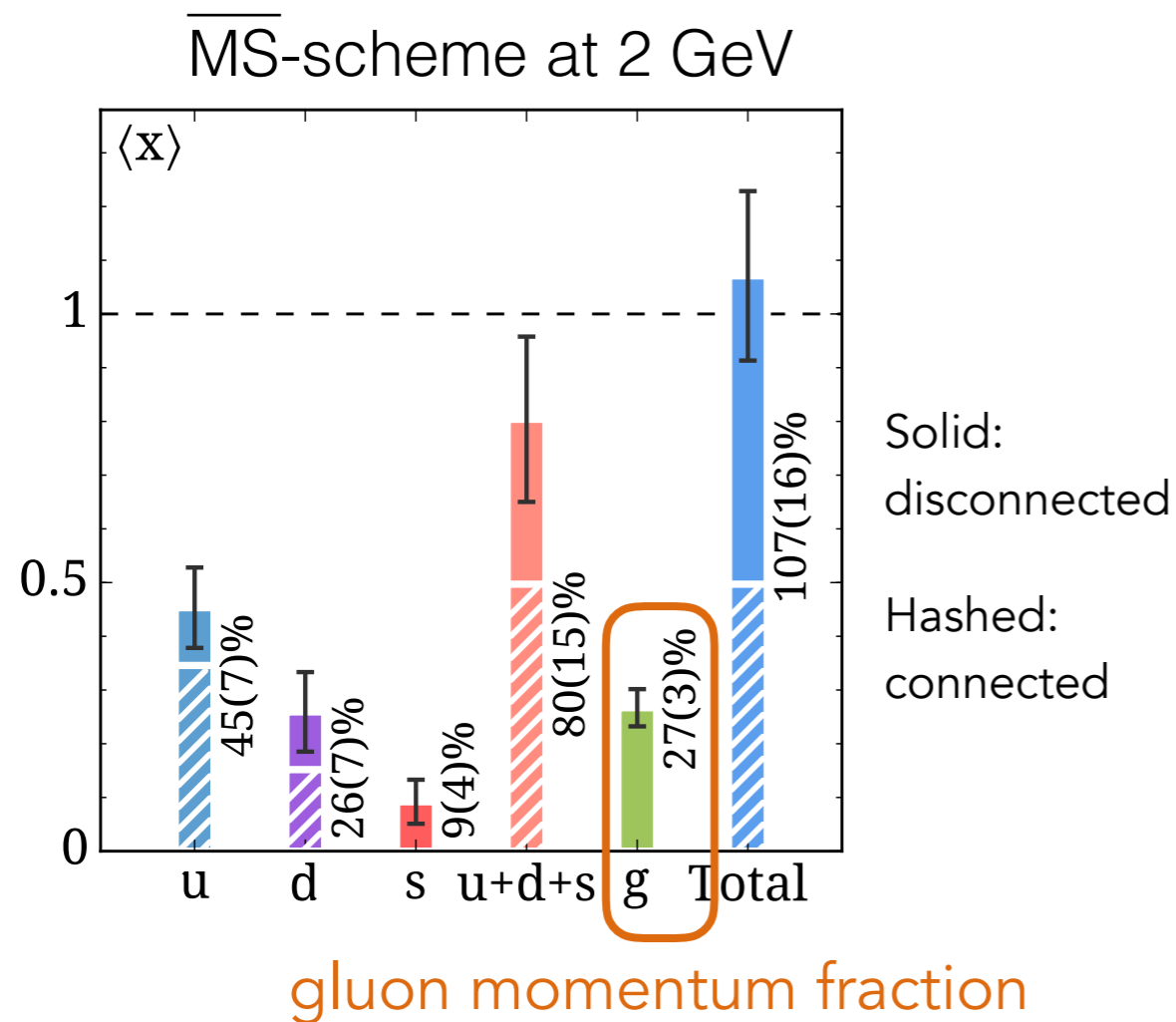
<https://web.pa.msu.edu/conf/Lattice2018/>

Nucleon gluon momentum fraction

Two direct calculations at the physical point since last year

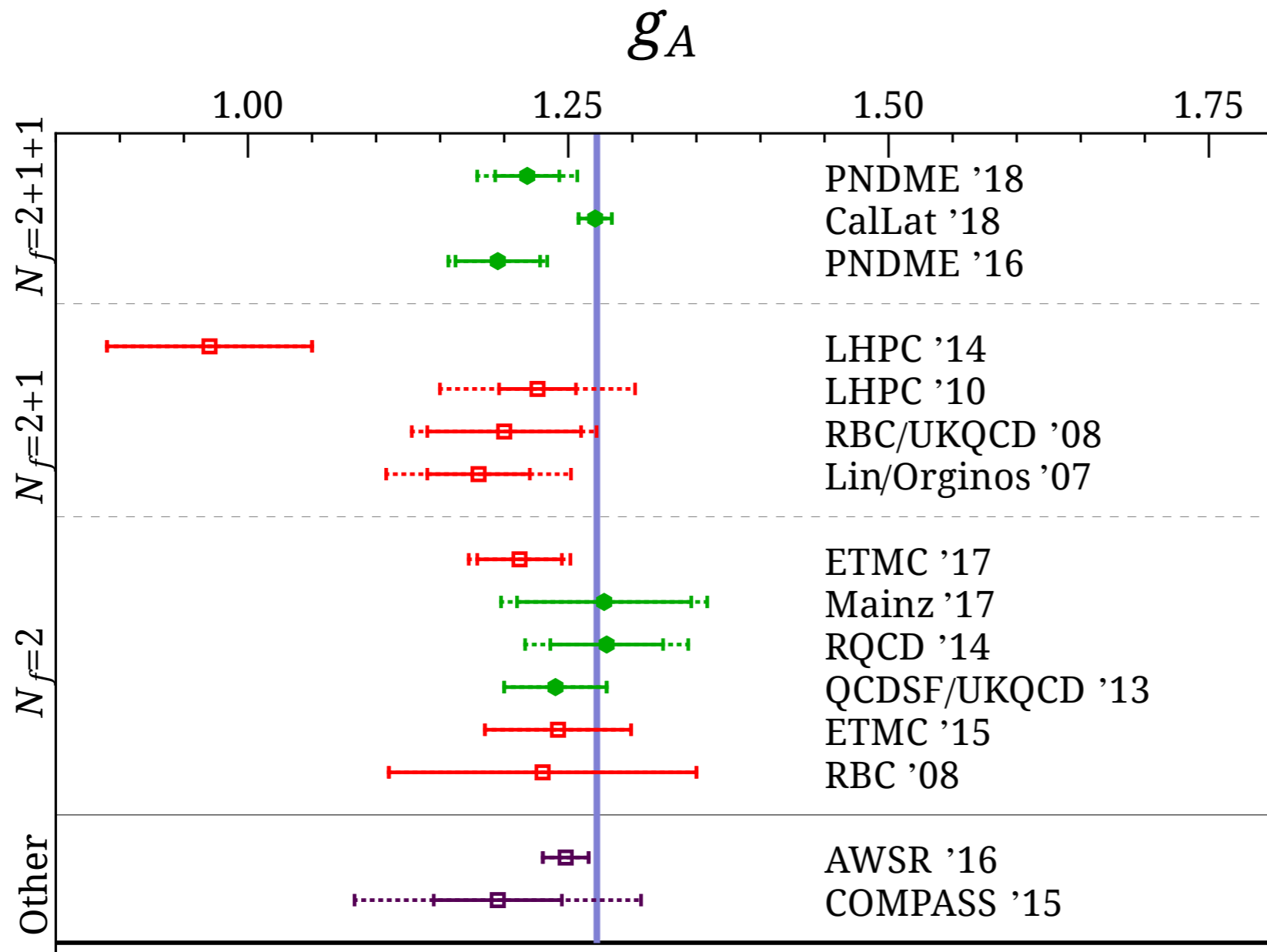
[C. Alexandrou et al., arXiv:1706.02973]

[Y-B. Yang et al., (χ QCD) arXiv:1805.00531]



Nucleon axial charge

Complete calculations with controlled uncertainties from multiple collaborations in 2018



Blue vertical line:

$$g_A = 1.2671$$

determined with high precision from nuclear beta decay

[Gupta et al, arXiv:1806.09006]

Nucleon axial form factor

Traditionally assumed to have dipole form

$$G_A(Q^2) = \frac{g_A}{(1 + Q^2 / M_A^2)}$$

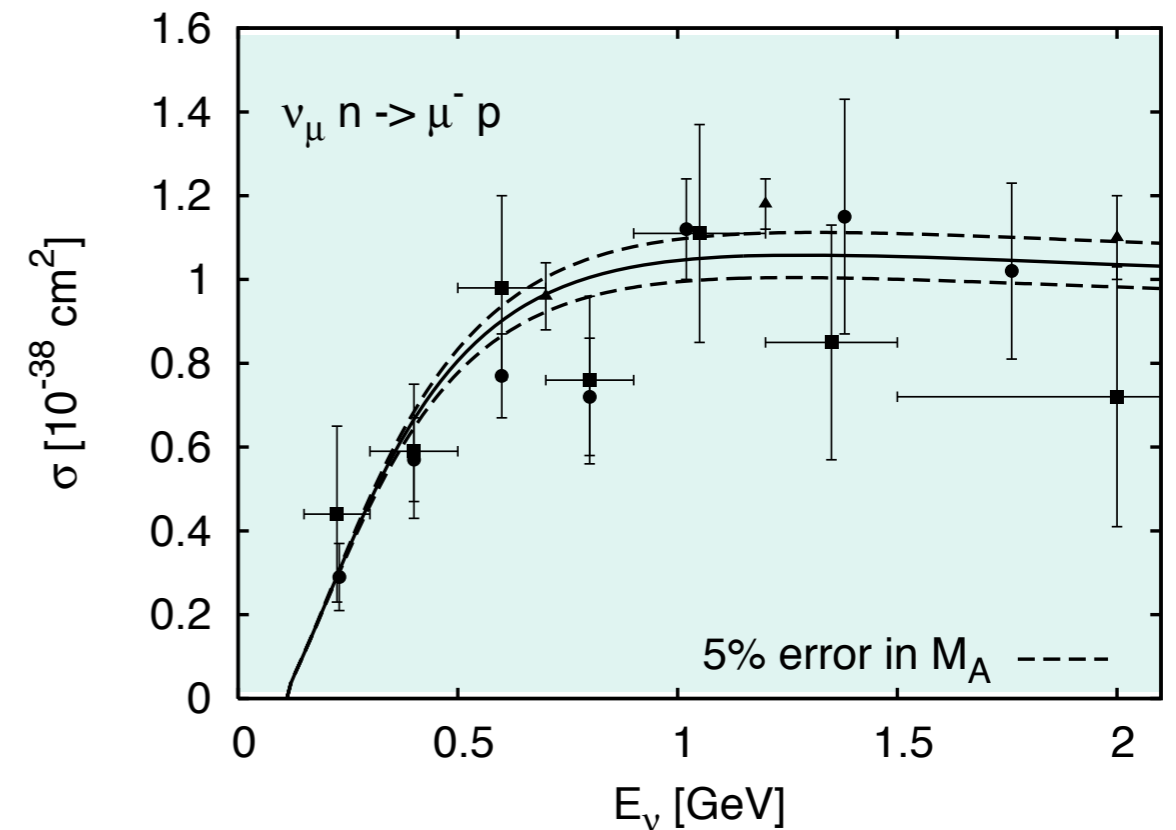
Axial charge determined with high precision from nuclear beta decay

Axial mass must be fit to data

Electromagnetic form factors show significant deviation from dipole parametrisation

➔ Model-indep. z-expansion commonly used to fit experimental or lattice QCD data [Hill & Paz (2010), Bhattacharya (2011)]

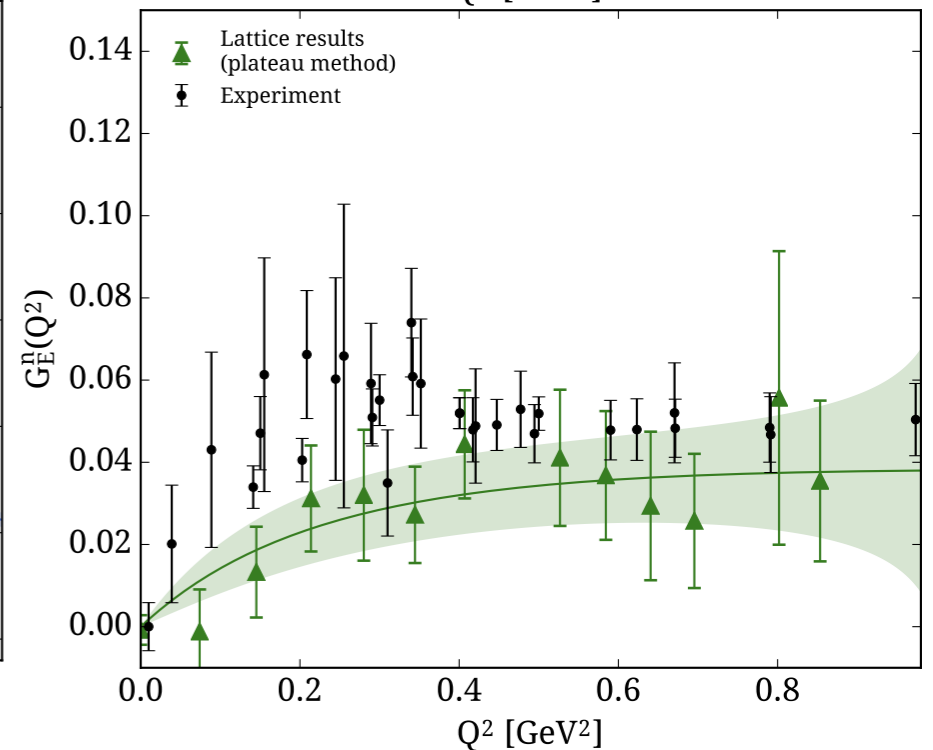
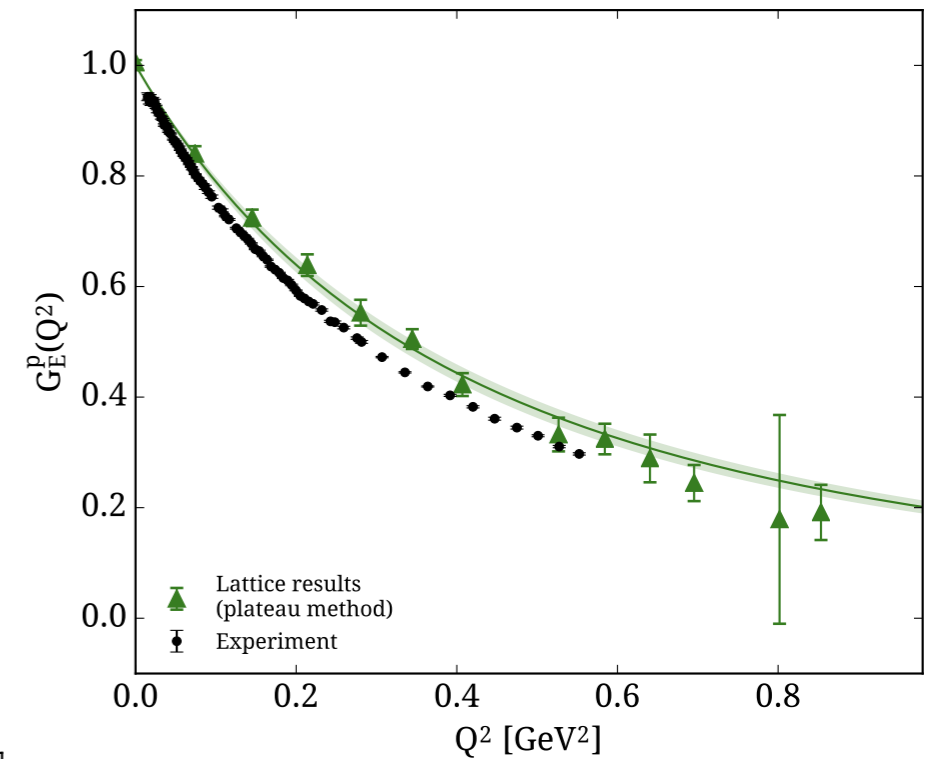
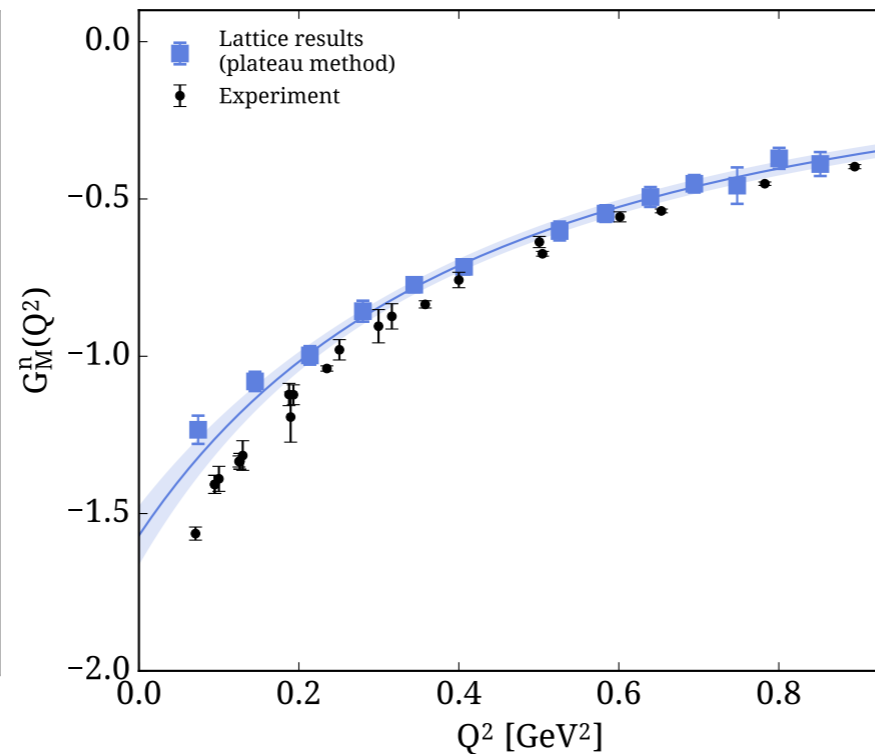
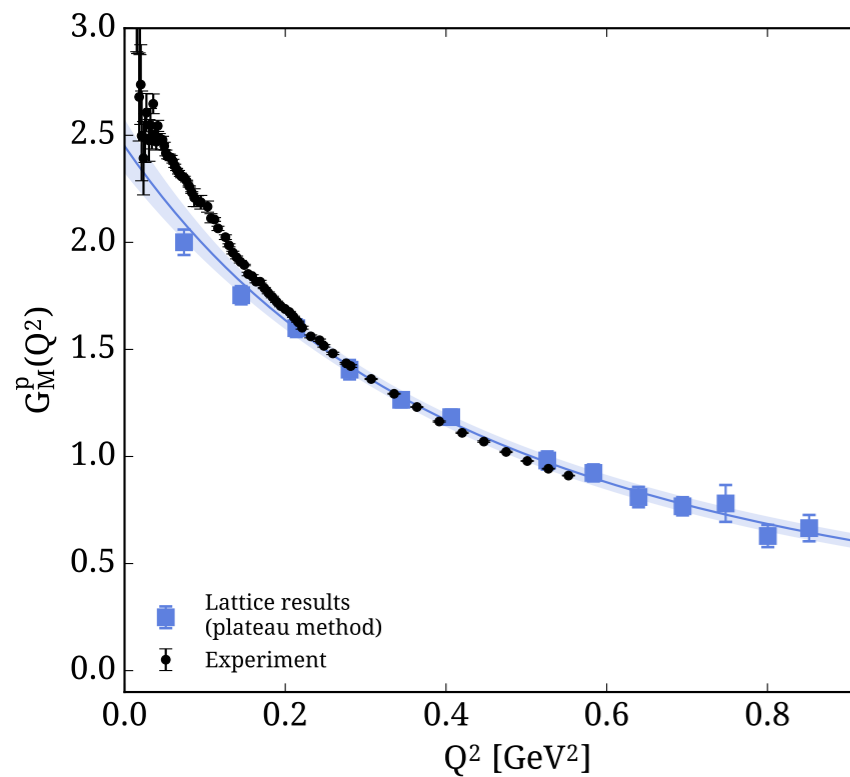
Total QE cross-section sensitive to the axial mass:



[Mosel, Ann. Rev. Nucl. Part. Sci. 66, 171 (2016)]

Nucleon electromagnetic form factors

State-of-the-art calculations have physical quark masses, large lattice volumes, and fine lattice spacings, but systematic uncertainties not yet controlled at a level comparable to flavour physics quantities



[Alexandrou et al., arXiv:1706.00469]

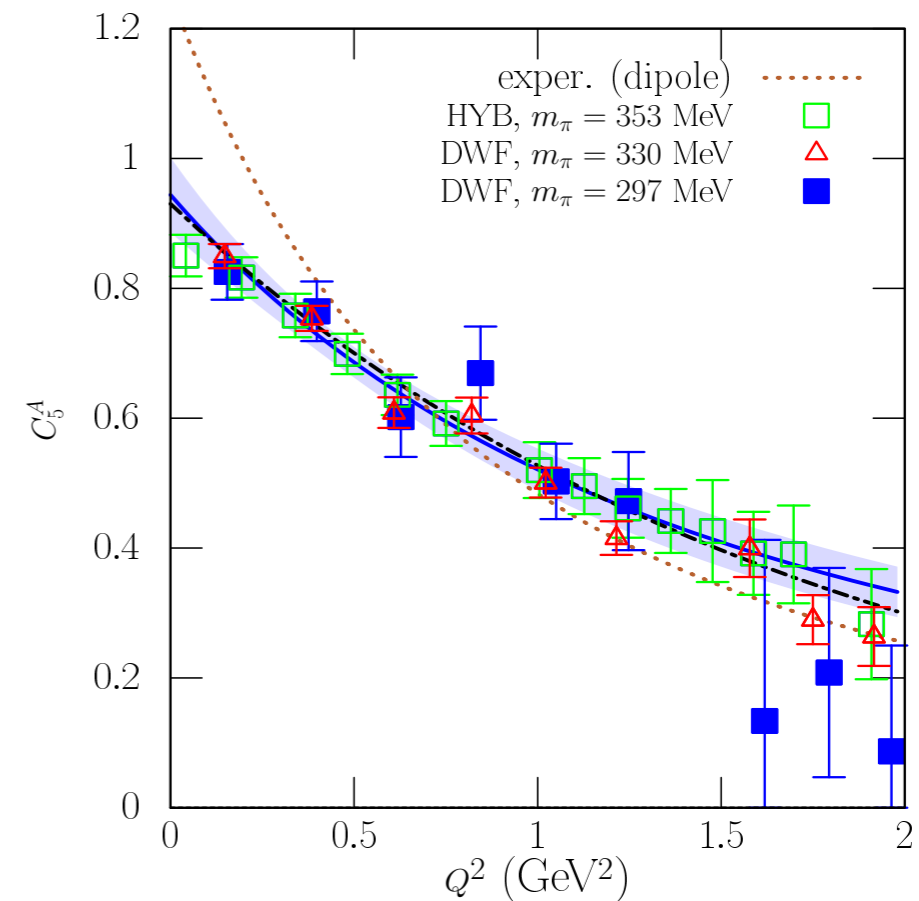
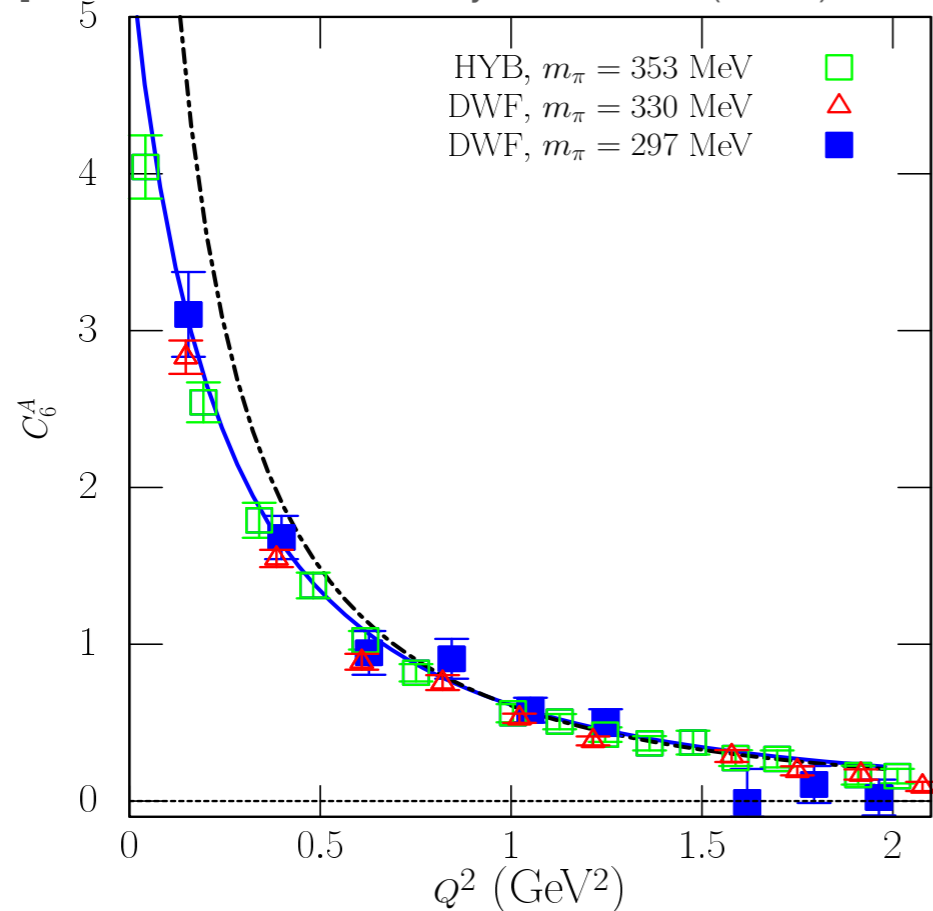
Phiala Shanahan, MIT

Transition form factors

Lattice QCD calculation of axial N Δ transition form factor:

$$\langle \Delta(p', s') | A_\mu^3 | N(p, s) \rangle = i \sqrt{\frac{2}{3}} \left(\frac{m_\Delta m_N}{E_\Delta(\mathbf{p}') E_N(\mathbf{p})} \right)^{1/2} \bar{u}_{\Delta+}^\lambda(p', s') \left[\left(\frac{C_3^A(q^2)}{m_N} \gamma^\nu + \frac{C_4^A(q^2)}{m_N^2} p'^\nu \right) (g_{\lambda\mu} g_{\rho\nu} - g_{\lambda\rho} g_{\mu\nu}) q^\rho + C_5^A(q^2) g_{\lambda\mu} + \frac{C_6^A(q^2)}{m_N^2} q_\lambda q_\mu \right] u_P(p, s)$$

[C Alexandrou et al., Phys.Rev. D83 (2011) 014501]



CAVEAT: Complexities at physical point with unstable resonances,
but formalism exists: [Lellouch-Lüscher hep-lat/0003023]

Muon $g-2$

Muon $g-2$: new lattice QCD results in last few years

Hadronic light-by-light

PRD93(2015)014503 (RBC)

PRL118(2016)022005 (RBC)

PRD96(2017)034515 (RBC)

Hadronic vacuum polarisation

PRL116(2015)232002 (RBC/UKQCD)

PRD93(2016)054508 (Aubin et al.)

JHEP04(2016)063 (RBC/UKQCD)

PRD96(2017)034516 (FNAL/HPQCD)

JHEP09(2017)153 (RBC/UKQCD)

PRL120(2017)152001 (FNAL/HPQCD/MILC)

arXiv:1801.07224 (RBC/UKQCD)