

Structure of the Nucleon in the Valence Quark Region

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- Structure of the nucleon; Inclusive spin responses
 - ➔ Valence region and quark orbital angular momentum
 - ➔ Moments of spin structure functions
 - ➔ Burkhardt-Cottingham sum rule
 - ➔ Color polarizabilities
- Structure of the nucleon; Semi-inclusive spin response
 - ➔ Transversity, Collins, Sivers, etc...



Introduction

Understand the nucleon structure in the valence quark region

→ Complete knowledge of parton distribution functions (PDFs).

→ Unpolarized, helicity dependent and transversity distribution functions...

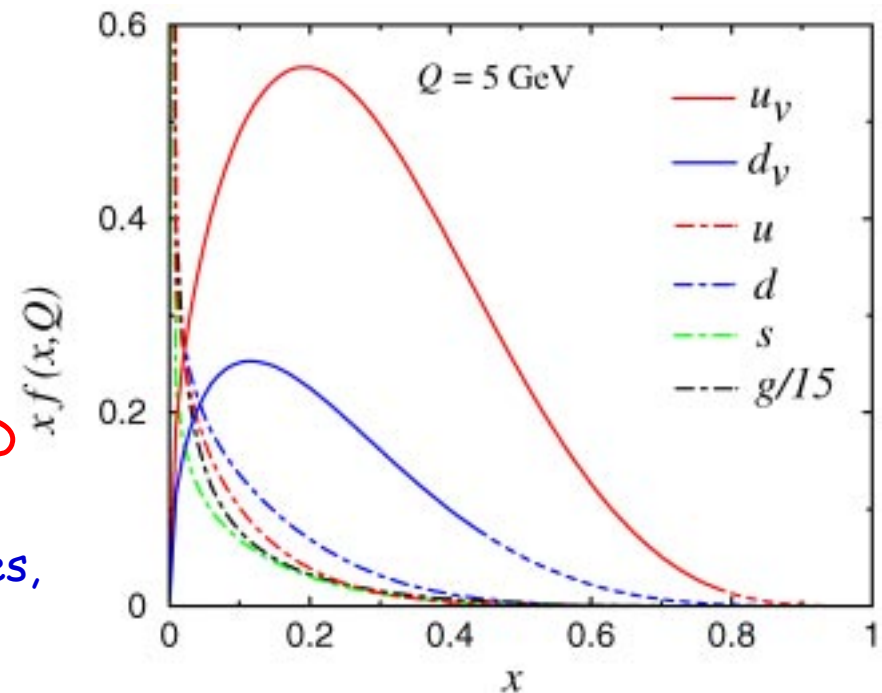
→ Why large x ?

→ large x exposes valence quarks
- free of sea effects

→ $x \rightarrow 1$ behavior - sensitive test of
spin-flavor symmetry breaking

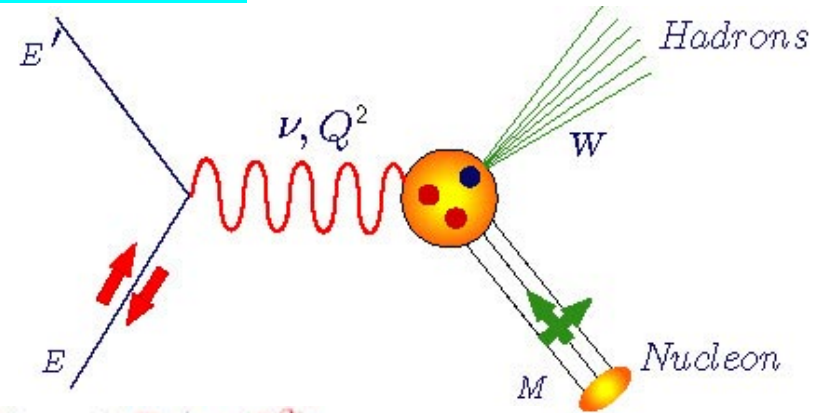
→ important for higher moments of
PDFs - compare with lattice QCD

→ intimately related with resonances,
quark-hadron duality



Inclusive DIS

- Unpolarized structure functions $F_1(x, Q^2)$ and $F_2(x, Q^2)$
 - ➔ Proton & neutron measurements provide d/u distributions ratio



$$U \quad \frac{d^2\sigma}{dE'd\Omega}(\downarrow\uparrow + \uparrow\uparrow) = \frac{8\alpha^2 \cos^2(\theta/2)}{Q^4} \left[\frac{F_2(x, Q^2)}{\nu} + \frac{2F_1(x, Q^2)}{M} \tan^2(\theta/2) \right]$$

- Polarized structure functions $g_1(x, Q^2)$ and $g_2(x, Q^2)$
 - ➔ Proton & neutron measurements combined with d/u provide the spin-flavor distributions $\Delta u/u$ & $\Delta d/d$

Q^2 : Four-momentum transfer
 x : Bjorken variable
 ν : Energy transfer
 M : Nucleon mass
 W : Final state hadrons mass

$$L \quad \frac{d^2\sigma}{dE'd\Omega}(\downarrow\uparrow - \uparrow\uparrow) = \frac{4\alpha^2 E'}{MQ^2 \nu E} \left[(E + E' \cos \theta) g_1(x, Q^2) - \frac{Q^2}{\nu} g_2(x, Q^2) \right]$$

$$T \quad \frac{d^2\sigma}{dE'd\Omega}(\downarrow\Rightarrow - \uparrow\Rightarrow) = \frac{4\alpha^2 \sin \theta E^2}{MQ^2 \nu^2 E} \left[\nu g_1(x, Q^2) + 2E g_2(x, Q^2) \right]$$

Virtual photon-nucleon asymmetries

$$A_1 = \frac{g_1(x, Q^2) - \gamma^2 g_2(x, Q^2)}{F_1(x, Q^2)}$$

$$A_2 = \frac{\gamma[g_1(x, Q^2) + g_2(x, Q^2)]}{F_1(x, Q^2)}$$

where $\gamma = \sqrt{Q^2}/\nu$

- Positivity constraints

$$|A_1| \leq 1 \quad \text{and} \quad |A_2| \leq \sqrt{R(1 + A_1)/2}$$

In the quark-parton model:

$$F_1(x, Q^2) = \frac{1}{2} \sum_f e_f^2 q_f(x, Q^2) \quad g_1(x, Q^2) = \frac{1}{2} \sum_f e_f^2 \Delta q_f(x, Q^2)$$

$$q_f(x) = q_f^\uparrow(x) + q_f^\downarrow(x) \quad \Delta q_f(x) = q_f^\uparrow(x) - q_f^\downarrow(x)$$

$q_f(x)$ quark momentum distributions of flavor f

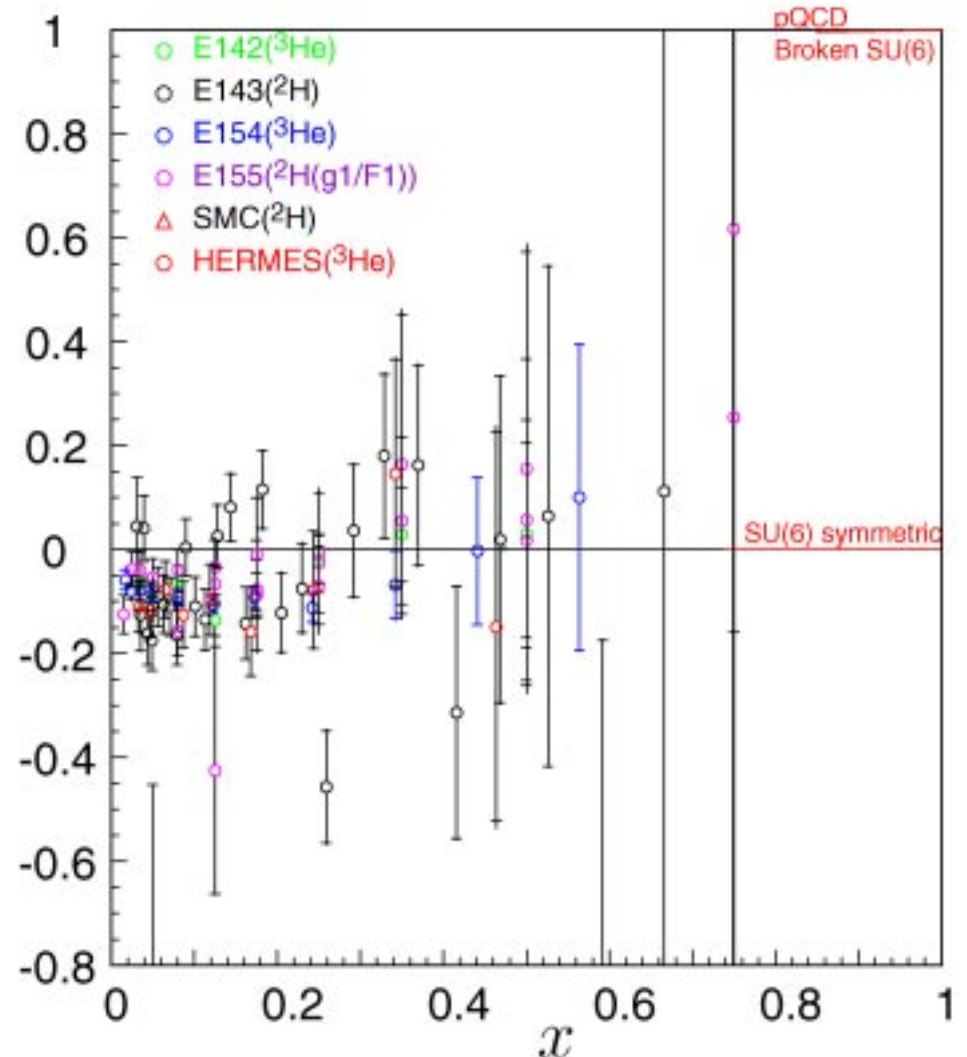
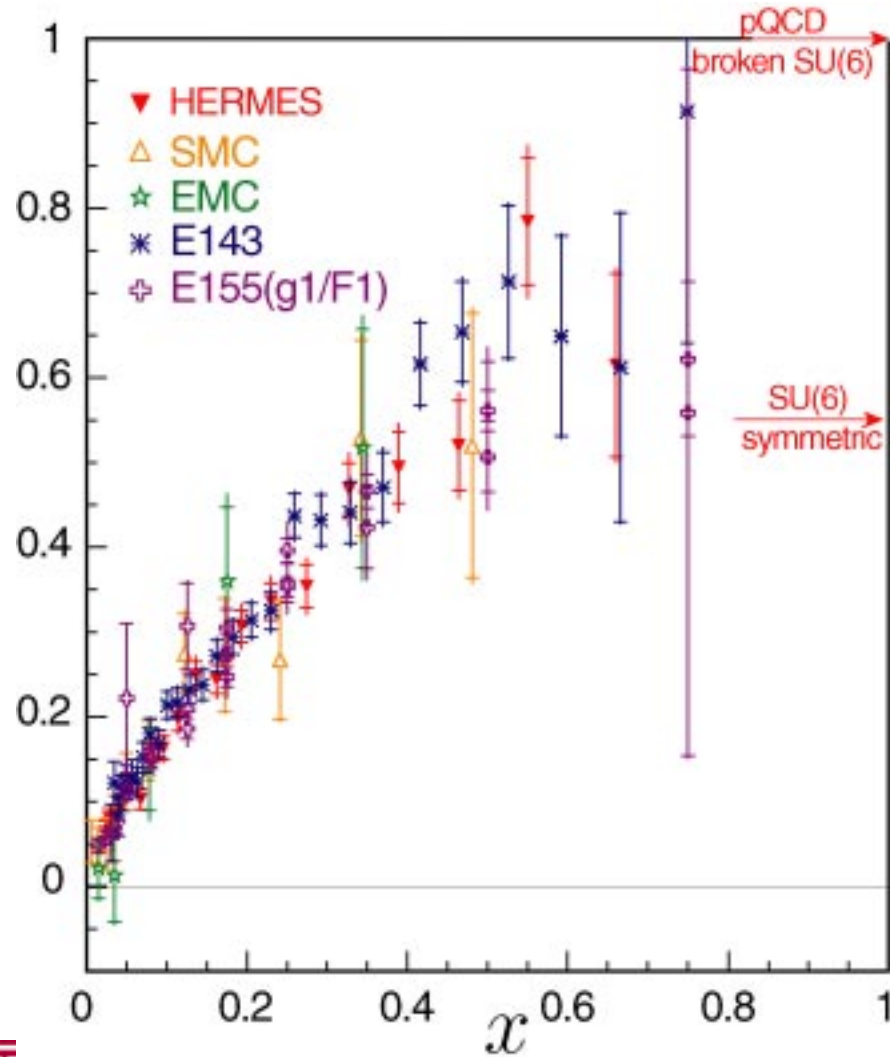
$\uparrow(\downarrow)$ parallel (antiparallel) to the nucleon spin



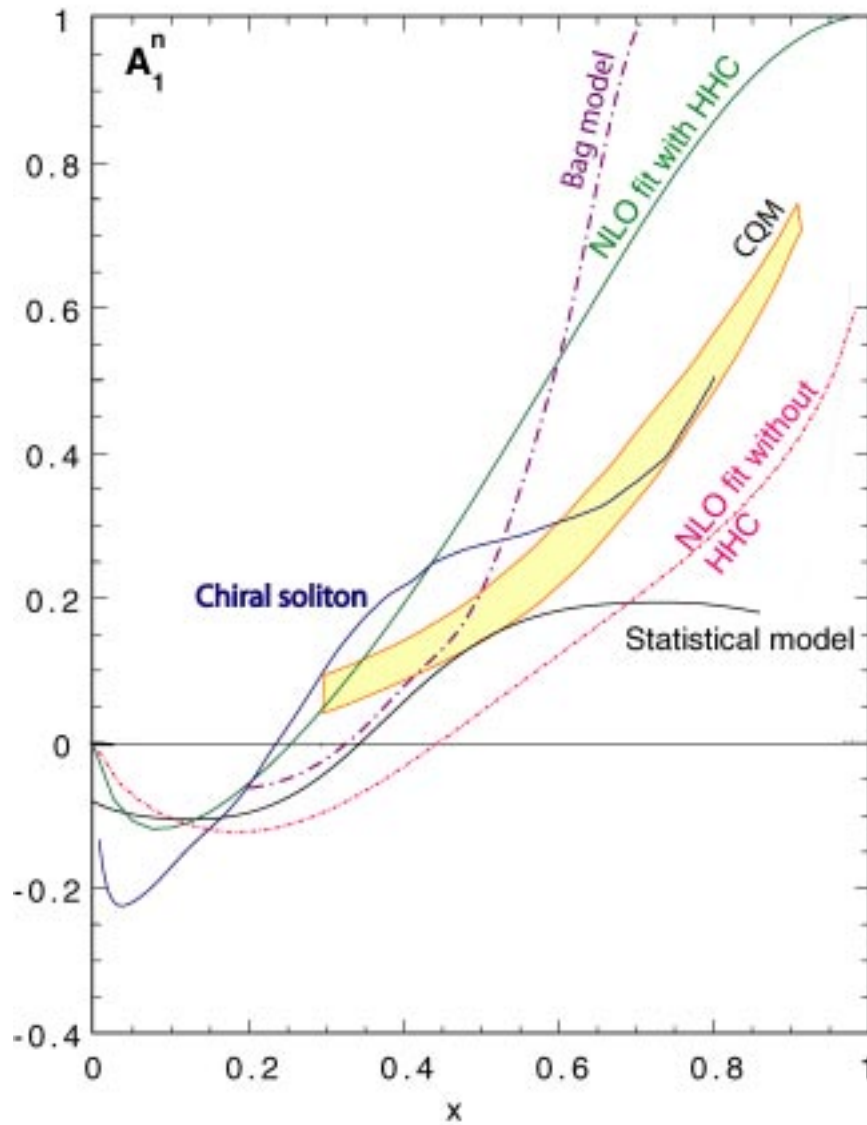
World data for A_1

Proton

Neutron



Polarized quarks as $x \rightarrow 1$



- SU(6) symmetry:

- $A_1^p = 5/9$ $A_1^n = 0$ $d/u = 1/2$

- $\Delta u/u = 2/3$ $\Delta d/d = -1/3$

- Broken SU(6) via scalar diquark dominance

- $A_1^p \rightarrow 1$ $A_1^n \rightarrow 1$ $d/u \rightarrow 0$

- $\Delta u/u \rightarrow 1$ $\Delta d/d \rightarrow -1/3$

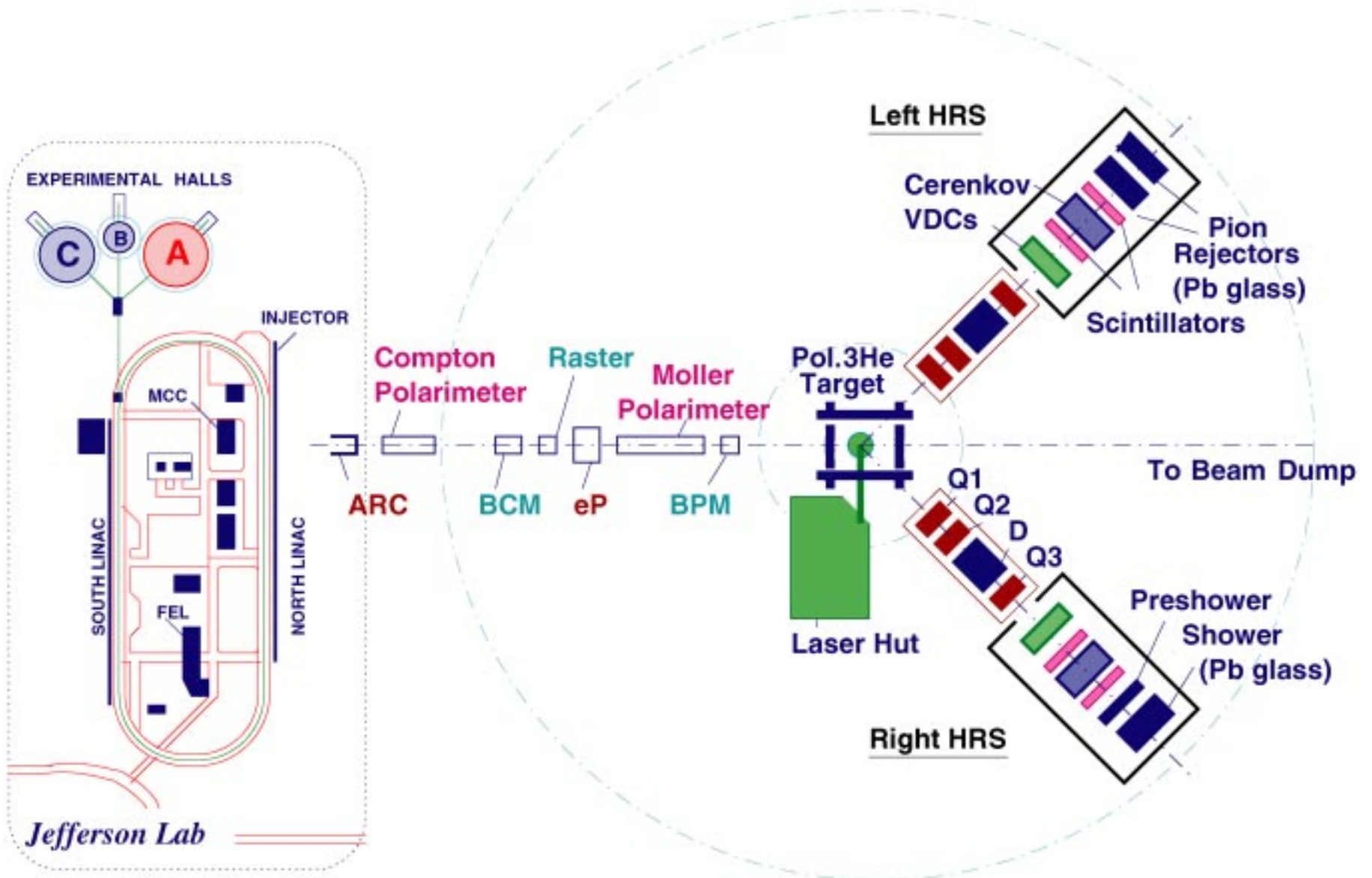
- Broken SU(6) via helicity conservation

- $A_1^p \rightarrow 1$ $A_1^n \rightarrow 1$ $d/u \rightarrow 1/5$

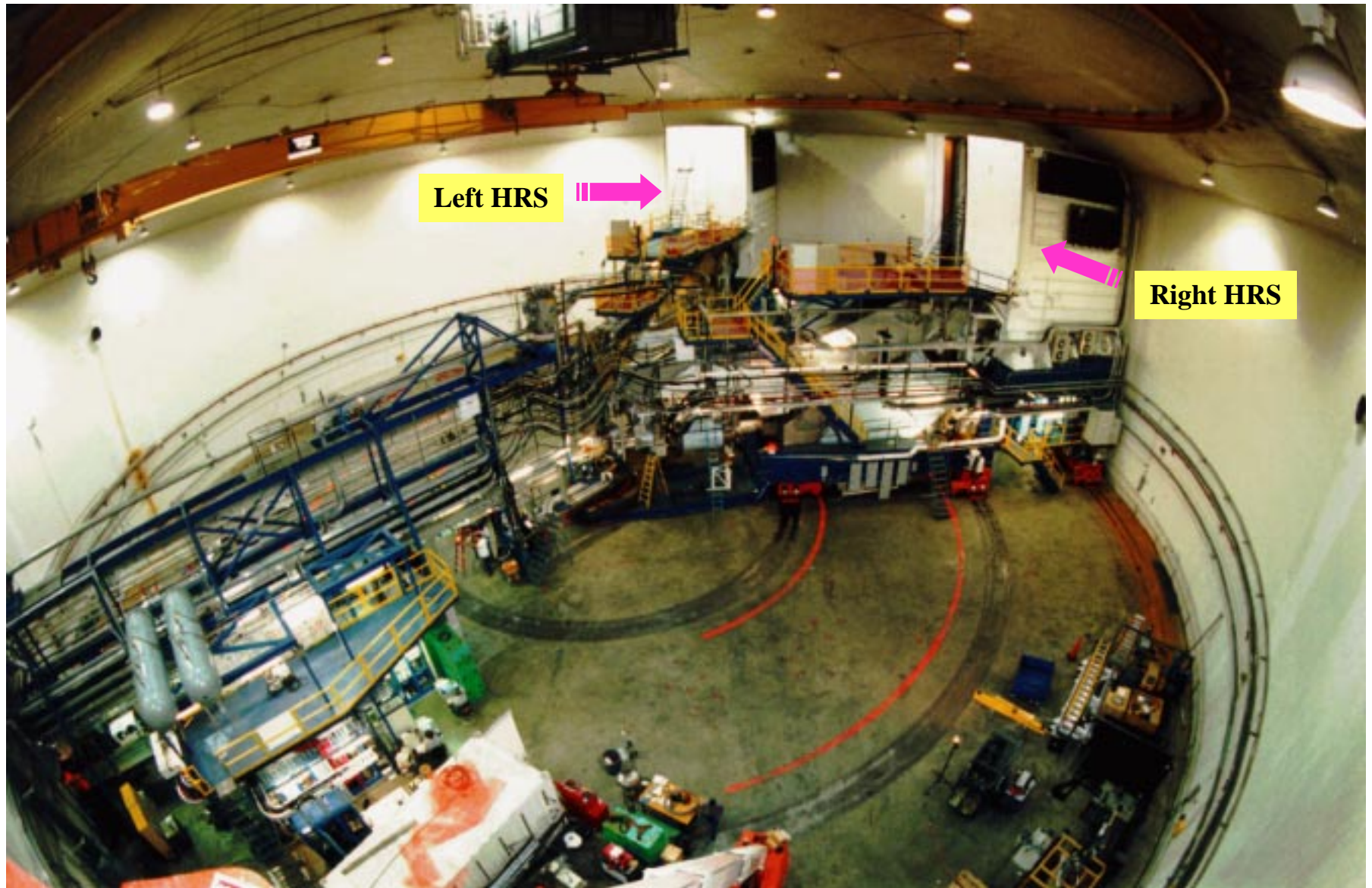
- $\Delta u/u \rightarrow 1$ $\Delta d/d \rightarrow 1$

Note that $\Delta q/q$ as $x \rightarrow 1$ is more sensitive to **spin-flavor** symmetry breaking effects than A_1

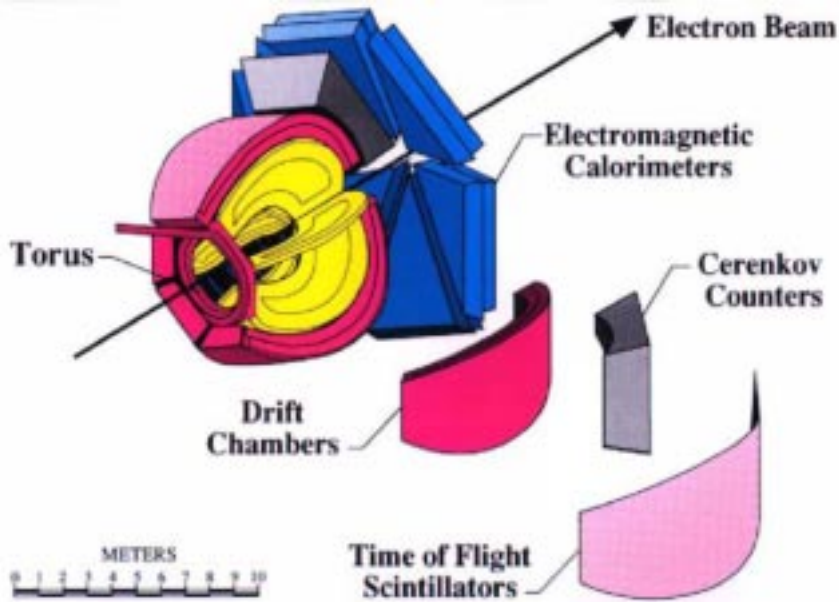
Hall A Experimental Setup



Hall A at Jefferson Lab

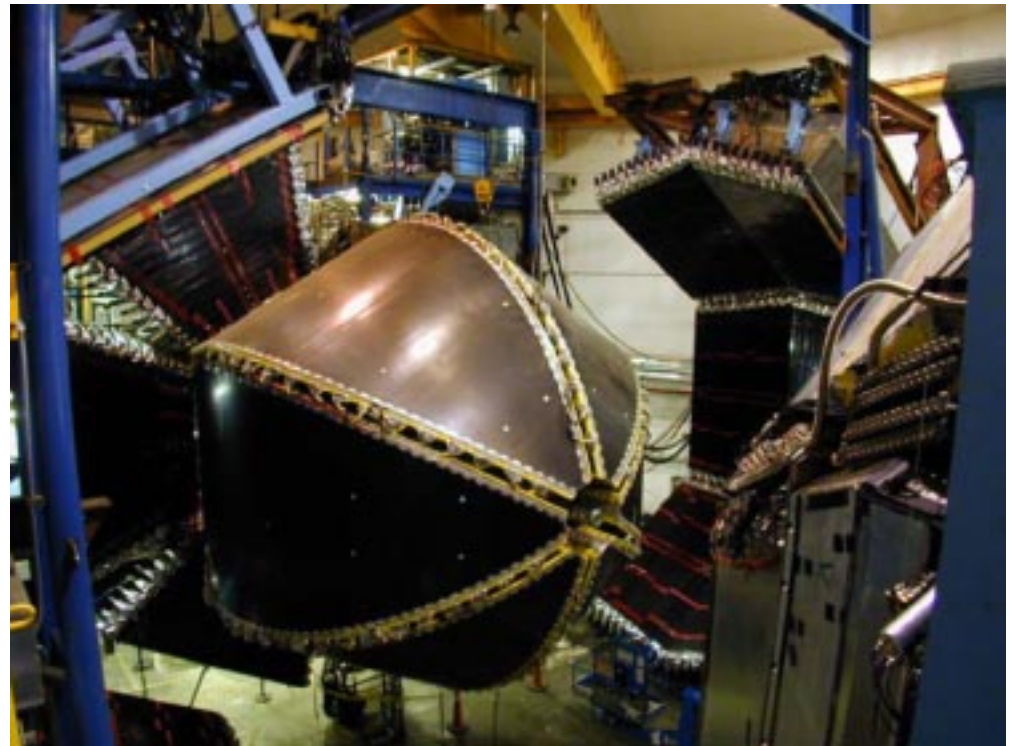


Hall B Experimental Setup

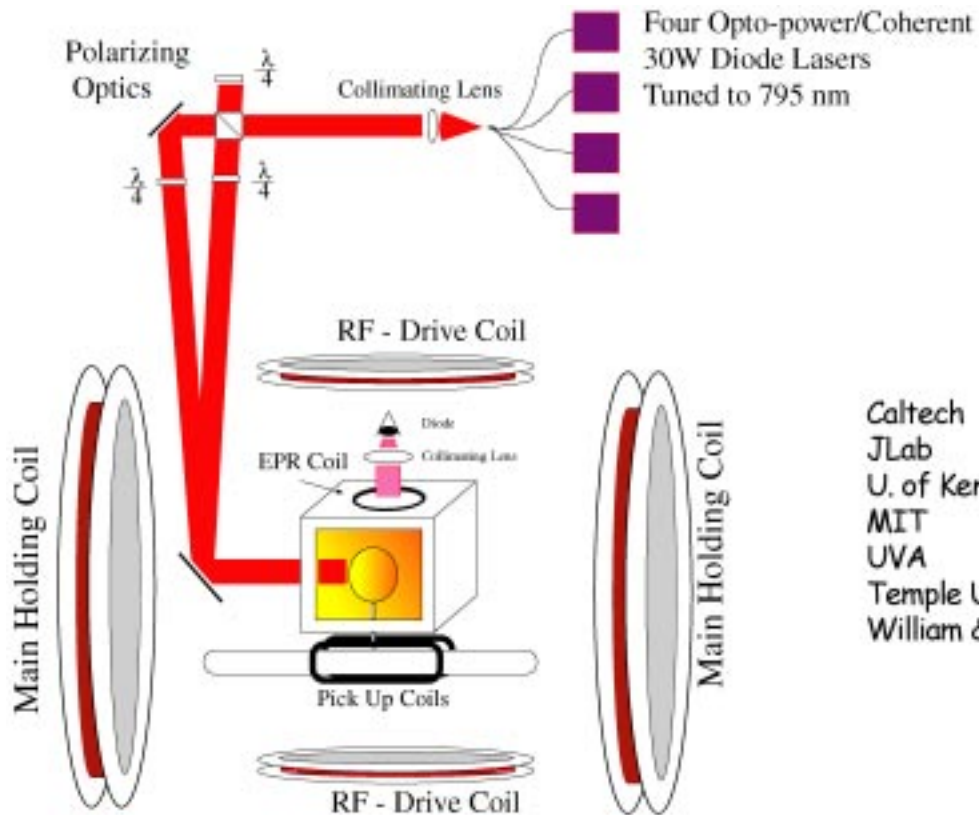


CEBAF
Large
Acceptance
Spectrometer

- Large kinematical coverage
- detection of charged and neutral particles
- Multiparticle final state



JLab Polarized ^3He target

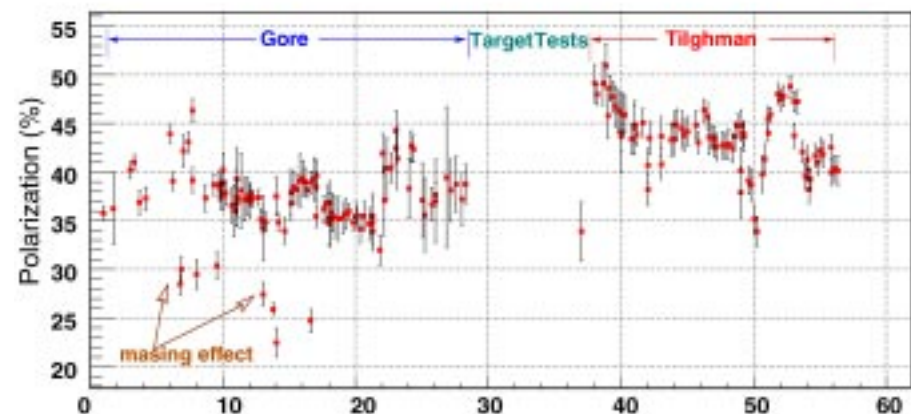


- NMR and EPR techniques for polarization monitoring.
- Elastic scattering for current induced depolarization.
- Target length 40 cm, window thickness 100 μm

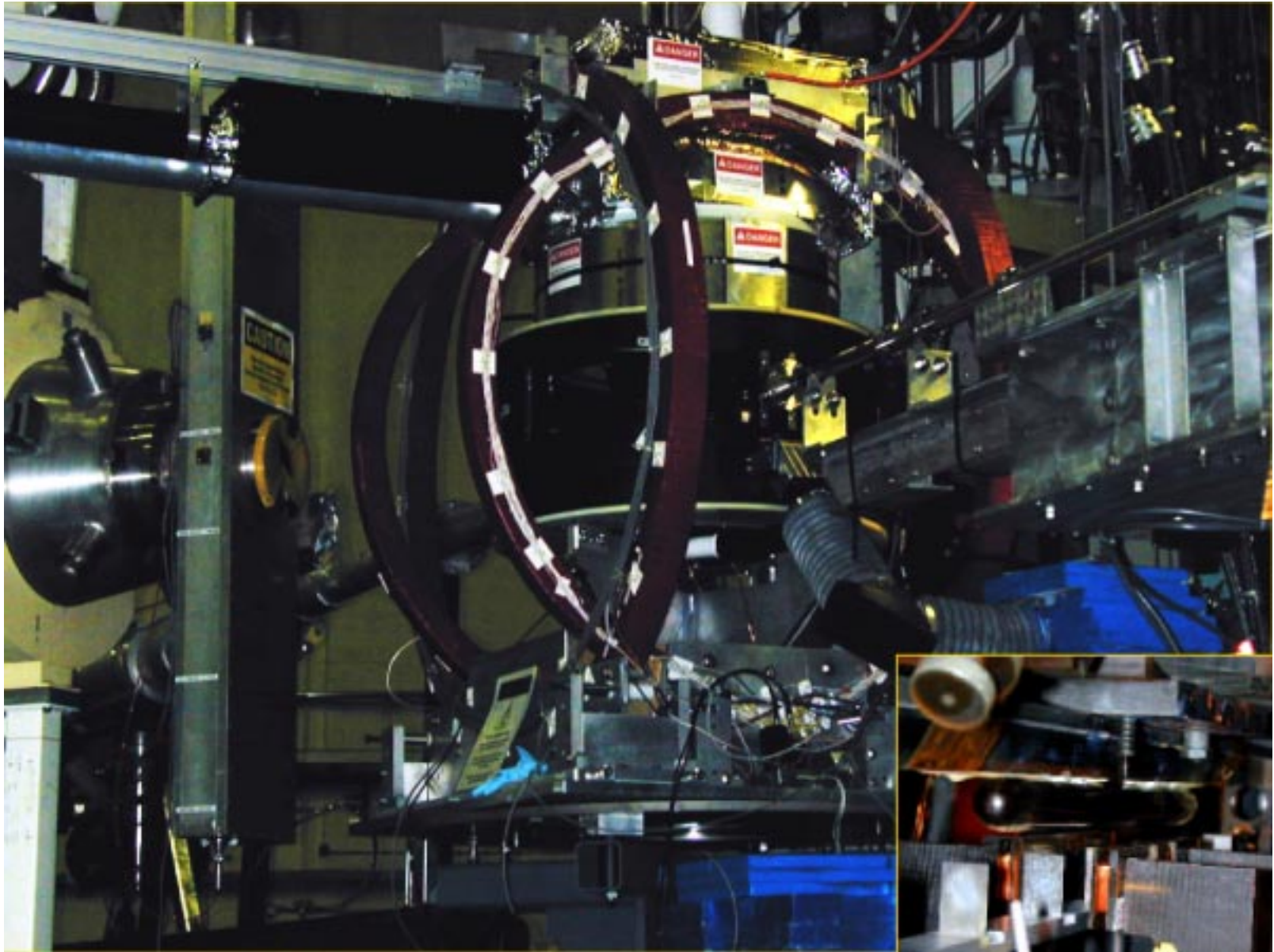
Caltech
JLab
U. of Kentucky
MIT
UVA
Temple U.
William & Mc

Performance of the target

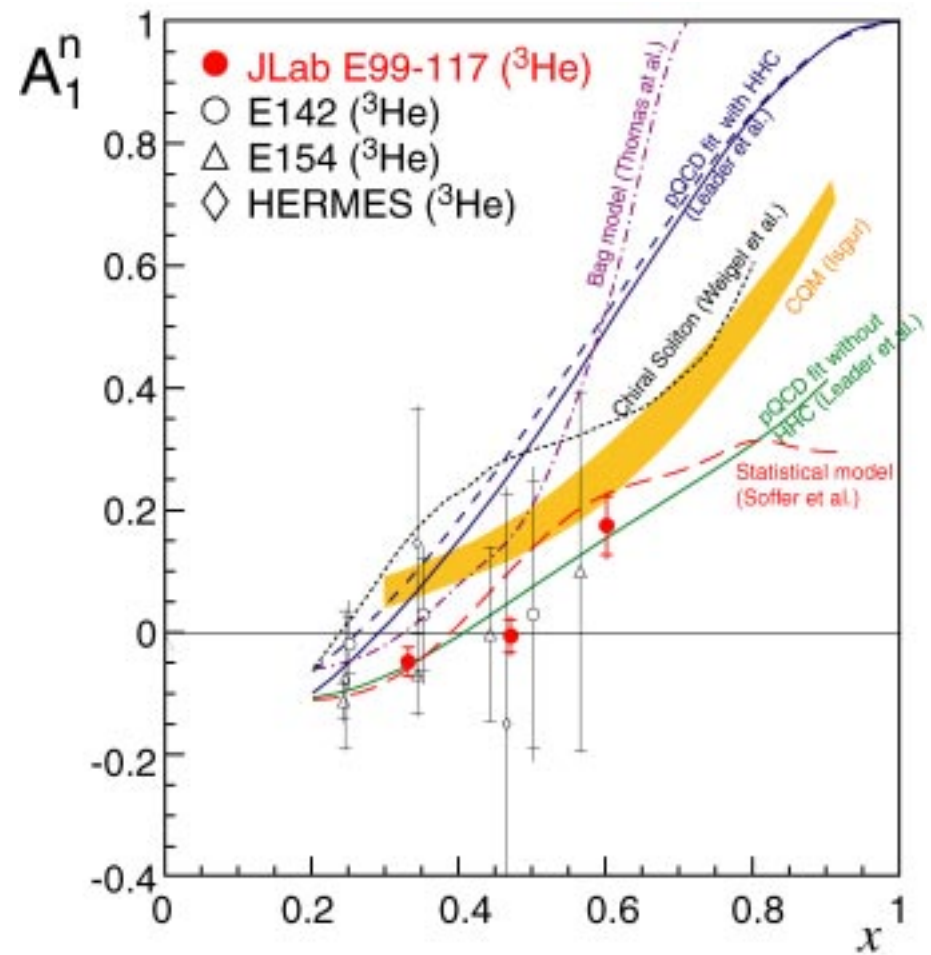
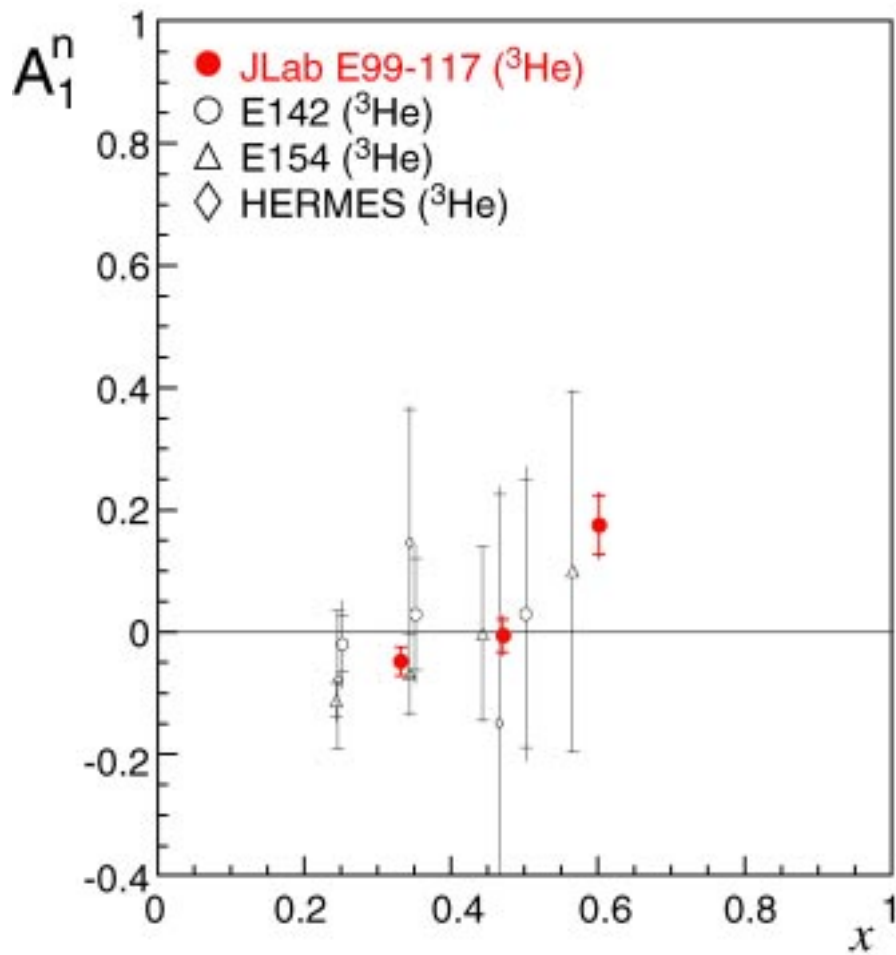
Cell Name	Field direction	0°	180°	270°
<i>Gore</i>	June 06 - July 03	37%	35%	43%
<i>Tilghman</i>	July 13 - July 31	45%	43%	39%



Polarized ^3He Target



JLab E99-117 A_1^n Results



Spokespeople: J.P. Chen,
Z.-E.M. & P. Souder

Thesis student: X. Zheng

Phys. Rev. Lett. 92, 012004 (2004)

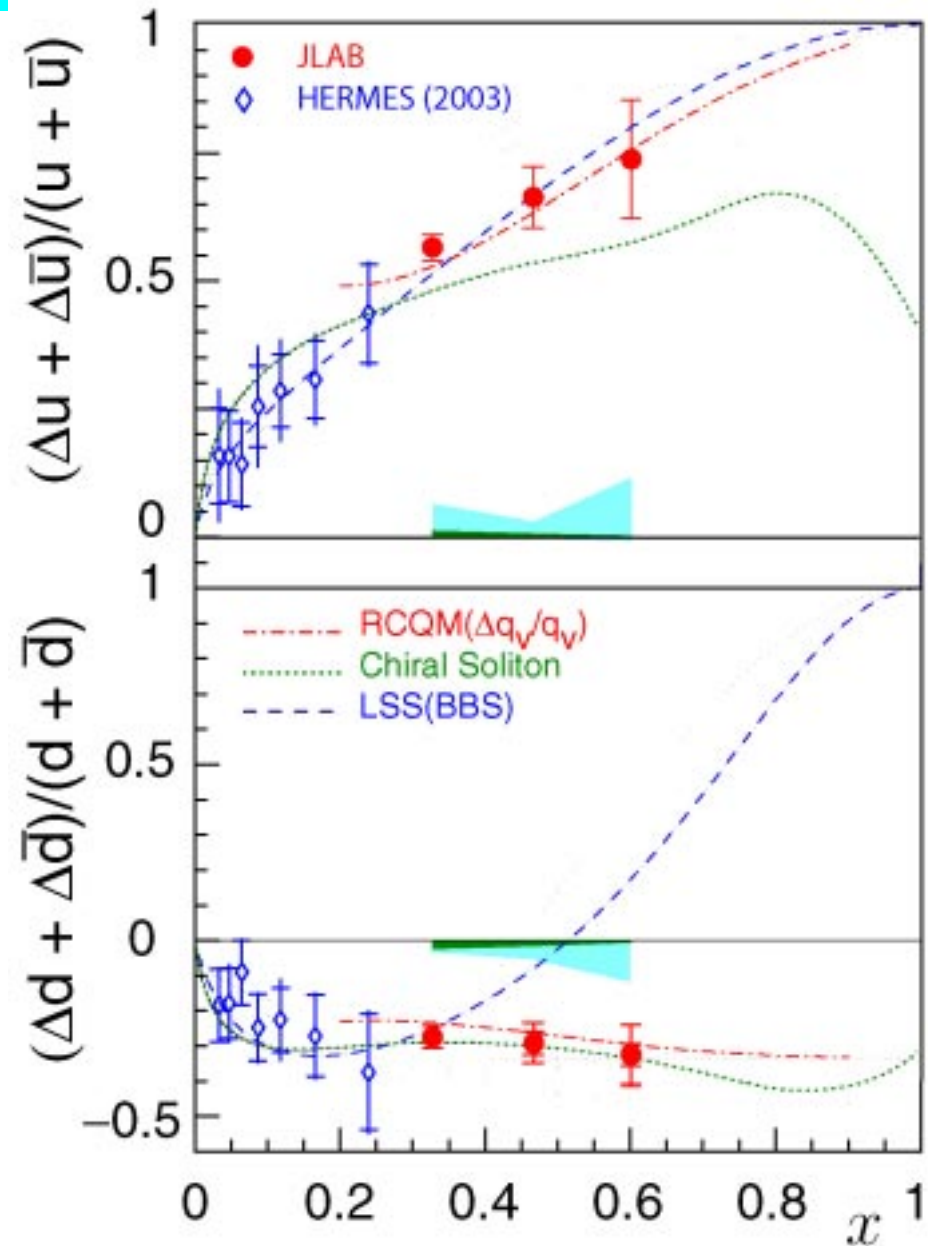
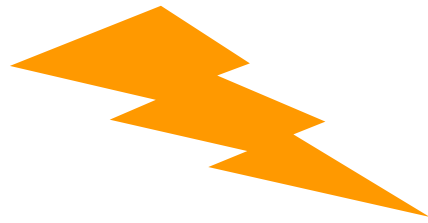


Helicity-Flavor Decomposition

$$\frac{\Delta u + \Delta \bar{u}}{u} = \frac{4}{15} \frac{g_1^p}{F_1^p} (4 + R^{du}) - \frac{1}{15} \frac{g_1^n}{F_1^n} (1 + 4R^{du})$$

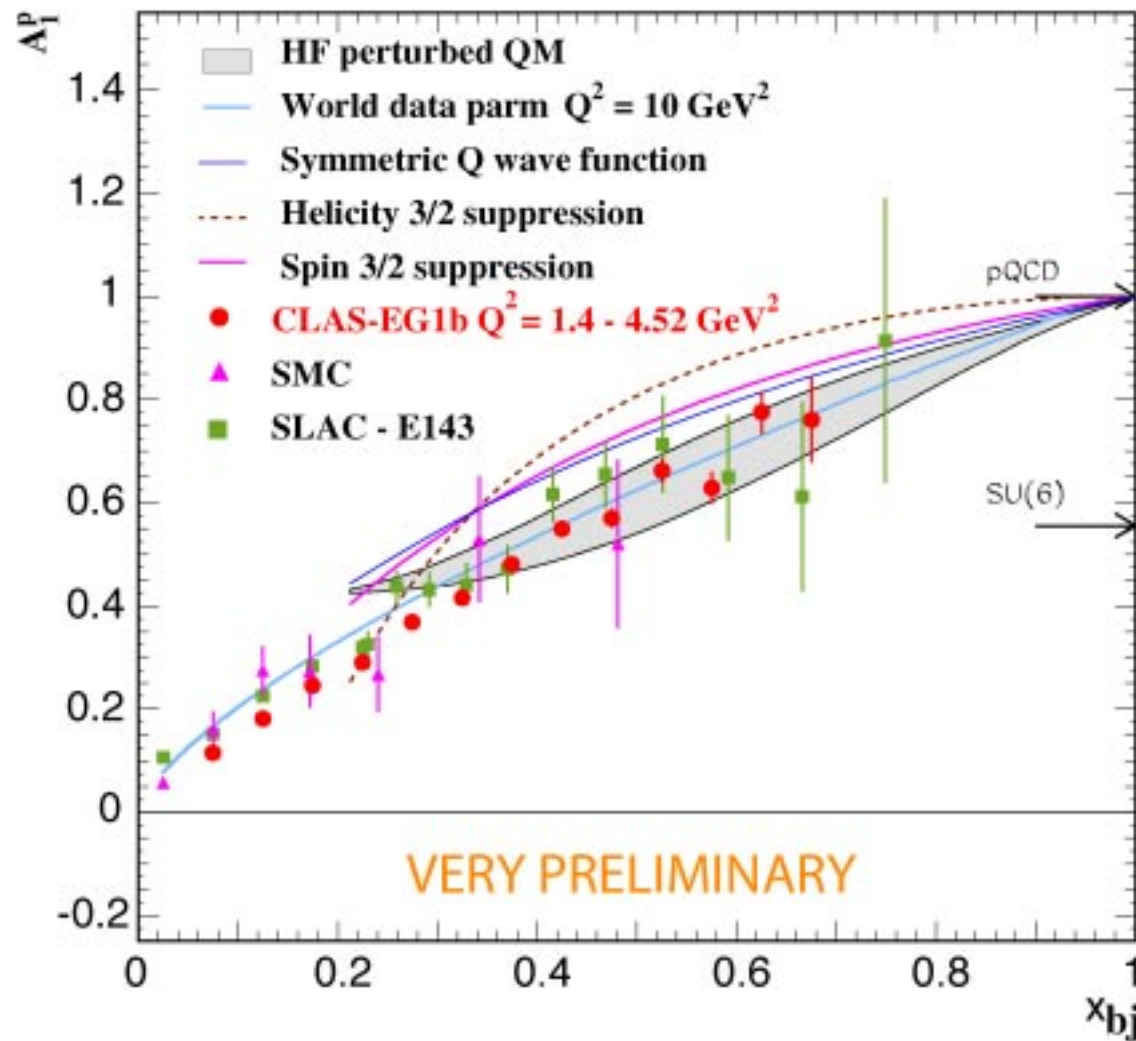
$$\frac{\Delta d + \Delta \bar{d}}{d} = \frac{4}{15} \frac{g_1^n}{F_1^n} (4 + \frac{1}{R^{du}}) - \frac{1}{15} \frac{g_1^p}{F_1^p} (1 + 4\frac{1}{R^{du}})$$

$$R^{du} = \frac{d + \bar{d}}{u + \bar{u}}$$

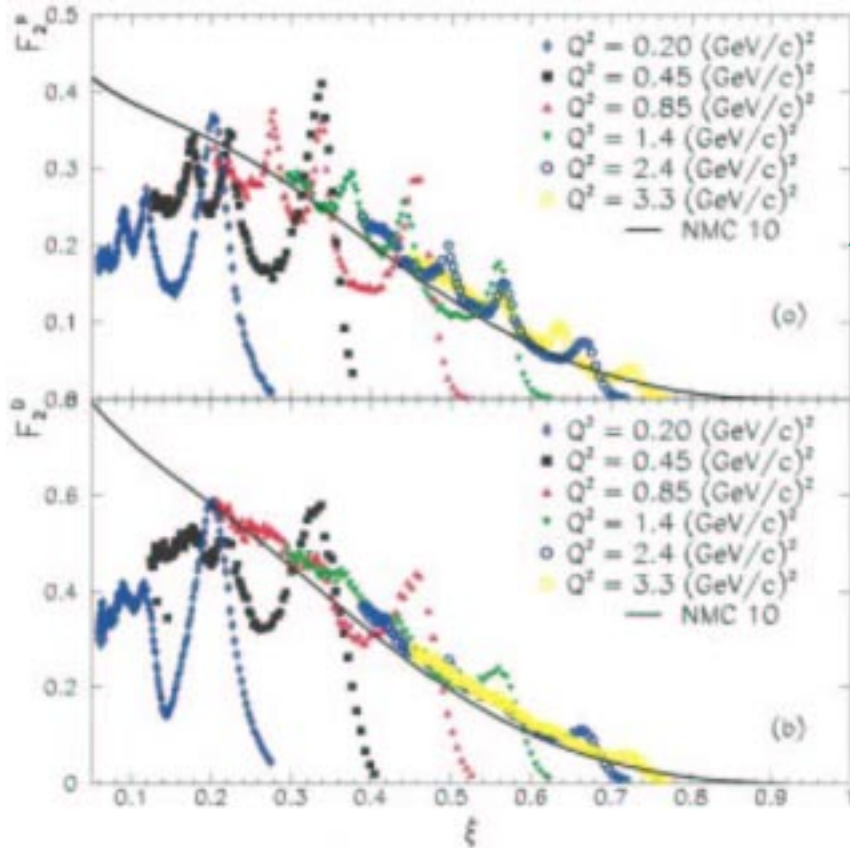


Hall B EG1b preliminary results

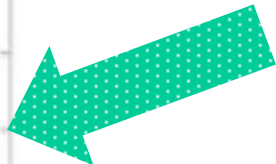
PROTON



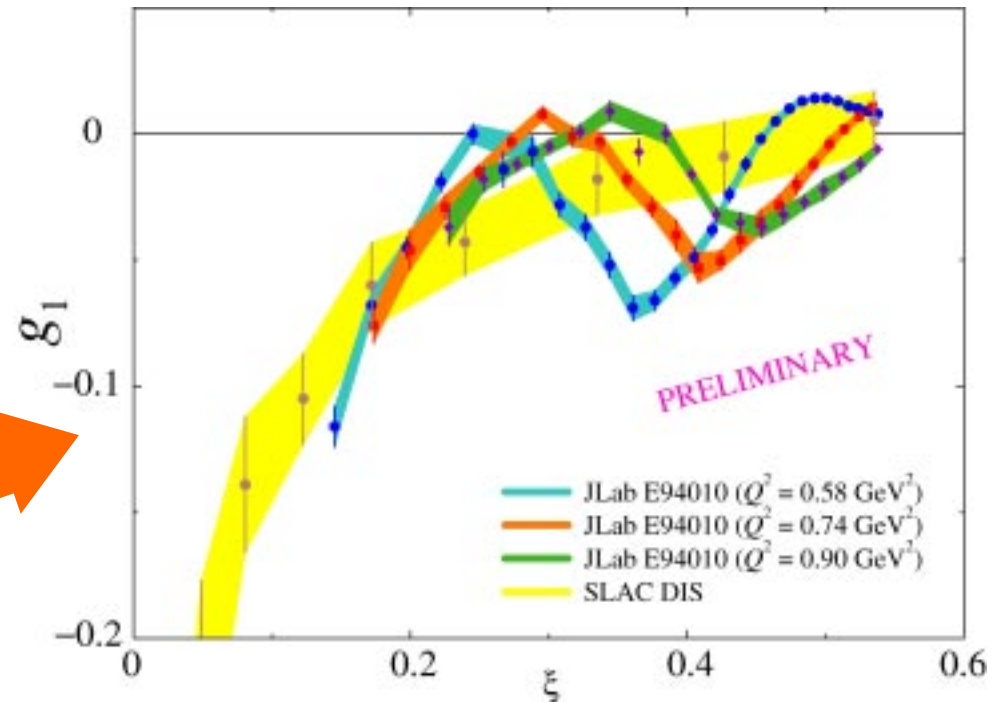
Duality in the neutron spin structure



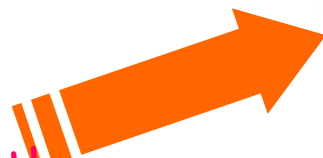
Helicity independent case
Proton and Deuteron
Niculescu et al. PRL 85 (2000) 1182



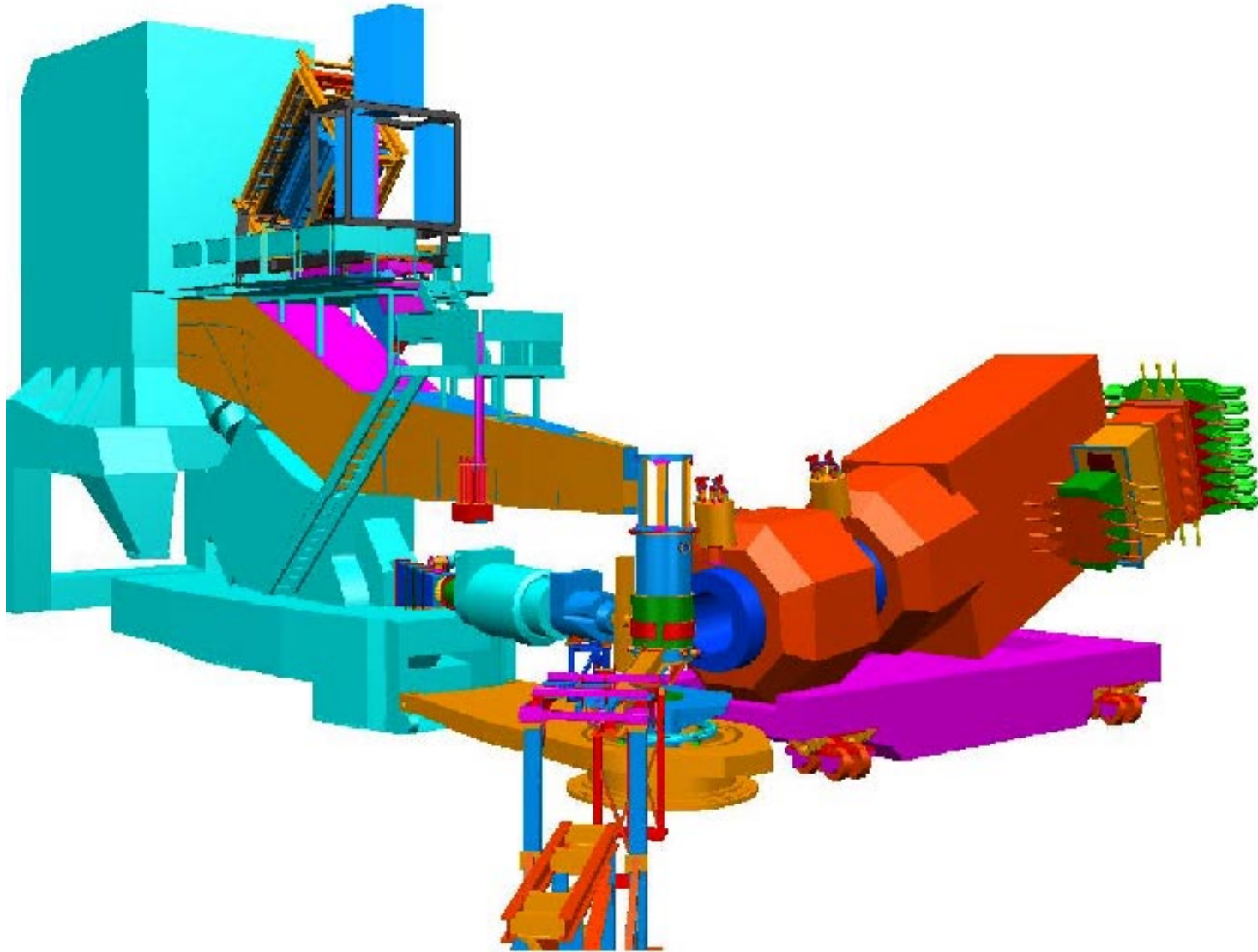
See W. Melnitchouk



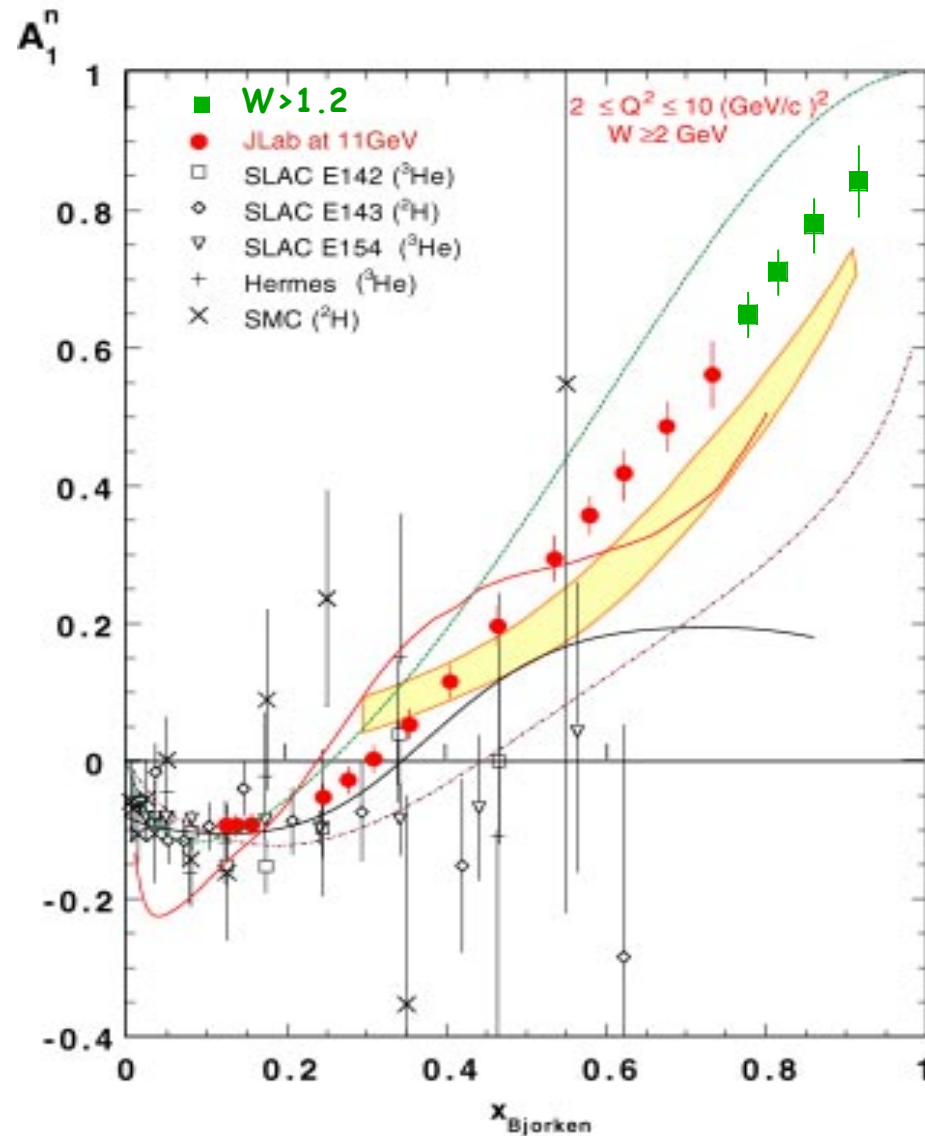
Hint of spin duality for the neutron!



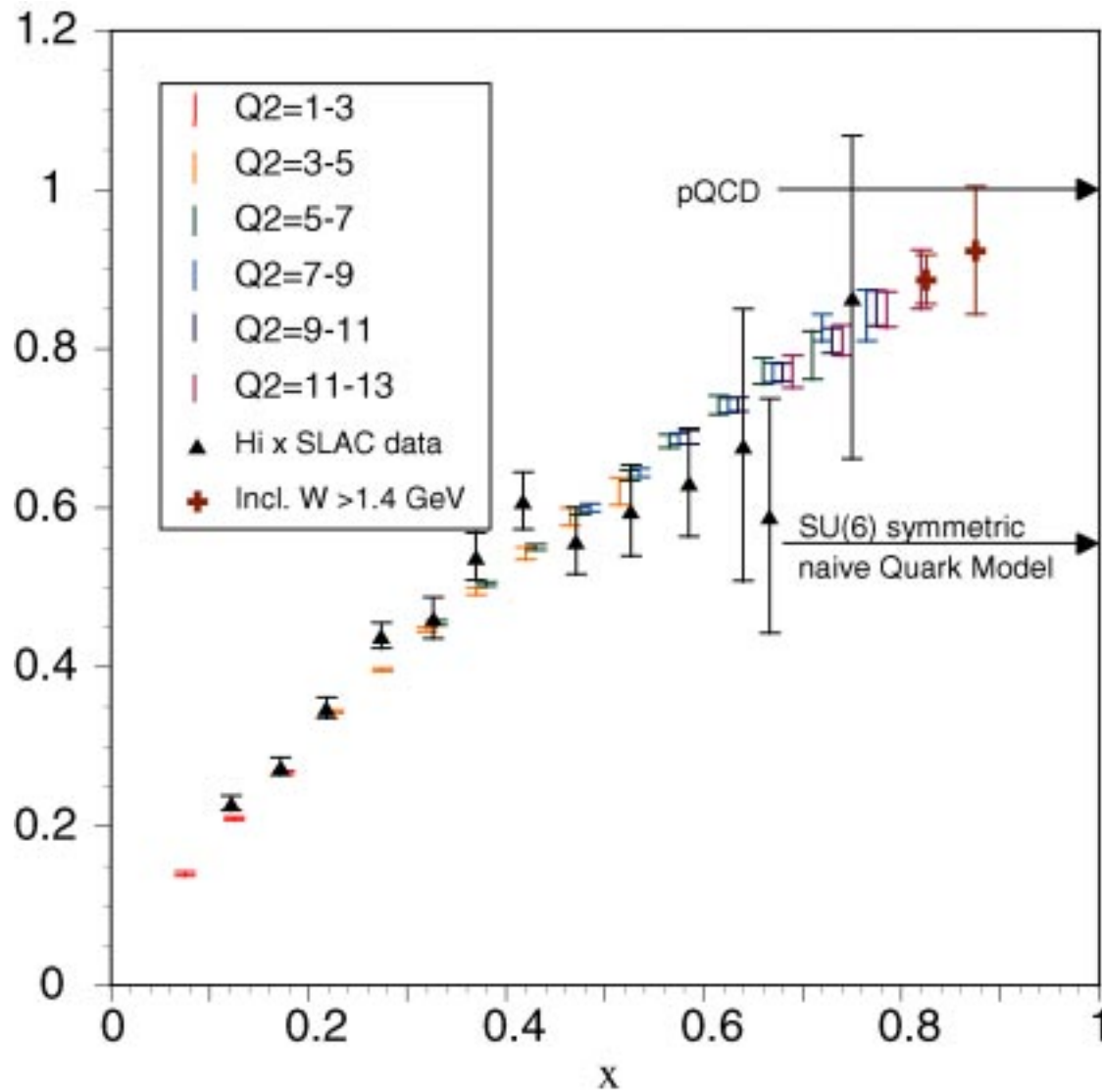
HRS+MAD spectrometers in Hall A



A_1^n at 11 GeV with MAD in Hall A



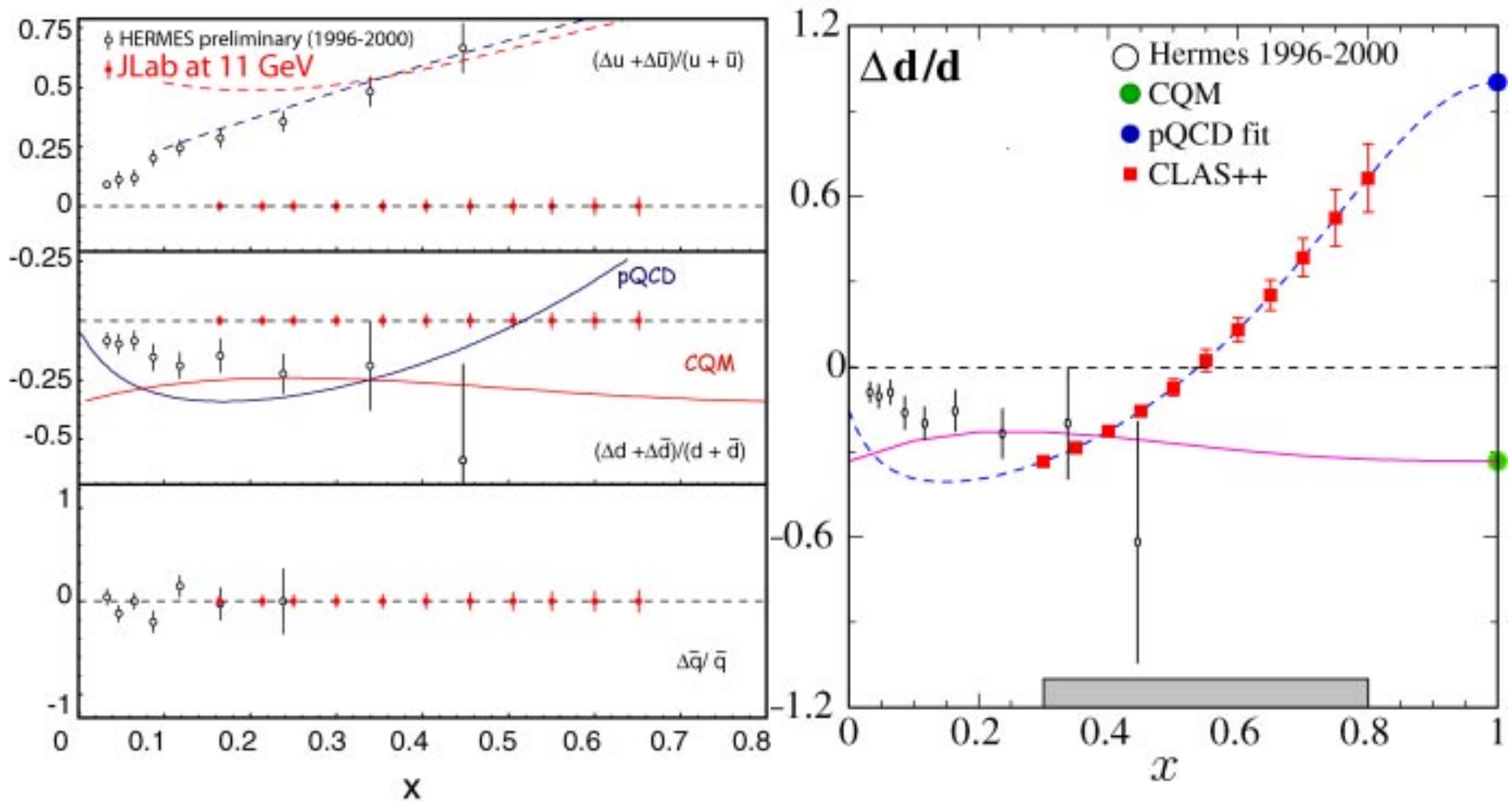
A_1^p at 11 GeV with CLAS++



Helicity-Flavor Decomposition

Hall A with MAD

Hall B with CLAS++



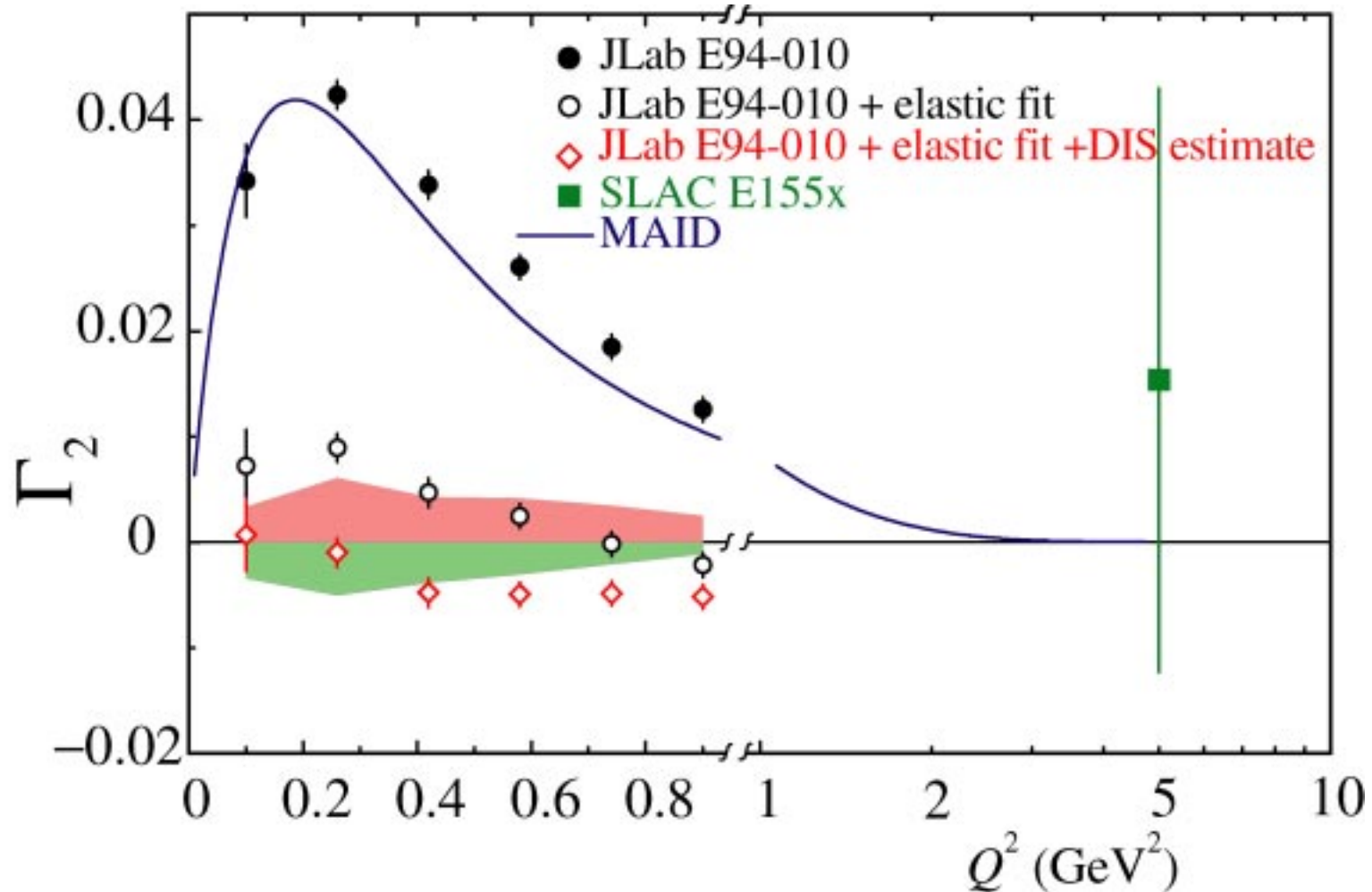
Burkhardt-Cottingham Sum Rule

$$\Gamma_2(Q^2) = \int_0^1 g_2(x, Q^2) dx = 0$$

- Dispersion relation for a spin-flip Compton amplitude
 - ➔ Causality
 - ➔ Analyticity
 - ➔ Absence of a $J=0$ pole with non polynomial residue
- Doesn't follow from Operator Product Expansion and **is valid at all Q^2 if valid at one Q^2**
- Many scenarios of g_2 's low x behavior which would invalidate the sum rule are discussed in the literature.



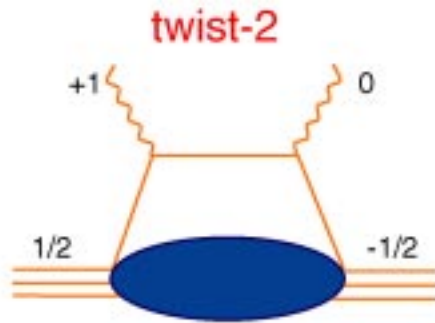
Γ_2 Results of Jlab E94-010



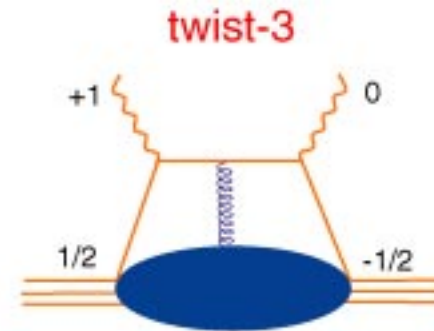
M. Amarian et al. Phys. Rev. Lett. 92, 022301 (2004)



Quark-Gluon Correlations and g_2



Carry one unit of orbital angular momentum



Couple to a gluon

$$g_2(x, Q^2) = g_2^{WW}(x, Q^2) + \bar{g}_2(x, Q^2)$$

- a twist-2 term (Wandzura & Wilczek, 1977):

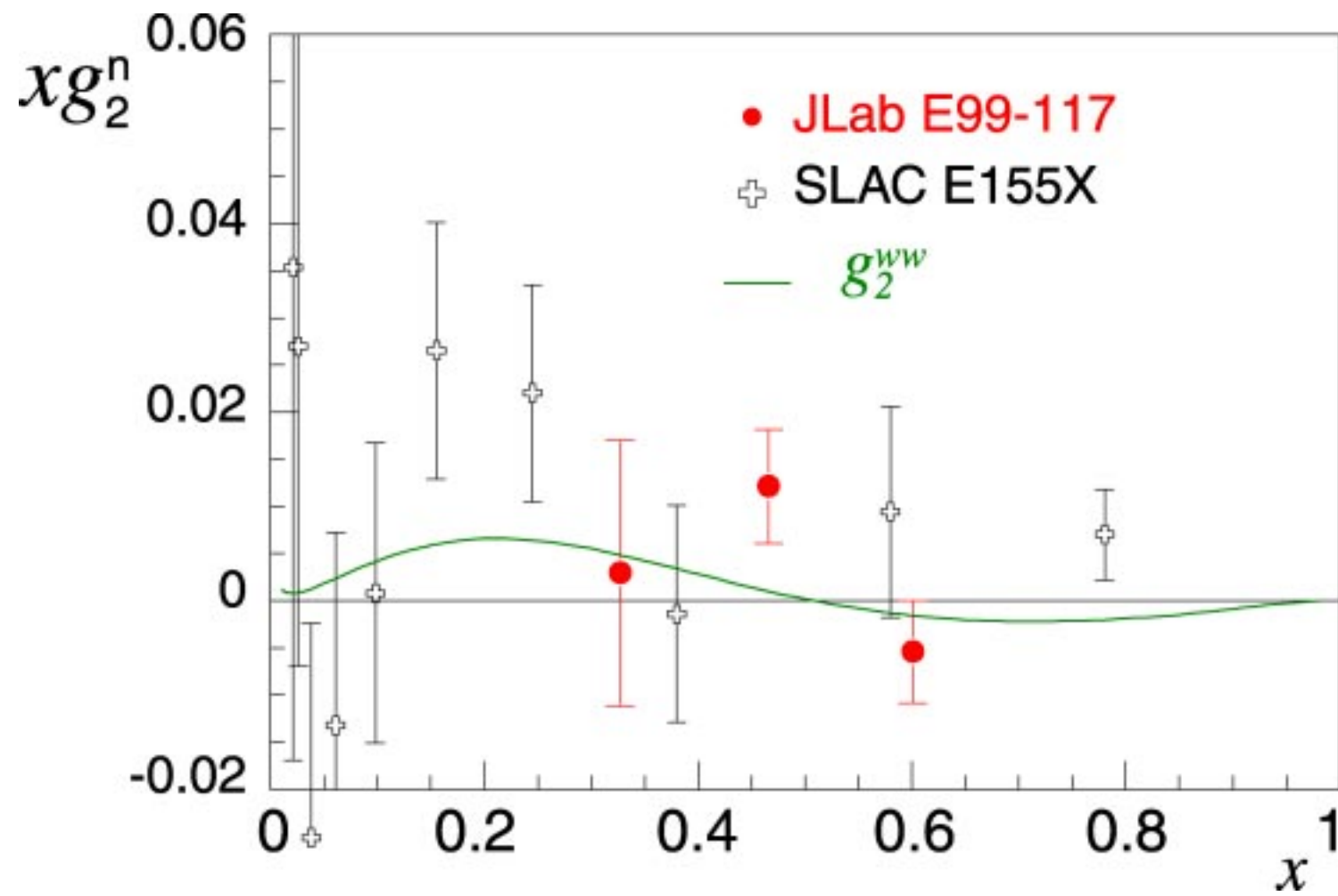
$$g_2^{WW}(x, Q^2) = -g_1(x, Q^2) + \int_0^1 g_1(y, Q^2) \frac{dy}{y}$$

- a twist-3 term with a suppressed twist-2 piece (Cortes, Pire & Ralston, 92):

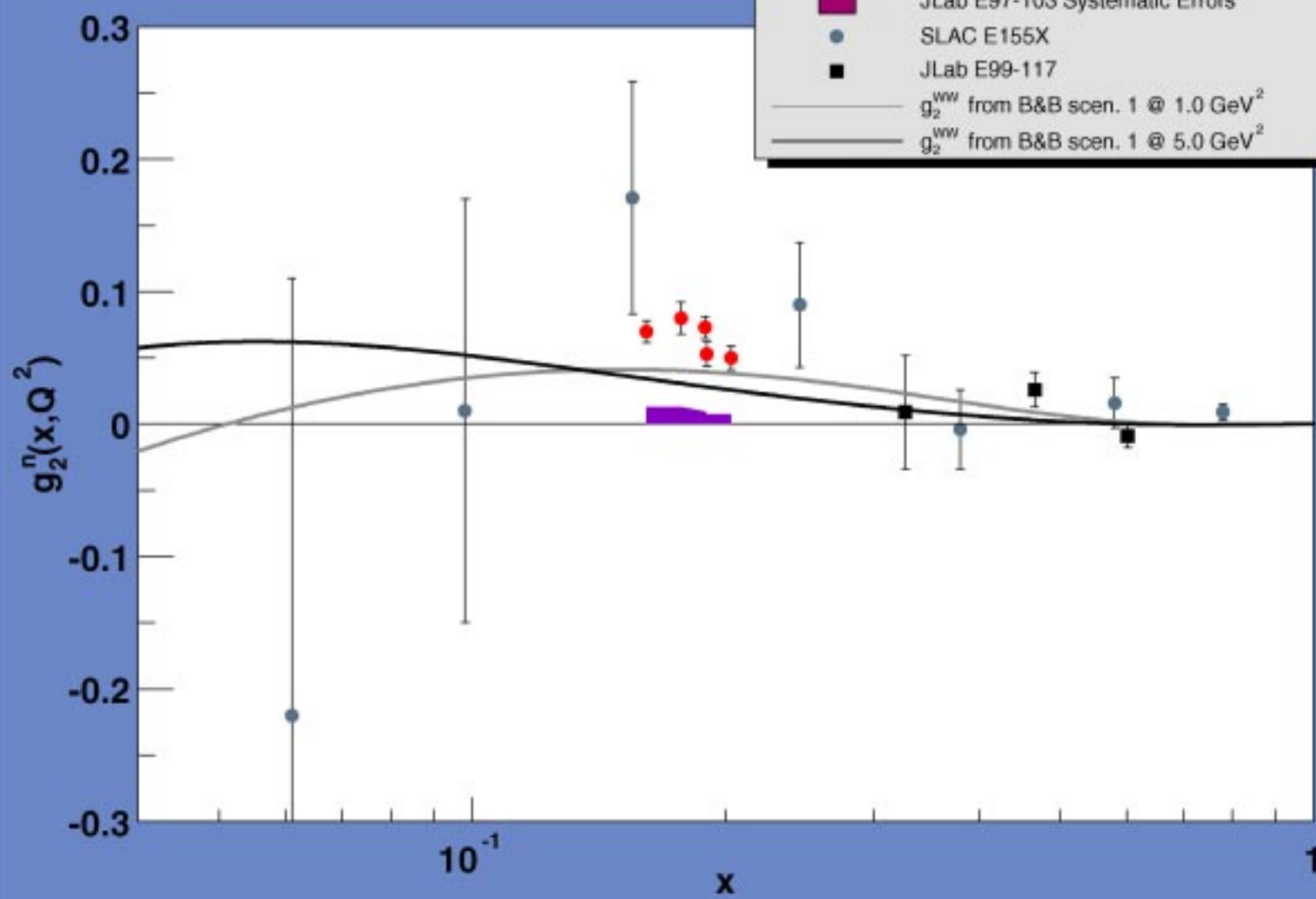
$$\bar{g}_2(x, Q^2) = -\int_x^1 \frac{\partial}{\partial y} \left(\frac{m_q}{M} h_T(y, Q^2) + \xi(y, Q^2) \right) \frac{dy}{y}$$

transversity

quark-gluon correlation



World Data on $g_2^n(x, Q^2)$



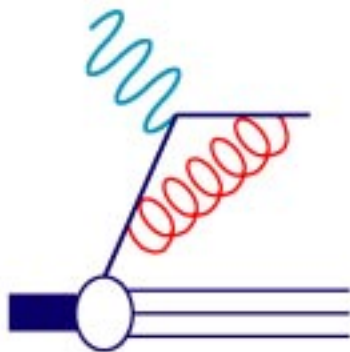
Moments of Structure Functions

$$\Gamma_1(Q^2) = \int_0^1 g_1(x, Q^2) dx = \underbrace{\mu_2}_{\text{leading twist}} + \frac{\mu_4}{Q^2} + \frac{\mu_6}{Q^4} + \dots$$

higher twist

$$\mu_2^{p,n}(Q^2) = \left(\pm \frac{1}{12} g_A + \frac{1}{36} a_8 \right) + \frac{1}{9} \Delta\Sigma + \text{pQCD corrections}$$

$g_A = 1.257$ and $a_8 = 0.579$ are the triplet and octet axial charge, respectively
 $\Delta\Sigma$ = singlet axial charge

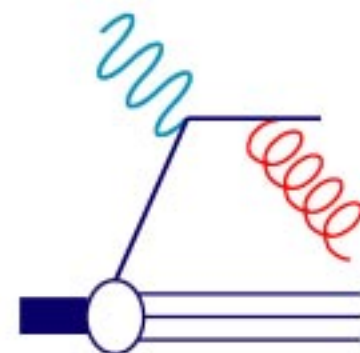


$$g_A = \Delta u - \Delta d$$

$$a_8 = \Delta u + \Delta d - 2\Delta s$$

$$\Delta\Sigma = \Delta u + \Delta d + \Delta s$$

pQCD radiative corrections



Study of Higher Twists (continued)

$$\mu_4(Q^2) = \frac{M^2}{9} [a_2(Q^2) + 4d_2(Q^2) + 4f_2(Q^2)]$$

Twist - 2 Twist - 3 Twist - 4
(TMC)

where a_2 , d_2 and f_2 are higher moments of g_1 and g_2

e.g. $d_2(Q^2) = \int_0^1 x^2 [2g_1(x, Q^2) + 3g_2(x, Q^2)] dx = \int_0^1 x^2 \overline{g_2}(x, Q^2) dx$

$$a_2(Q^2) = \int_0^1 x^2 g_1(x, Q^2) dx$$

- To extract f_2 , d_2 needs to be determined first.
- Both d_2 and f_2 are required to determine the color polarizabilities



Color "polarizabilities"

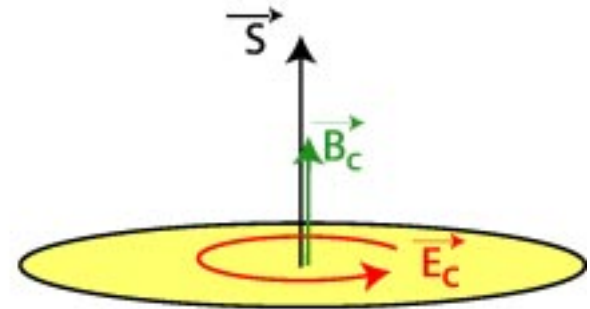
How does the gluon field respond when a nucleon is polarized ?

Define color magnetic and electric polarizabilities (in nucleon rest frame):

$$\chi_{B,E} 2M^2 \vec{S} = \langle PS | \vec{O}_{B,E} | PS \rangle$$

where $\vec{O}_B = \psi^\dagger g \vec{B} \psi$

$$\vec{O}_E = \psi^\dagger \vec{\alpha} \times g \vec{E} \psi$$

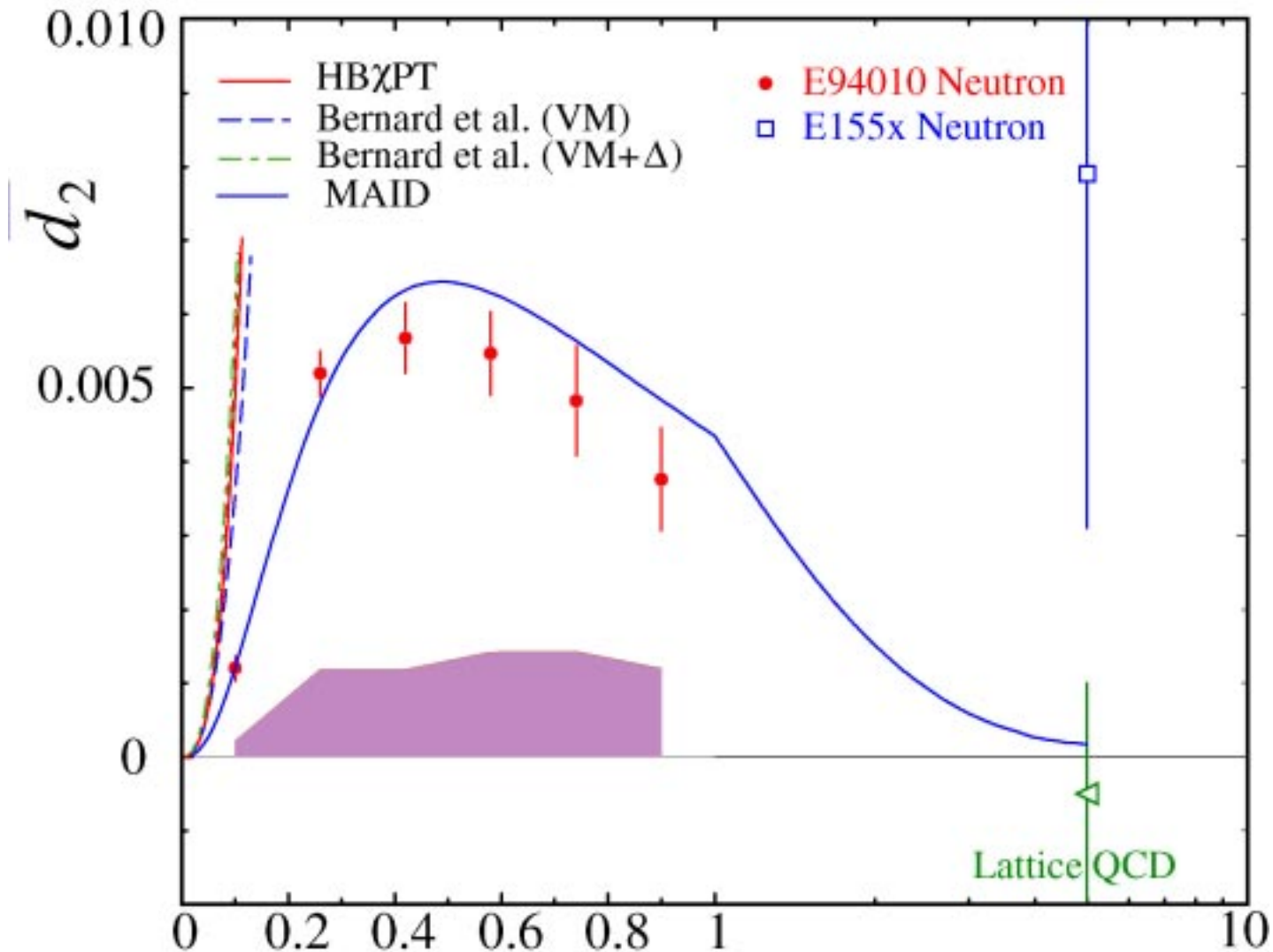


$$\chi_E^n = (4d_2^n + 2f_2^n)/3$$

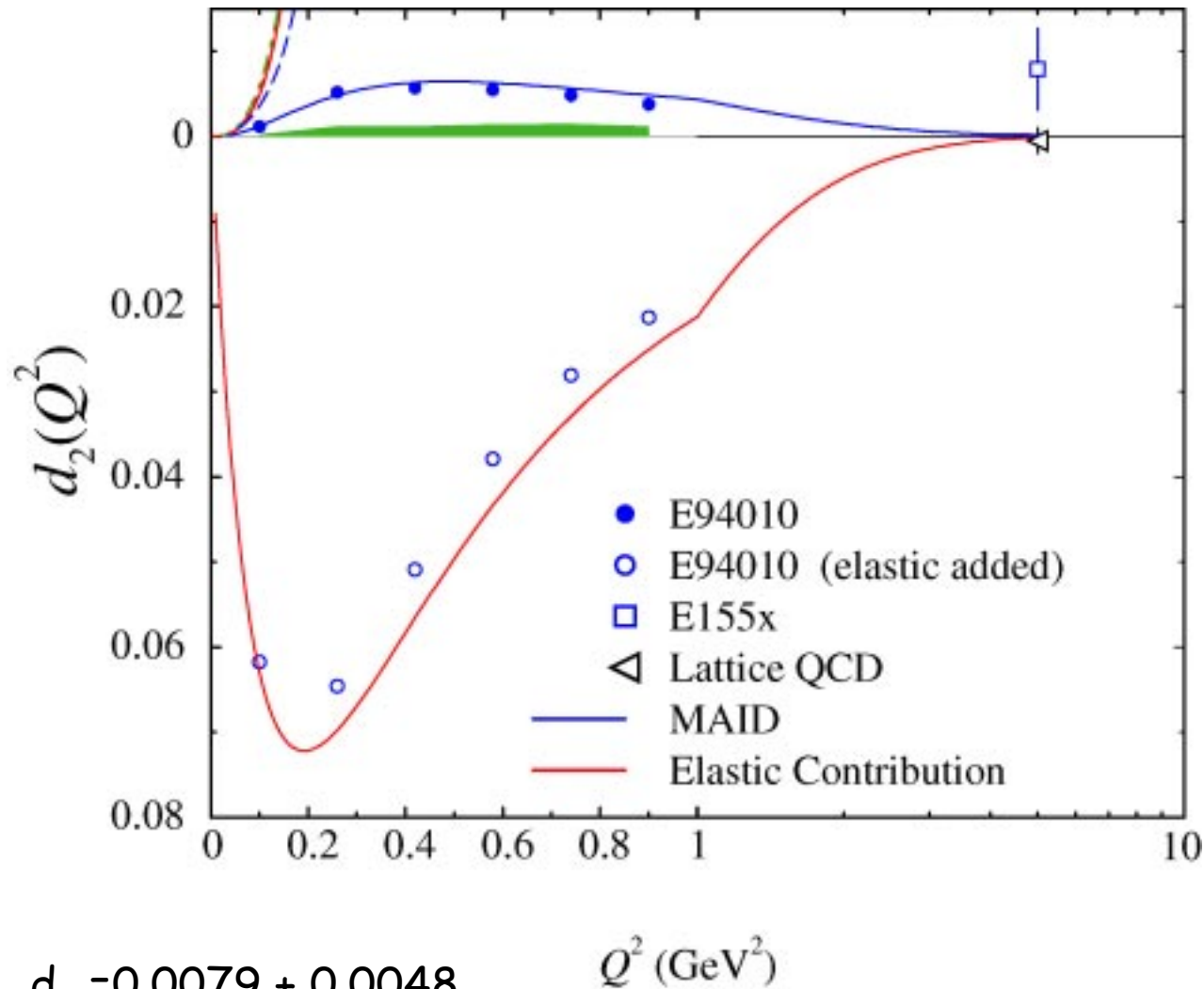
$$\chi_B^n = (4d_2^n - f_2^n)/3$$

χ_E and χ_B represent the response of the color \vec{B} & \vec{E} fields to the nucleon polarization

d_2 results of Jlab E94-010



Adding the elastic contribution

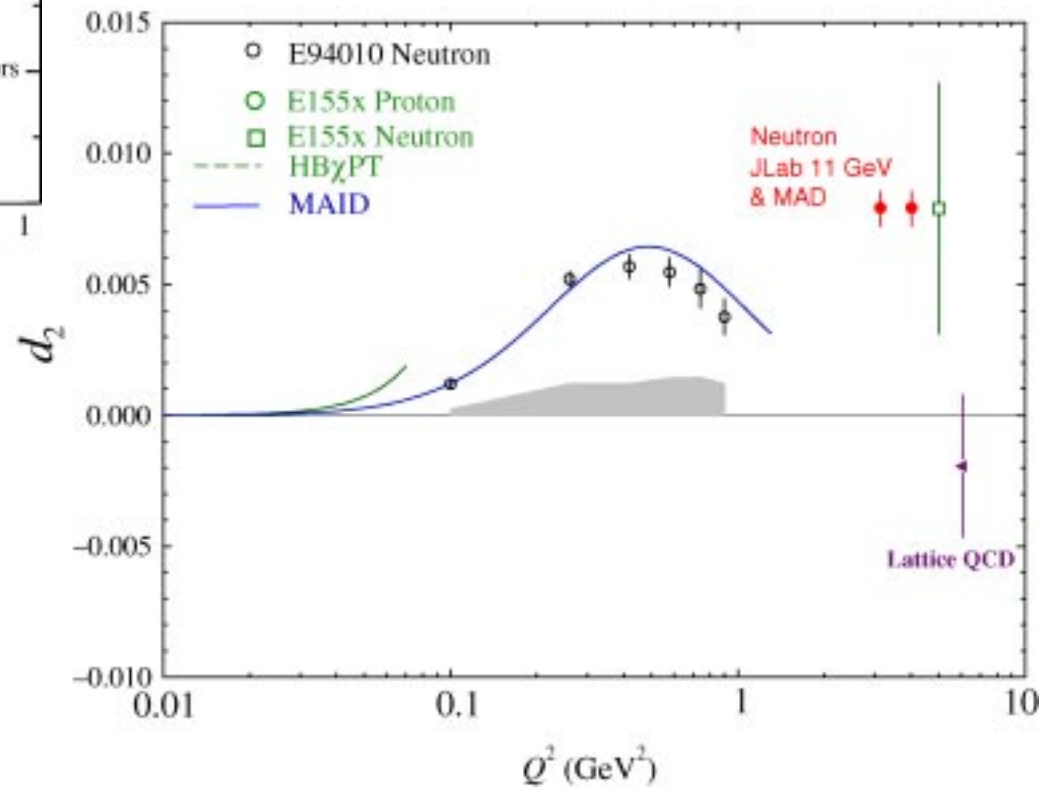
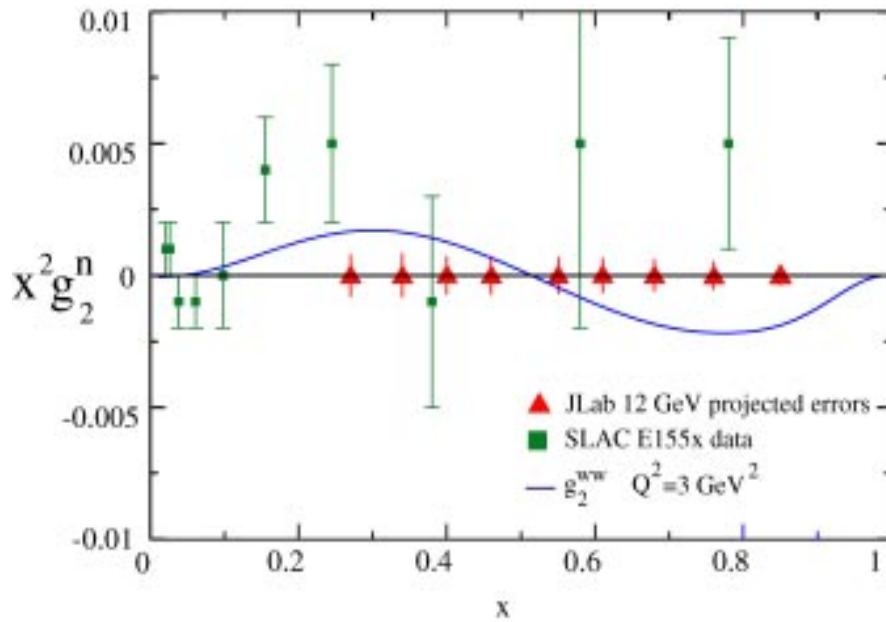


E155x: $d_2 = 0.0079 \pm 0.0048$

Updated value using E99-117: $d_2 = 0.0062 \pm 0.0028$



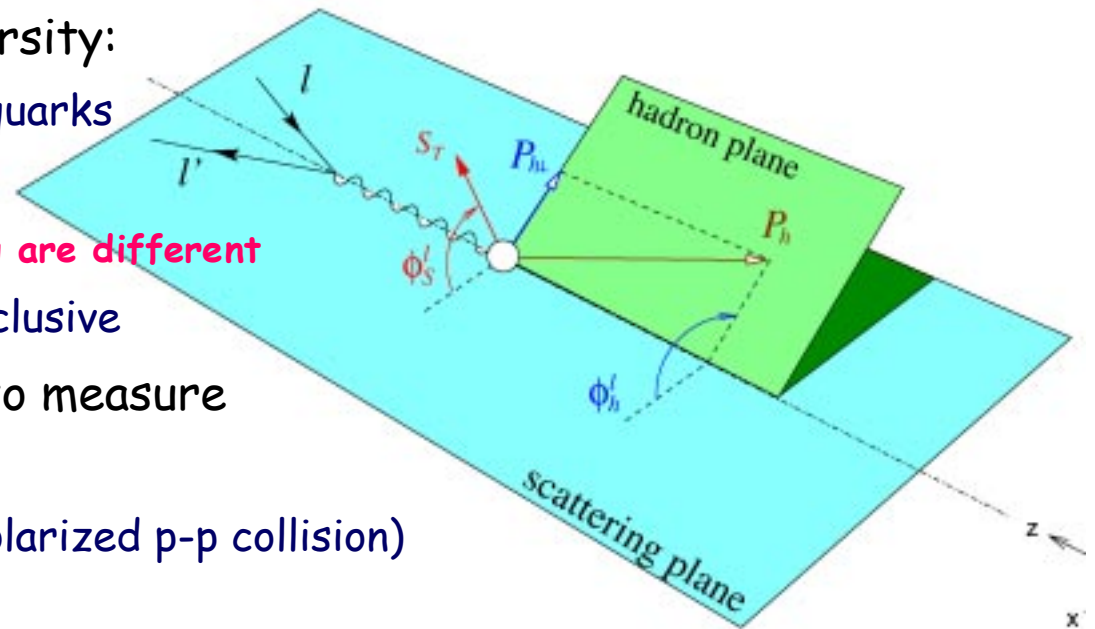
g_2 and d_2 with 11 GeV at JLab



Semi-Inclusive/Transversity

- Three twist-2 quark distributions:
 - ➔ Density distribution: $q(x, Q^2) = q^\uparrow(x) + q^\downarrow(x)$
 - ➔ Helicity distribution: $\Delta q(x, Q^2) = q^\uparrow(x) - q^\downarrow(x)$
 - ➔ Transversity distribution: $\delta q(x, Q^2) = q^\perp(x) + q_\perp(x)$
- Some characteristics of transversity:
 - ➔ $\delta q(x) = \Delta q(x)$ for non relativistic quarks
 - ➔ δq and gluons do not mix
 - ↪ --> Q^2 evolution for δq and Δq are different
 - ➔ Chiral-odd --> not accessible in inclusive
- It takes two chiral-odd objects to measure transversity
 - ➔ Drell-Yan (doubly transversely polarized p-p collision)
 - ➔ Semi-inclusive DIS
 - ↪ Chiral-odd distribution function (**transversity**)
 - ↪ Chiral-Odd fragmentation function (**Collins function**)

See N. Makins

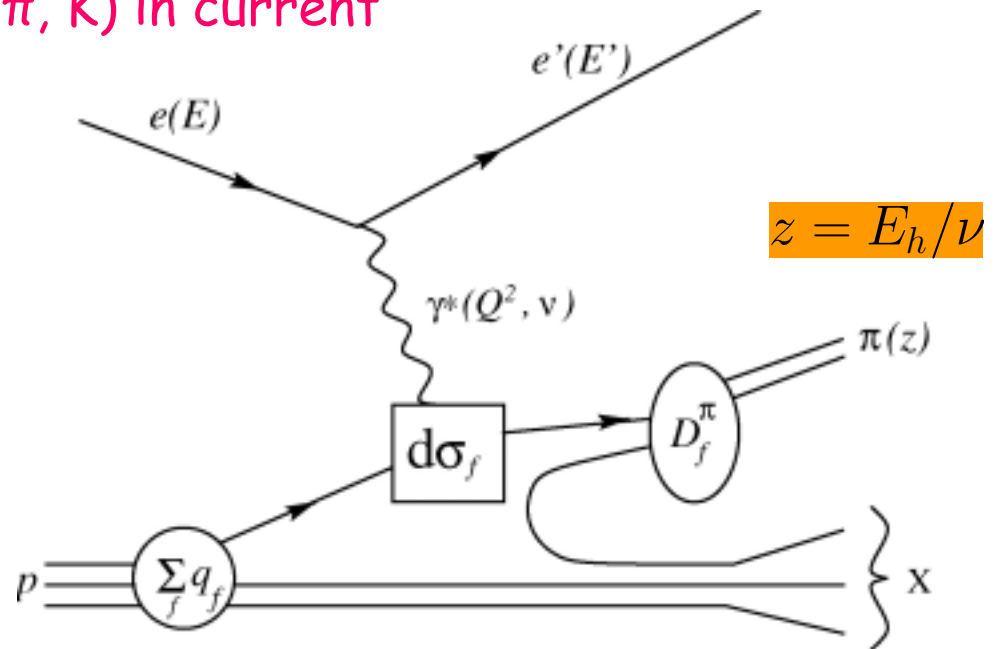


Semi-inclusive DIS

- Spin-flavor decomposition of valence and sea quarks by tagging hadron (e.g. π , K) in current fragmentation region

$$d\sigma = \sum_f e_f^2 q_f(x) D_f^h(z)$$

(z) quark \rightarrow hadron fragmentation function



- unpolarized or polarized beam and target
- mass of unobserved X system, $W_X > 2 \text{ GeV}$

All Eight Quark Distributions Are Probed in Semi-Inclusive DIS

$$d^6\sigma = \frac{4\pi\alpha^2 sx}{Q^4} \times$$

$f_1 =$	$\{ [1 + (1-y)^2] \sum_{q,\bar{q}} e_q^2 f_1^q(x) D_1^q(z, P_{h\perp}^2)$	Unpolarized
$h_1^\perp =$	$+ (1-y) \frac{P_{h\perp}^2}{4z^2 M_N M_h} \cos(2\phi_h^l) \sum_{q,\bar{q}} e_q^2 h_1^{\perp(1)q}(x) H_1^{\perp q}(z, P_{h\perp}^2)$	
$h_{1L}^\perp =$	$- S_L (1-y) \frac{P_{h\perp}^2}{4z^2 M_N M_h} \sin(2\phi_h^l) \sum_{q,\bar{q}} e_q^2 h_{1L}^{\perp(1)q}(x) H_1^{\perp q}(z, P_{h\perp}^2)$	
Transversity $h_{1T} =$	$+ S_T (1-y) \frac{P_{h\perp}}{zM_h} \sin(\phi_h^l + \phi_S^l) \sum_{q,\bar{q}} e_q^2 h_1^q(x) H_1^{\perp q}(z, P_{h\perp}^2)$	Polarized target
Sivers $f_{1T}^\perp =$	$+ S_T (1-y + \frac{1}{2}y^2) \frac{P_{h\perp}}{zM_N} \sin(\phi_h^l - \phi_S^l) \sum_{q,\bar{q}} e_q^2 f_{1T}^{\perp(1)q}(x) D_1^q(z, P_{h\perp}^2)$	
$h_{1T}^\perp =$	$+ S_T (1-y) \frac{P_{h\perp}^3}{6z^3 M_N^2 M_h} \sin(3\phi_h^l - \phi_S^l) \sum_{q,\bar{q}} e_q^2 h_{1T}^{\perp(2)q}(x) H_1^{\perp q}(z, P_{h\perp}^2)$	
$g_{1L} =$	$+ \lambda_e S_L y (1 - \frac{1}{2}y) \sum_{q,\bar{q}} e_q^2 g_1^q(x) D_1^q(z, P_{h\perp}^2)$	Polarized beam and target
$g_{1T} =$	$+ \lambda_e S_T y (1 - \frac{1}{2}y) \frac{P_{h\perp}}{zM_N} \cos(\phi_h^l - \phi_S^l) \sum_{q,\bar{q}} e_q^2 g_{1T}^{(1)q}(x) D_1^q(z, P_{h\perp}^2) \}$	

S_L and S_T : Target Polarizations; λ_e : Beam Polarization



- **Transversity** “Collins effect”: Finding in a polarized target nucleon a transverse-polarized quark, which fragments with a transverse momentum correlated with that quark polarization

$$h_1 = \begin{array}{c} \uparrow \\ \circlearrowleft \\ \uparrow \\ \bullet \\ \downarrow \end{array} - \begin{array}{c} \uparrow \\ \circlearrowright \\ \uparrow \\ \bullet \\ \downarrow \end{array} \quad \times \quad H_1^\perp = \begin{array}{c} \uparrow \\ \circlearrowleft \\ \uparrow \\ \bullet \\ \downarrow \end{array} - \begin{array}{c} \uparrow \\ \circlearrowright \\ \uparrow \\ \bullet \\ \downarrow \end{array}$$

chiral-odd, T-even chiral-odd, T-odd

- “Sivers effect”: Finding in a transverse-polarized target nucleon a quark with correlated primordial transverse momentum

$$f_{1T}^\perp = \begin{array}{c} \uparrow \\ \circlearrowleft \\ \bullet \end{array} - \begin{array}{c} \uparrow \\ \circlearrowright \\ \bullet \end{array} \quad \times \quad D_1 = \begin{array}{c} \bullet \end{array}$$

chiral-even, T-odd chiral-even, T-even

- (Boer & Mulders) Finding in an unpolarized target nucleon a quark with correlated transverse polarization and primordial transverse momentum

$$h_1^\perp = \begin{array}{c} \uparrow \\ \circlearrowleft \\ \uparrow \\ \bullet \\ \downarrow \end{array} - \begin{array}{c} \uparrow \\ \circlearrowright \\ \uparrow \\ \bullet \\ \downarrow \end{array} \quad \times \quad H_1^\perp = \begin{array}{c} \uparrow \\ \circlearrowleft \\ \uparrow \\ \bullet \\ \downarrow \end{array} - \begin{array}{c} \uparrow \\ \circlearrowright \\ \uparrow \\ \bullet \\ \downarrow \end{array}$$

chiral-odd, T-odd chiral-odd, T-odd

Jlab Hall A E03-004 / $^3\text{He}^\uparrow (e, e'\pi^-)X$

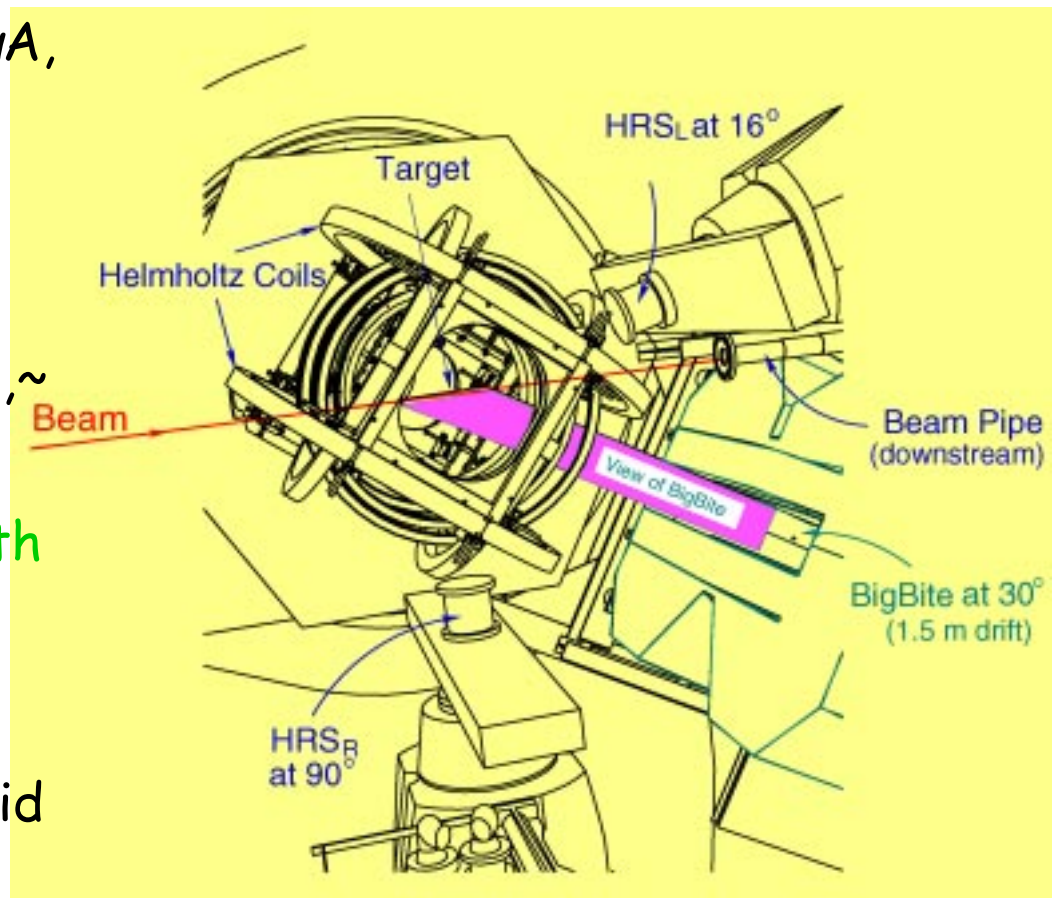
- Beam

- ➔ Polarized ($P \sim 80\%$) e^- , $15 \mu\text{A}$, helicity flip at 60Hz

Spokespeople: J.-P. Chen, X. Jiang & J.-C. Peng

- Target

- ➔ Optically pumped Rb+spin exchange ^3He , 50 mg/cm^2 , $\sim 40\%$ polarization
- ➔ Transversely polarized with tunable direction



- Electron detection

- ➔ Bigbite spectrometer, Solid angle 60 msr , $\theta = 30^\circ$

- Charged pion detection

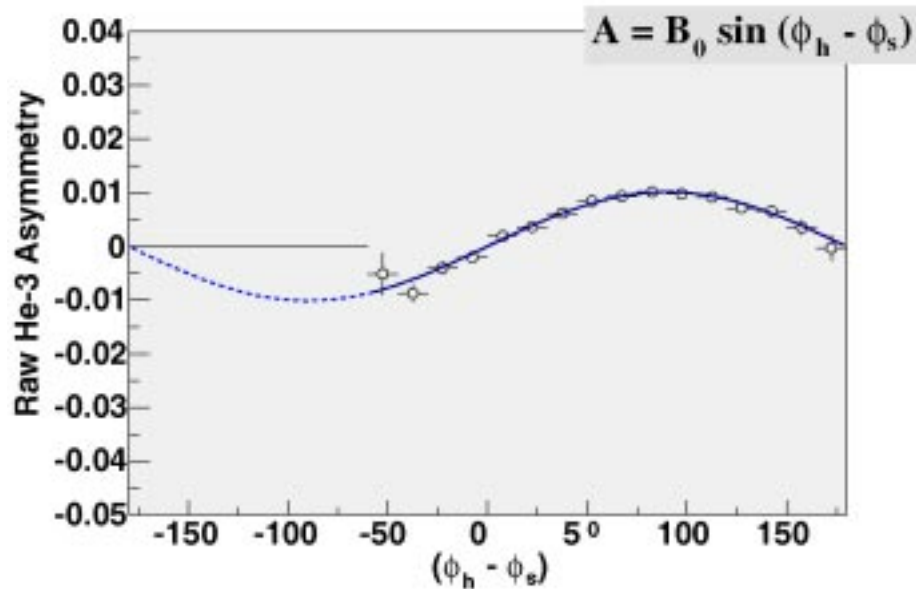
- ➔ HRS spectrometer, $\theta = 16^\circ$



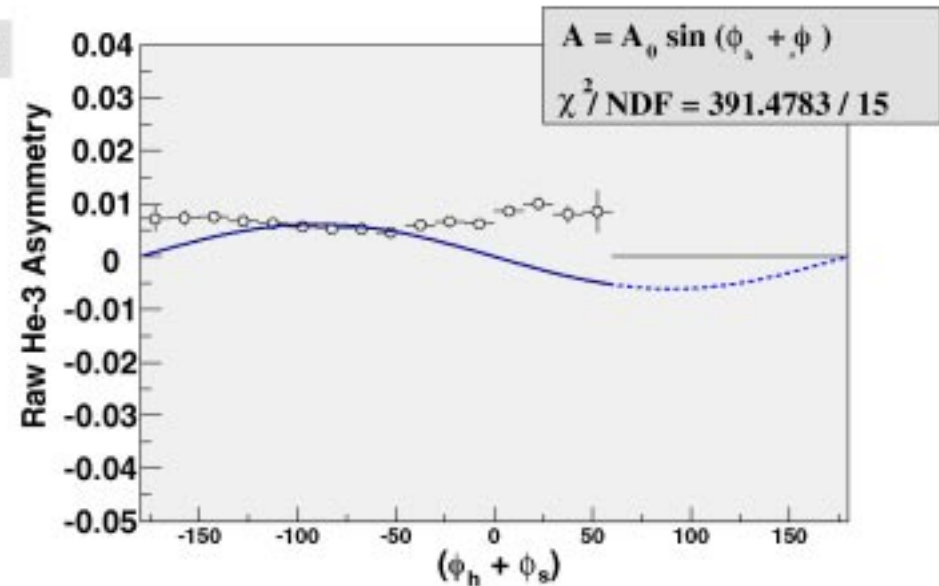
Disentangling Collins and Sivers Effects

Monte Carlo assuming 1.0% asymmetry due to **Sivers** effect

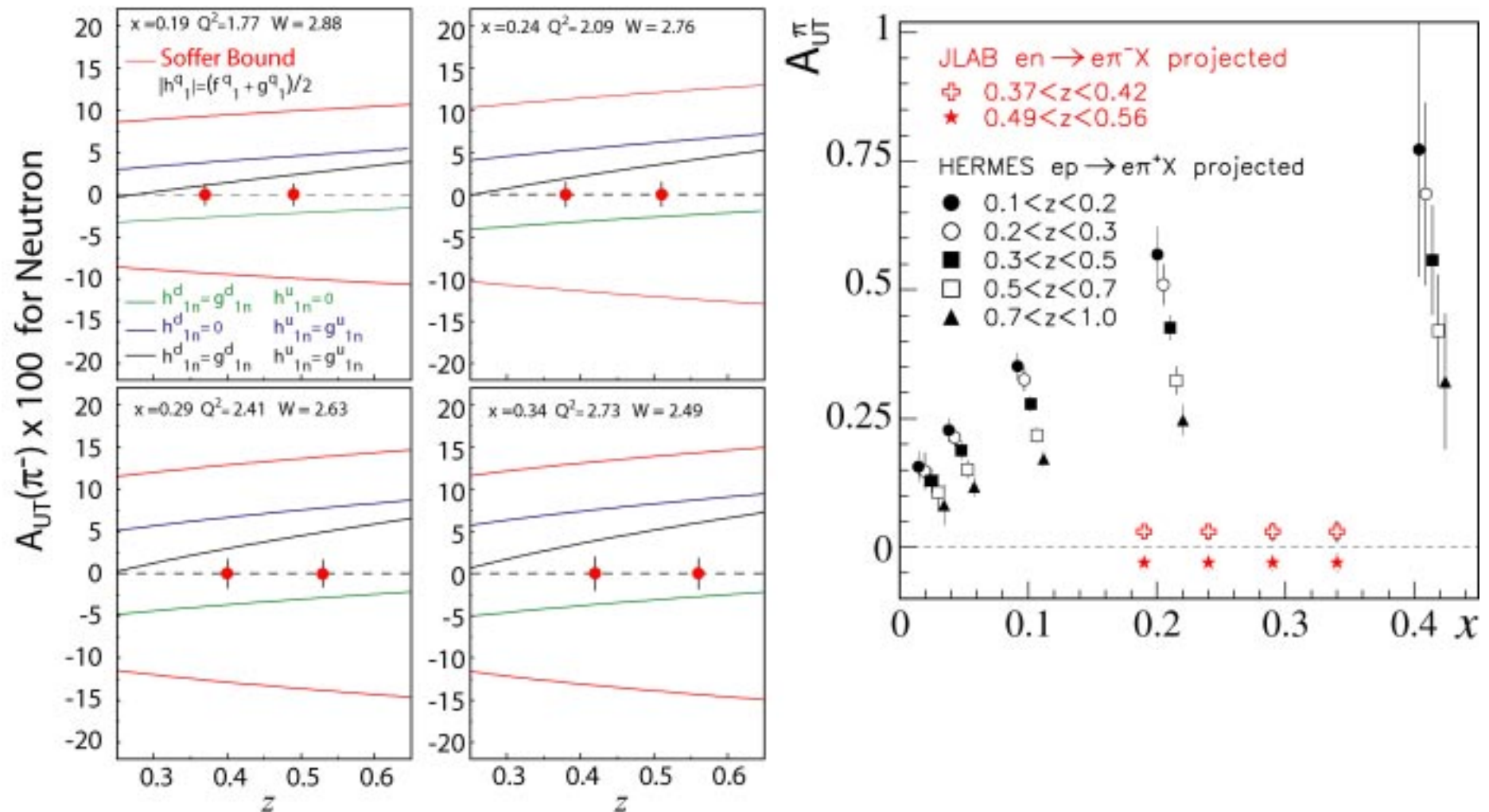
Asymmetry versus **Sivers** angle



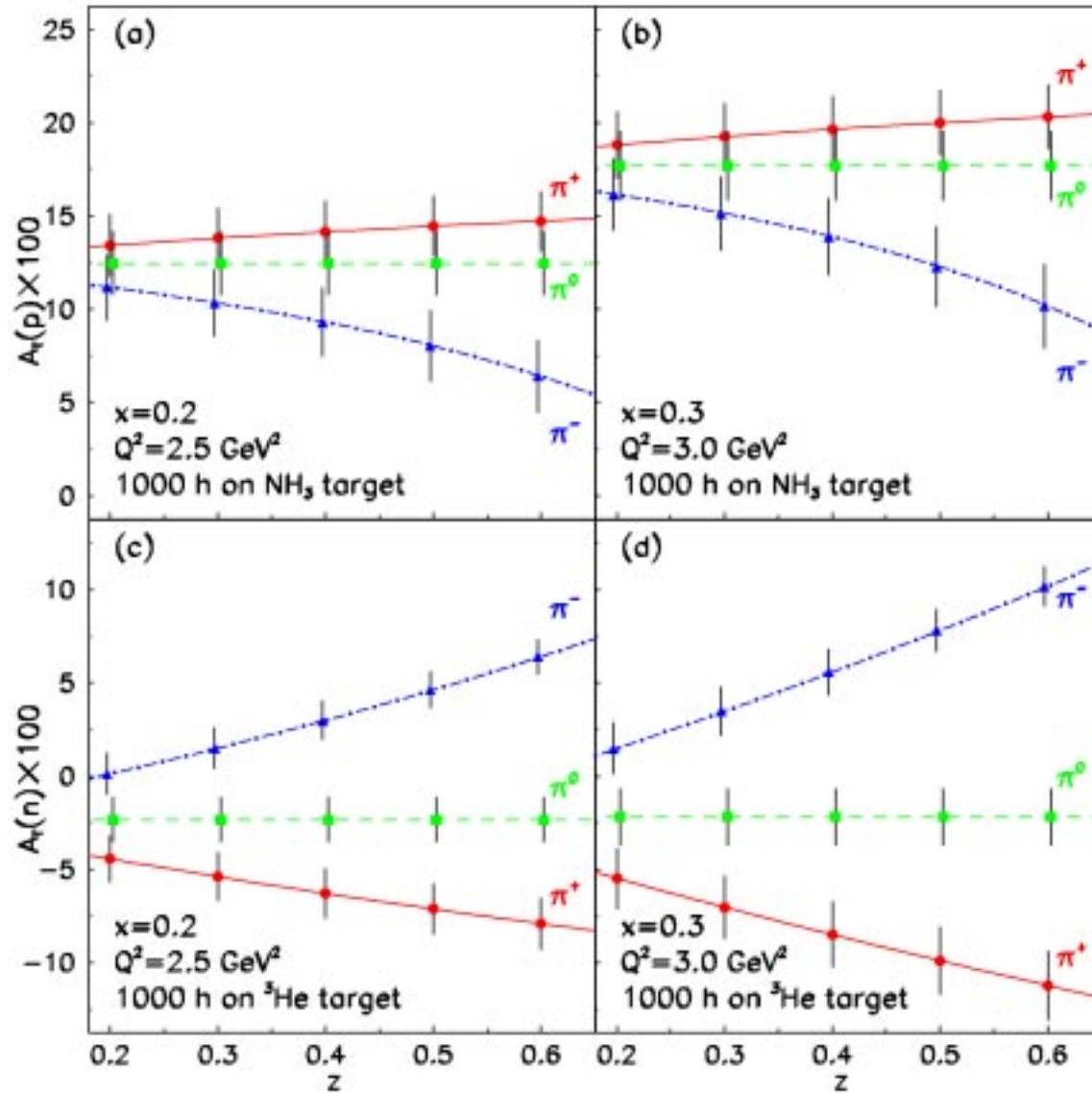
Asymmetry versus **Collins** Angle



Expected result from JLab at 6 GeV



Jlab 12 GeV projections

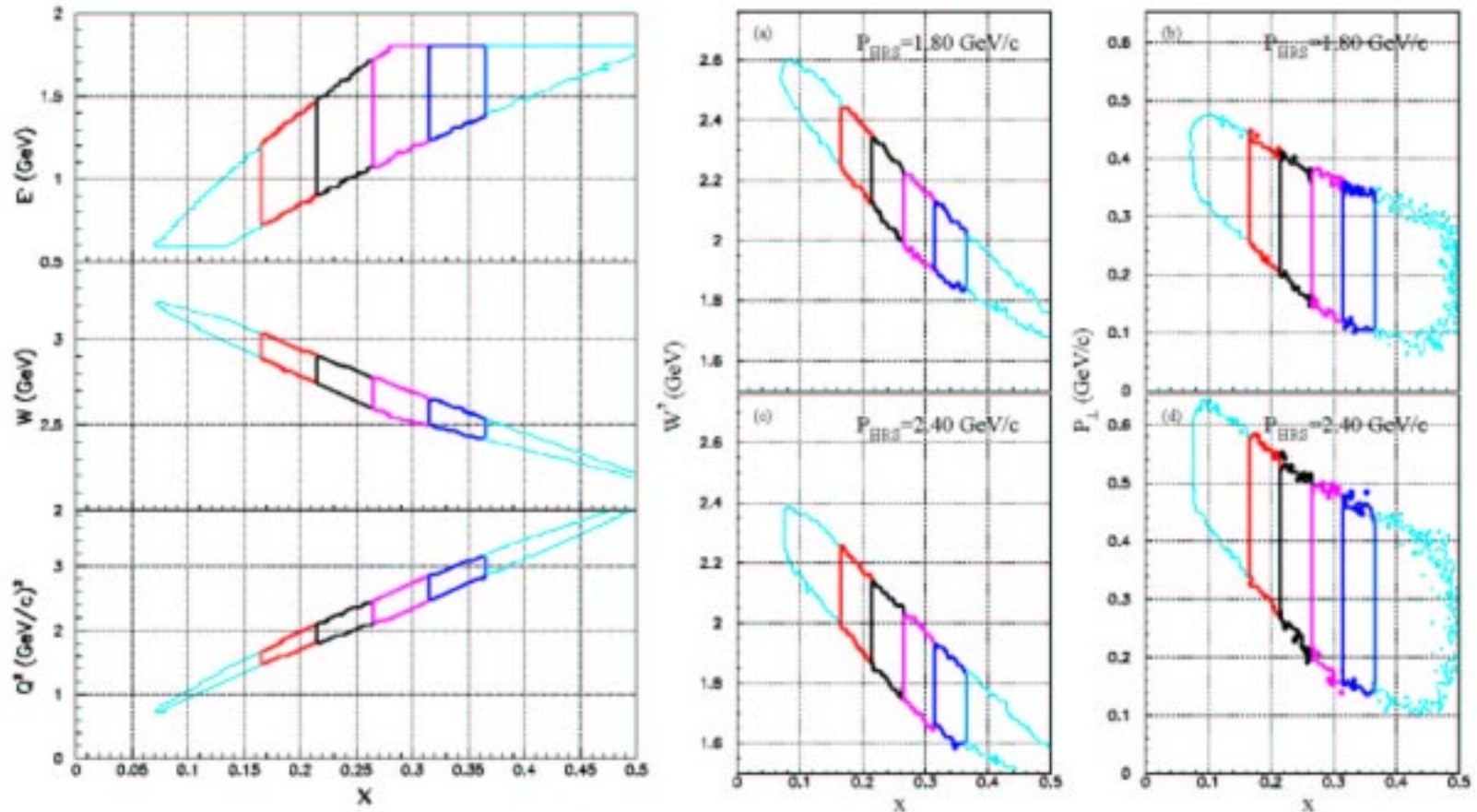


Summary

- Neutron A_1 crosses zero around $x \sim 0.4$ and becomes positive at large x .
- $\Delta d/d$ is negative up to $x=0.6$ consistent with RCQM models but not with hadron helicity conservation assumption.
- The Burkhardt-Cottingham sum rule seems verified within errors for $Q^2 < 1 \text{ GeV}^2$.
- Neutron d_2 is small but finite in the resonance region. Precision measurements of g_2 in the range $1 < Q^2 < 4 \text{ GeV}^2$ are required for an accurate extraction of color polarizabilities.
- 11 GeV at JLab will allow to extend the investigation of the nucleon spin structure in the large x region through inclusive and semi-inclusive reactions, this includes transversity and other transverse momentum distributions.



Kinematic Acceptance



Hall-A : $x: 0.19 - 0.34, Q^2: 1.8 - 2.7 \text{ GeV}^2, W : 2.5 - 2.9 \text{ GeV}, z: 0.37 - 0.56$

HERMES: $\langle Q^2 \rangle = 2.5 \text{ GeV}^2$

