

Latest results on the strong coupling at the LHC



Seminar at the
Department of Physics
University of Oxford

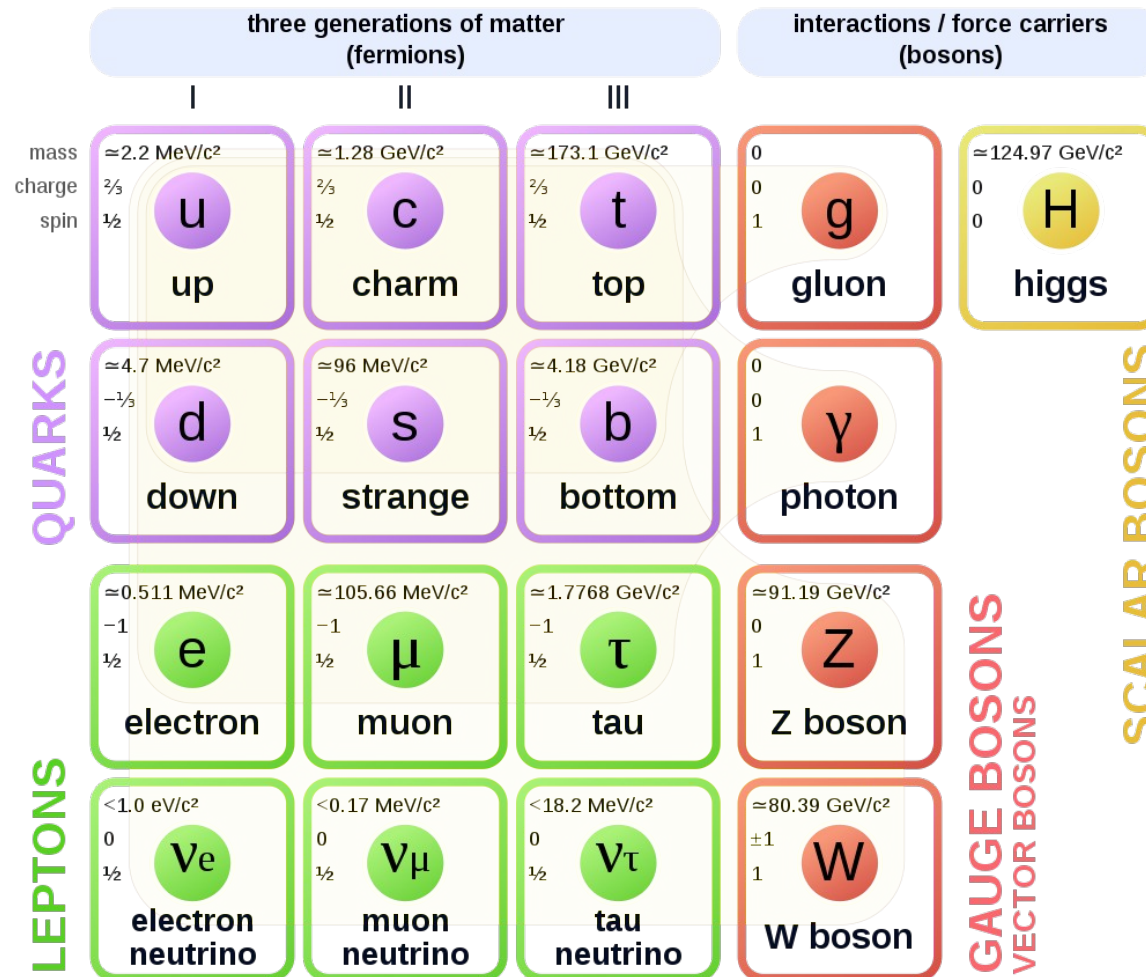
Outline

- Motivation
- Status of α_s
- New results from LHC
 - Jet cross sections
 - Normalised distributions
 - Ratio observables
- Outlook





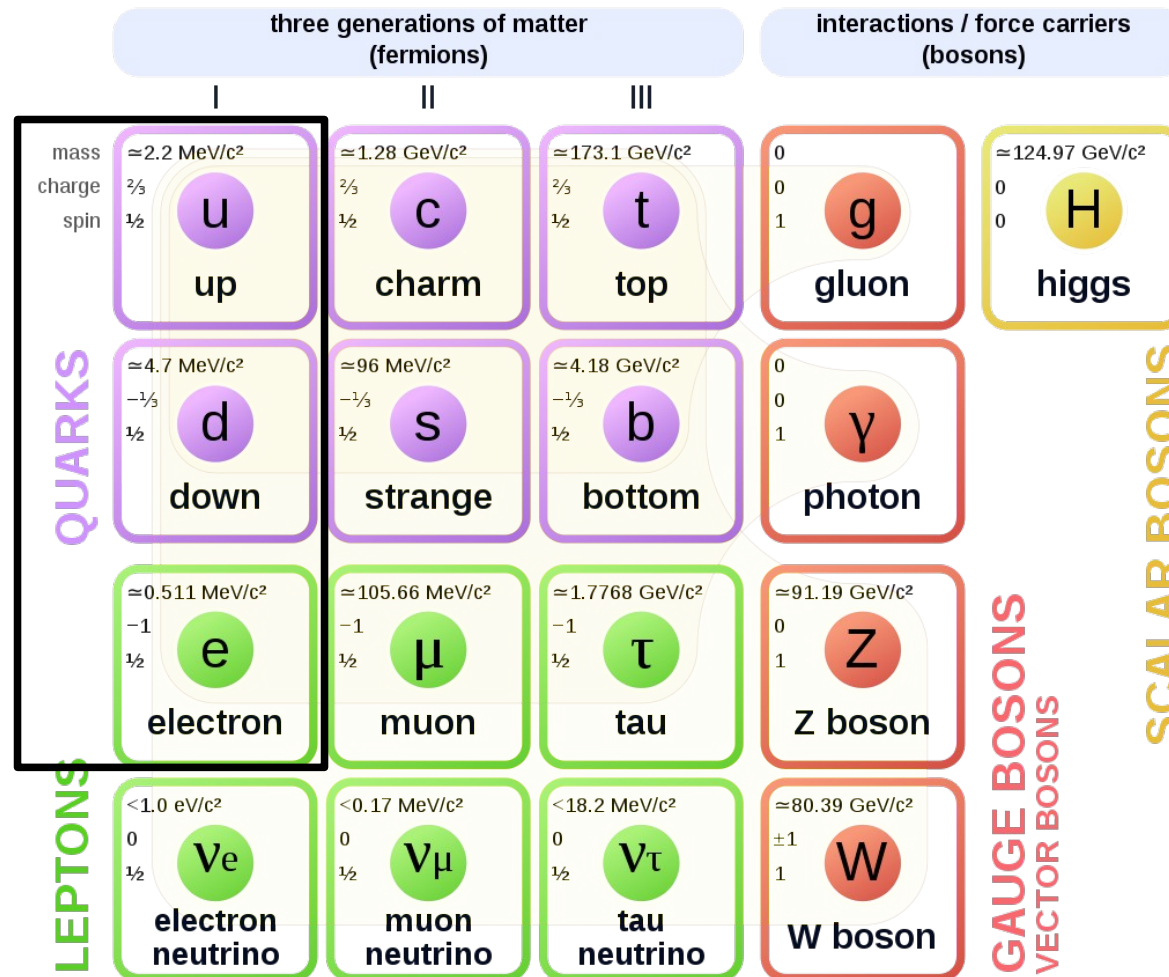
Standard Model of Elementary Particles



Cush, Wikipedia.



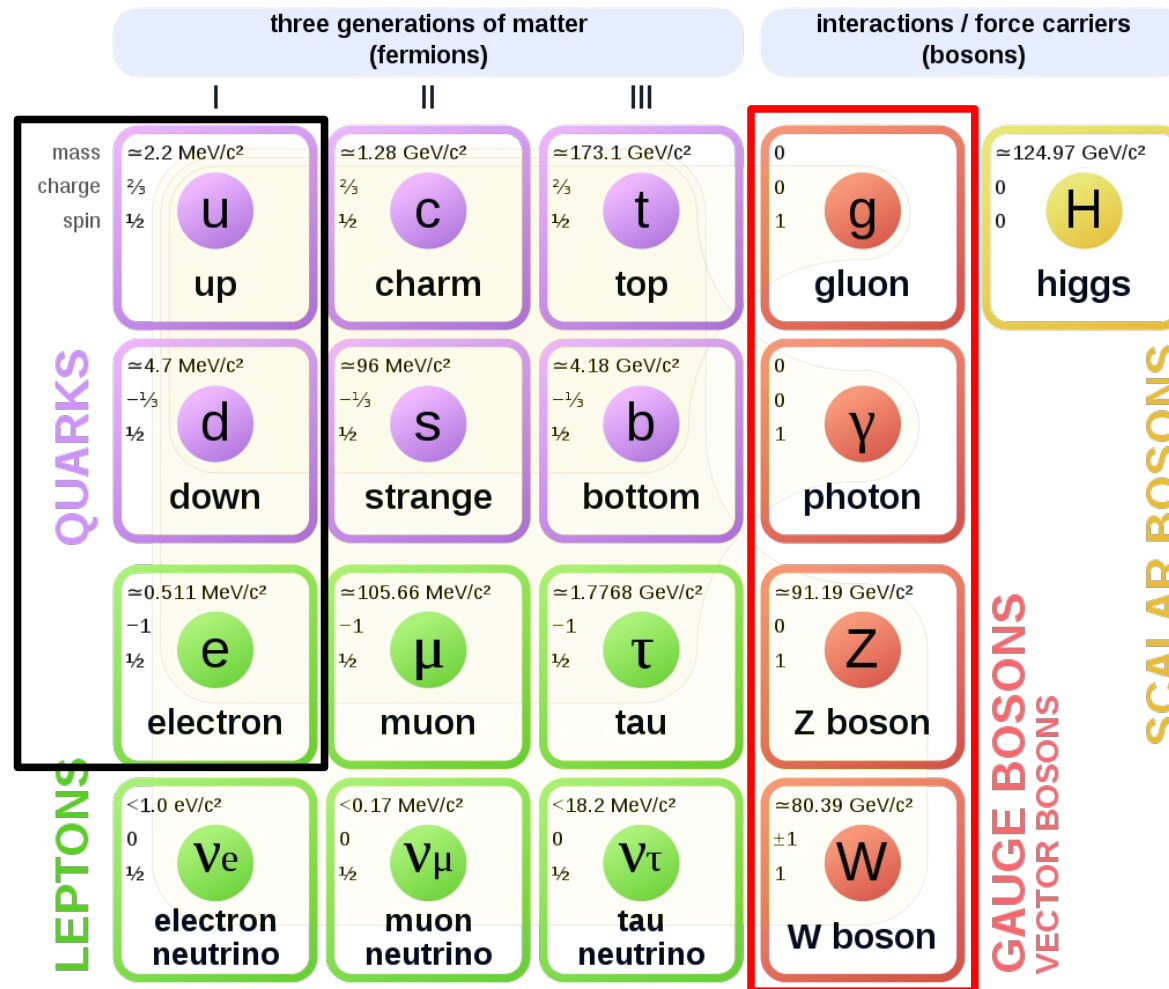
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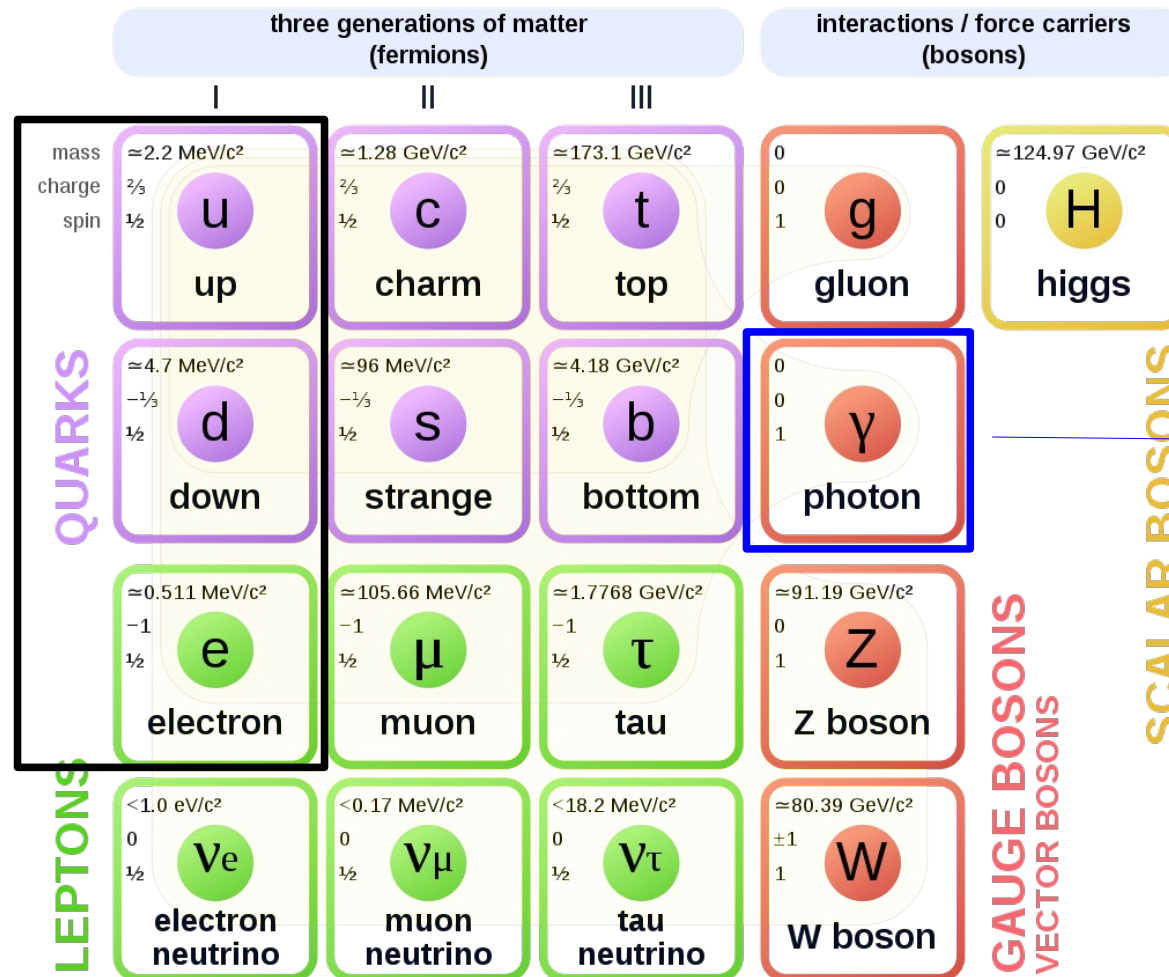


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... and three fundamental interactions.
(no gravity)



Standard Model of Elementary Particles



Electromagnetic interaction
(magnets, electricity, ...)

$$\alpha \approx 1/137$$

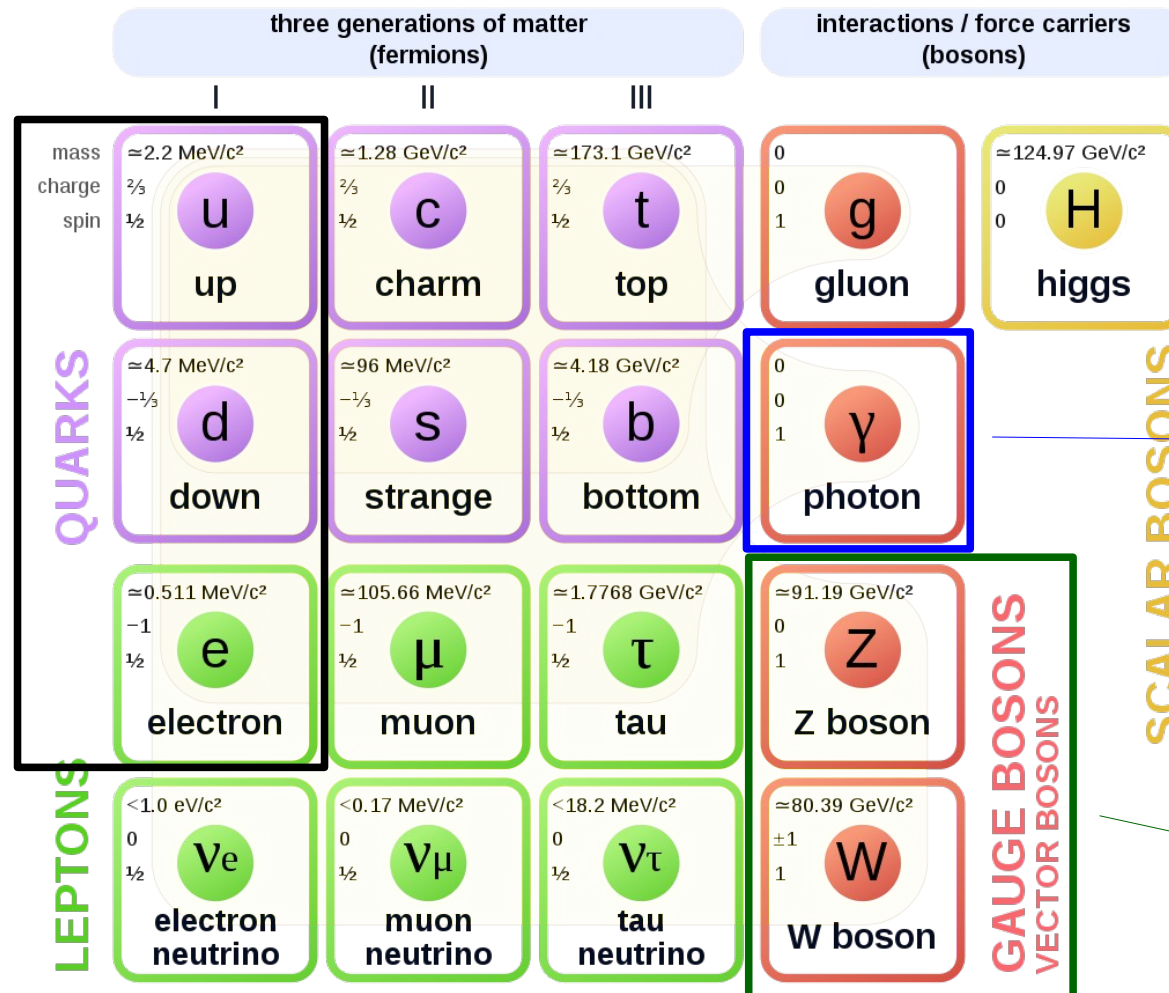
$$\Delta\alpha/\alpha = 0.15 \cdot 10^{-9}$$

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Standard Model of Elementary Particles



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Weak interaction
(β decays, sun, ...)

$$G_F \approx 1.17 \cdot 10^{-5} / \text{GeV}^2$$

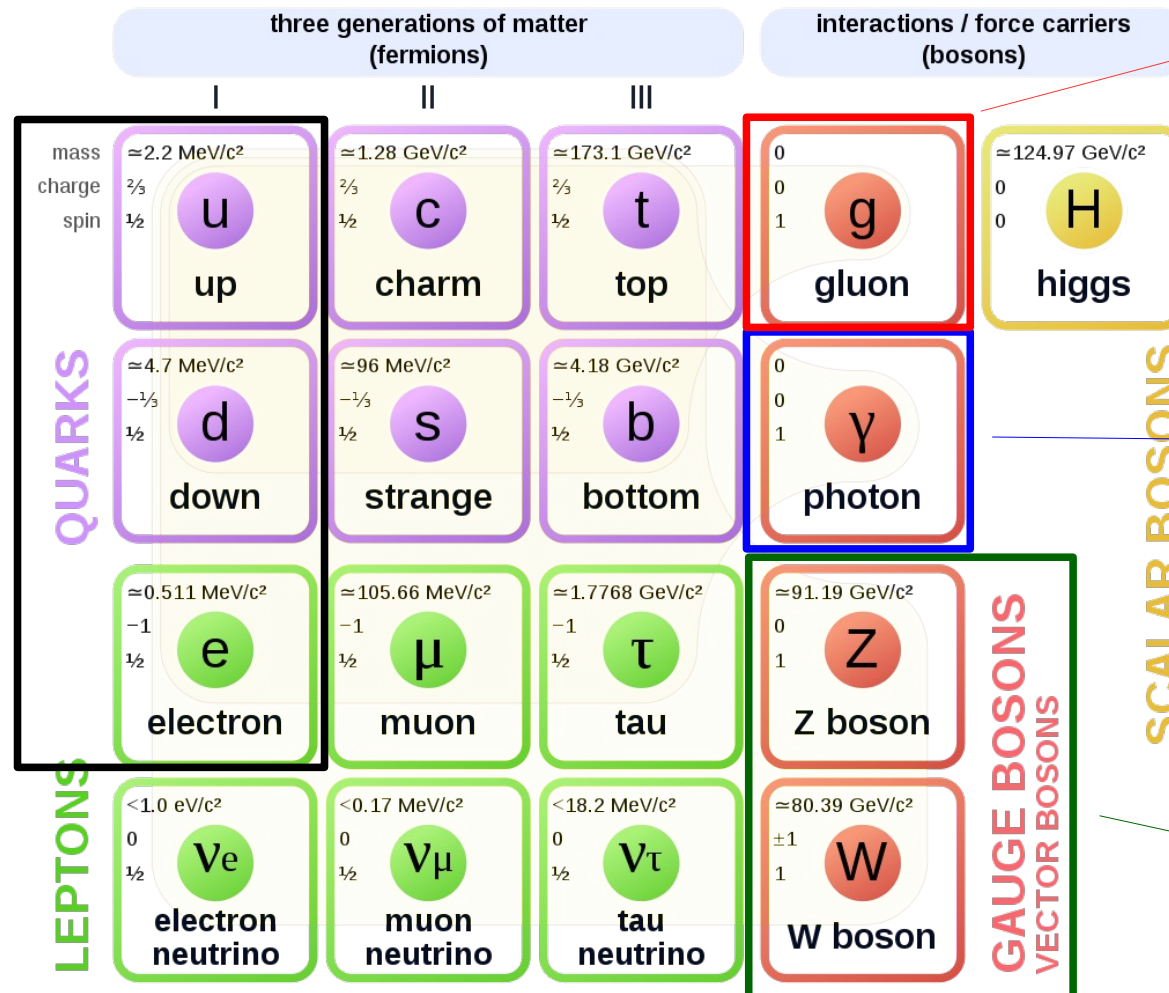
$$\Delta G_F / G_F = 0.51 \cdot 10^{-6}$$

... and three fundamental interactions.
(no gravity)

Cush, Wikipedia.



Standard Model of Elementary Particles



Cush, Wikipedia.

Strong interaction
(nuclear forces, ...)

$$\alpha_s \approx 0.118$$

$$\Delta\alpha_s/\alpha_s = 0.76 \cdot 10^{-2}$$

Electromagnetic interaction
(magnets, electricity, ...)

$$\alpha \approx 1/137$$

$$\Delta\alpha/\alpha = 0.15 \cdot 10^{-9}$$

Weak interaction
(β decays, sun, ...)

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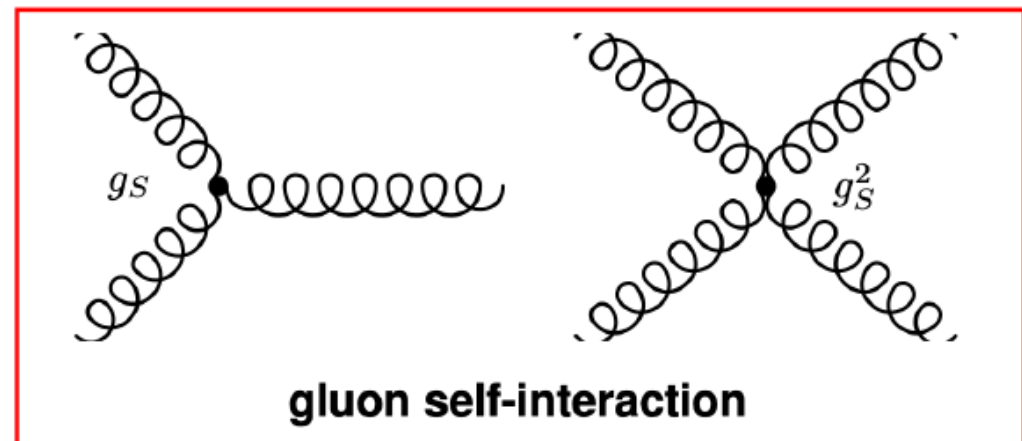
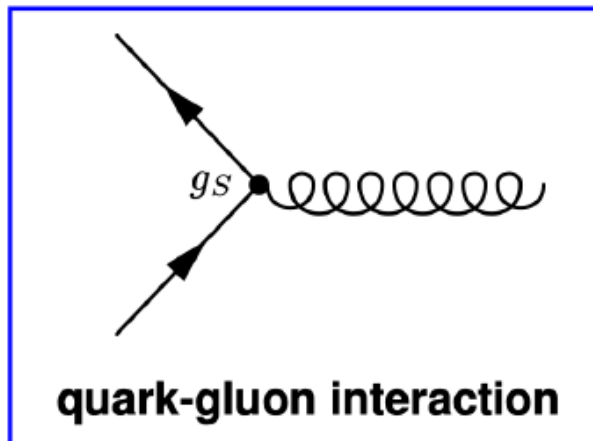
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Invariance under local $SU(3)_c$ transformations

- Three color charges $a = 1, 2, 3 \rightarrow$ **Red**, **Green**, **Blue**
(as analogue to electric charge in QED)
- Eight vector fields (gluons) $\mathcal{A}_\mu^A$ carry color charge and color anti-charge
- The gluons are massless
 → exact symmetry
 → in principle infinite range of strong force

$$\mathcal{G}_{\mu\nu}^A = \partial_\mu \mathcal{A}_\nu^A - \partial_\nu \mathcal{A}_\mu^A - g_s f^{ABC} \mathcal{A}_\mu^B \mathcal{A}_\nu^C$$

- Non-zero commutator leads to gluon self-interactions via triple and quartic gauge couplings





■ In (renormalisable) QFT the beta function encodes the dependence of the coupling parameter g on the energy (or distance) scale μ :

$$\alpha_i := \frac{g_i^2}{4\pi}$$

$$\beta(g) = \frac{\partial g}{\partial \log(\mu^2)}$$



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- **Beta function of QED (1-loop):** $\beta(\alpha) = \frac{1}{3\pi}\alpha^2$

- The coupling increases with energy scale
- The coupling decreases with larger distances
 - ➔ Infinite range, Coulomb potential: $V(r) \propto \frac{1}{r}$



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- **Beta function of QCD (1-loop):** $\beta(\alpha_s) = - \left(\frac{11N_C - 2N_f}{12\pi} \right) \alpha_s^2$

- ➔ The coupling decreases with energy scale, if $N_C = 3, \quad N_f \leq 16$
 - ➔ **Asymptotic freedom**
- ➔ The coupling increases with larger distances
 - ➔ **Confinement**, string potential: $V(r) \approx \sigma \cdot r$ with tension $\sigma \approx 1 \text{ GeV/fm}$



Nobel prize 2004

Theory:

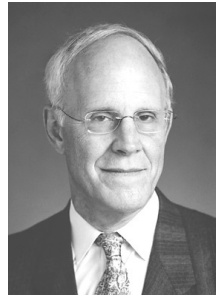
- Renormalisation group equation (RGE)
- Solution of 1-loop equation
- **Running coupling constant**

$$\alpha_s(Q^2) = \frac{\alpha_s(\mu^2)}{1 + \alpha_s(\mu^2)\beta_0 \ln\left(\frac{Q^2}{\mu^2}\right)}$$
$$\alpha_s(Q^2) = \frac{1}{\beta_0 \ln\left(\frac{Q^2}{\Lambda^2}\right)}$$

Towards small distances, $Q^2 \rightarrow \infty$

- “Strong” coupling becomes weak
- Perturbative methods usable
- Asymptotic freedom

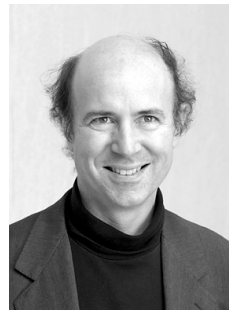
Physik Journal 3 (2004) Nr. 12



D. Gross



D. Politzer



F. Wilczek

nobelprize.org



Nobel prize 2004

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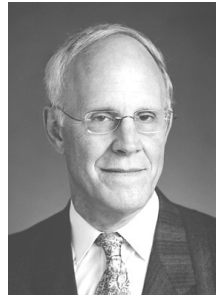
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Towards large distances, $Q^2 \rightarrow 0$

- Strong coupling, confinement
- Perturbative methods **not** usable for $Q^2 \rightarrow \Lambda^2$
- Lattice gauge theory

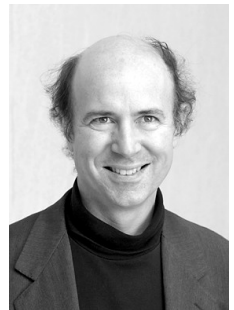
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Running coupling constant

$$\alpha_s(Q^2) = \frac{1}{\beta_0 \ln \left(\frac{Q^2}{\Lambda^2} \right)}$$

with Λ typically $\approx 200 - 300$ MeV

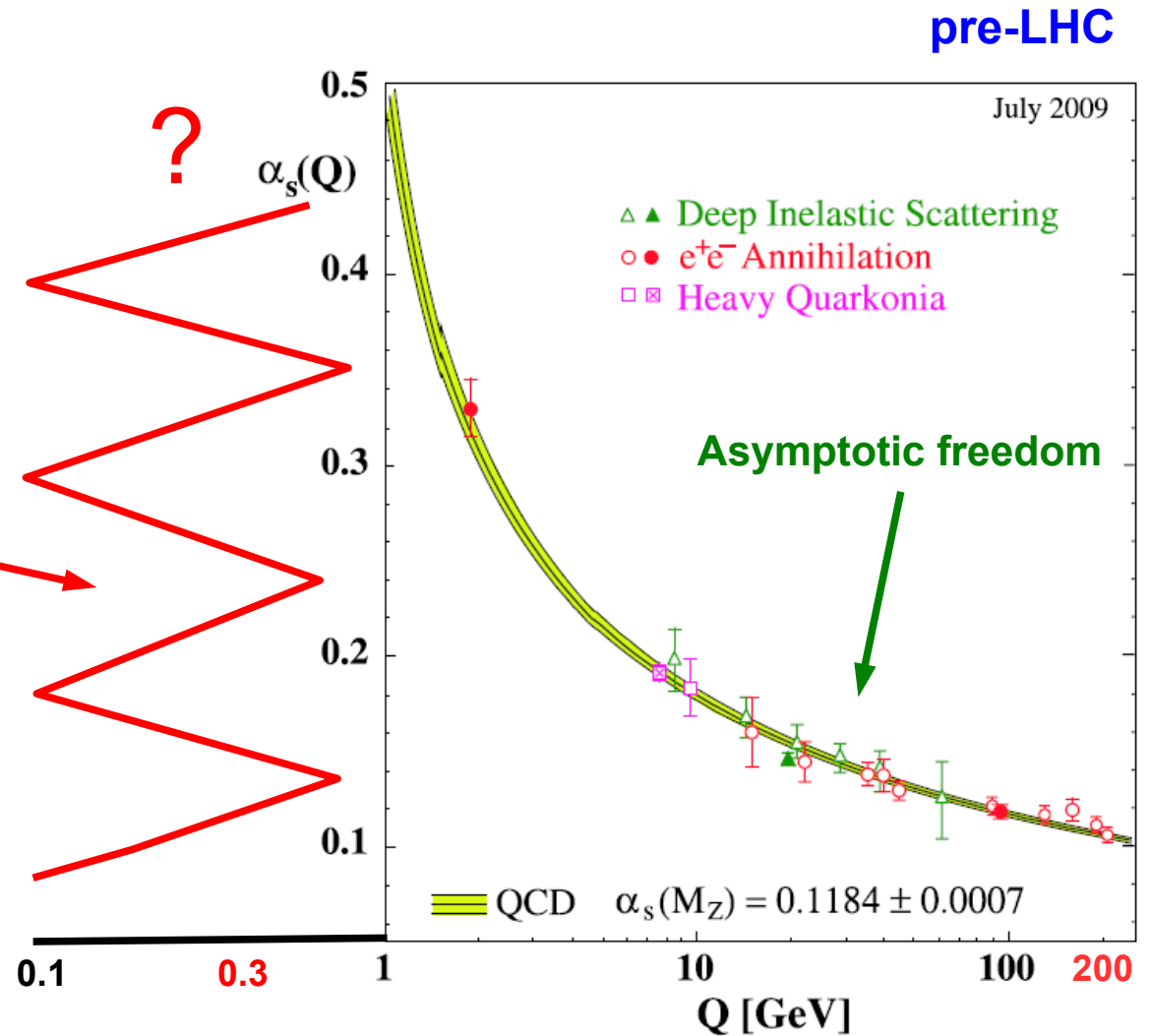
Non-perturbative regime

QCD potential grows linearly
with larger distances:

$$V = \sigma \cdot r \approx 1 \text{ GeV/fm} \cdot r$$

→ No free quarks (or gluons)

→ Confinement



S. Bethke, EPJC 64 (2009).



Status of α_s in PDG review

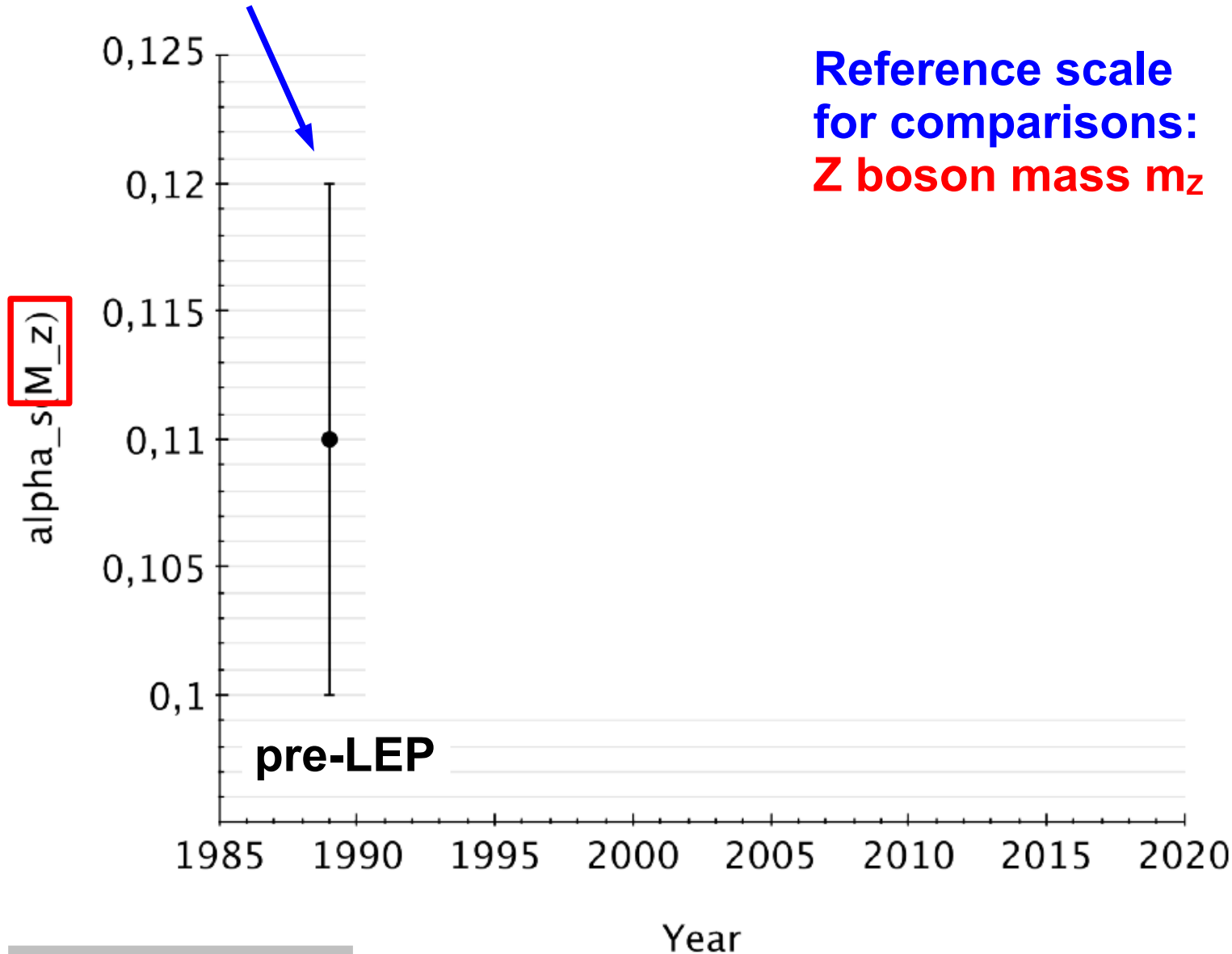


Particle Data Group
<https://pdg.lbl.gov>



$\alpha_s(m_Z)$ average versus time

1st estimate from G. Altarelli

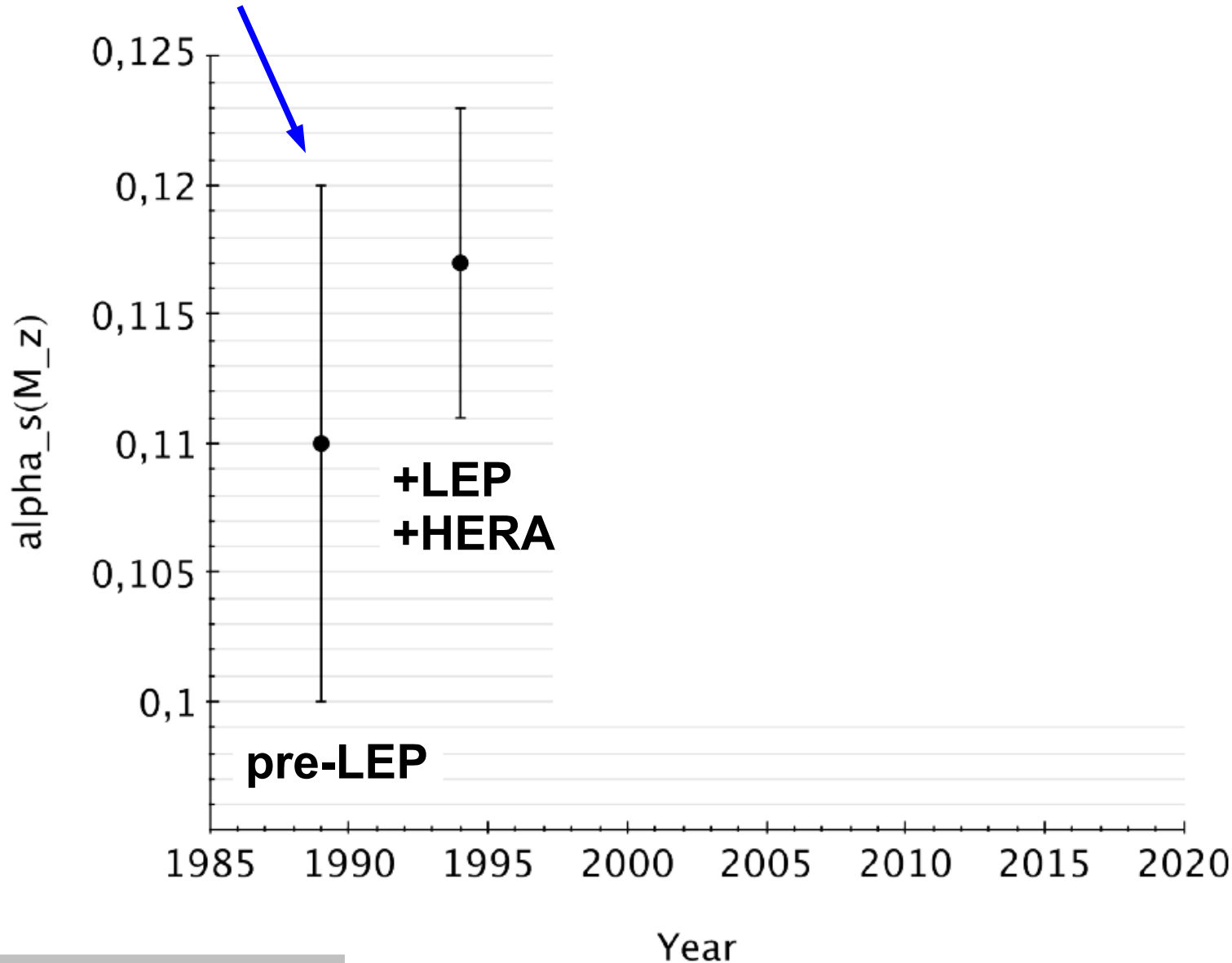


S. Bethke, arXiv:1907.01435.



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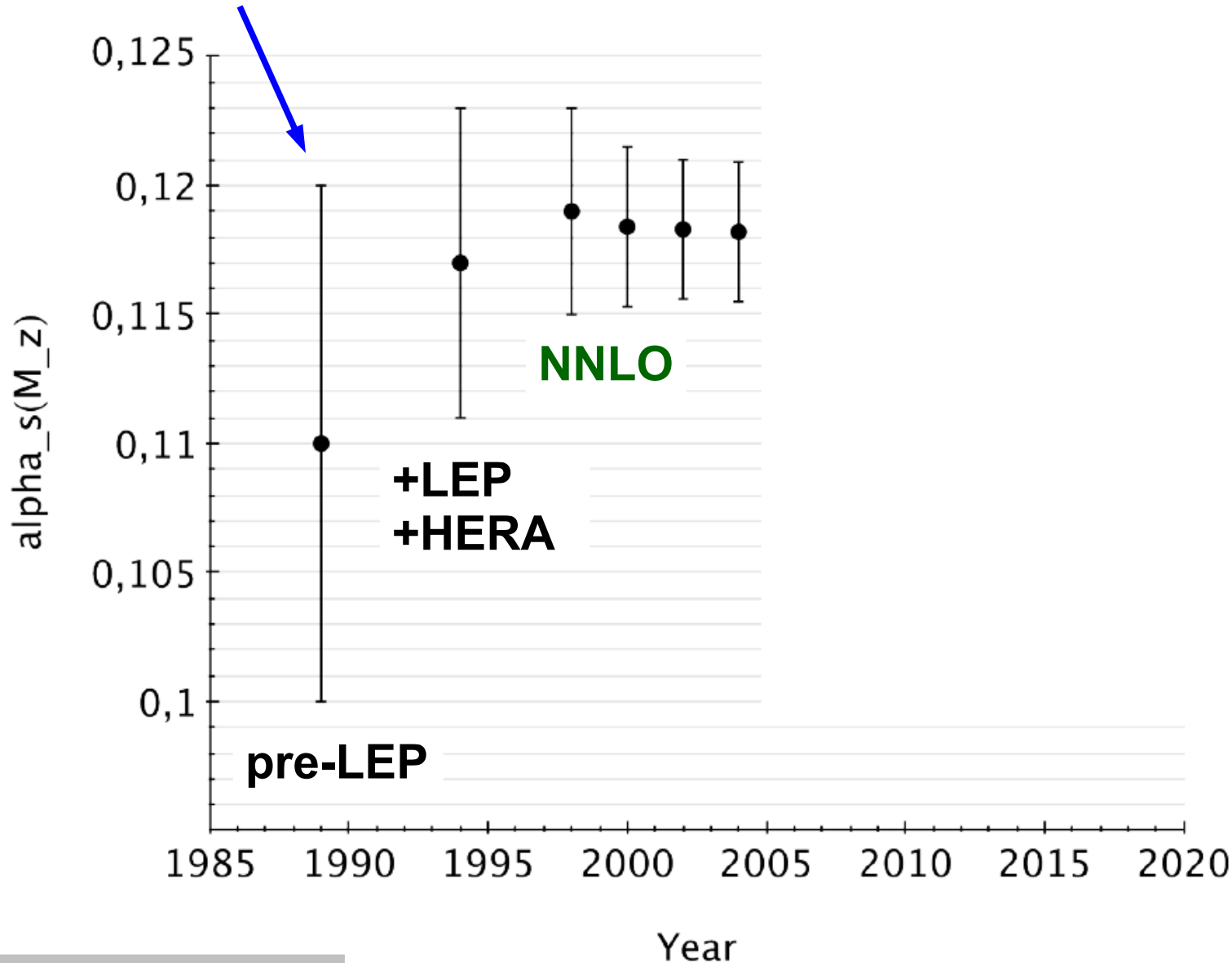


S. Bethke, [arXiv:1907.01435](https://arxiv.org/abs/1907.01435).



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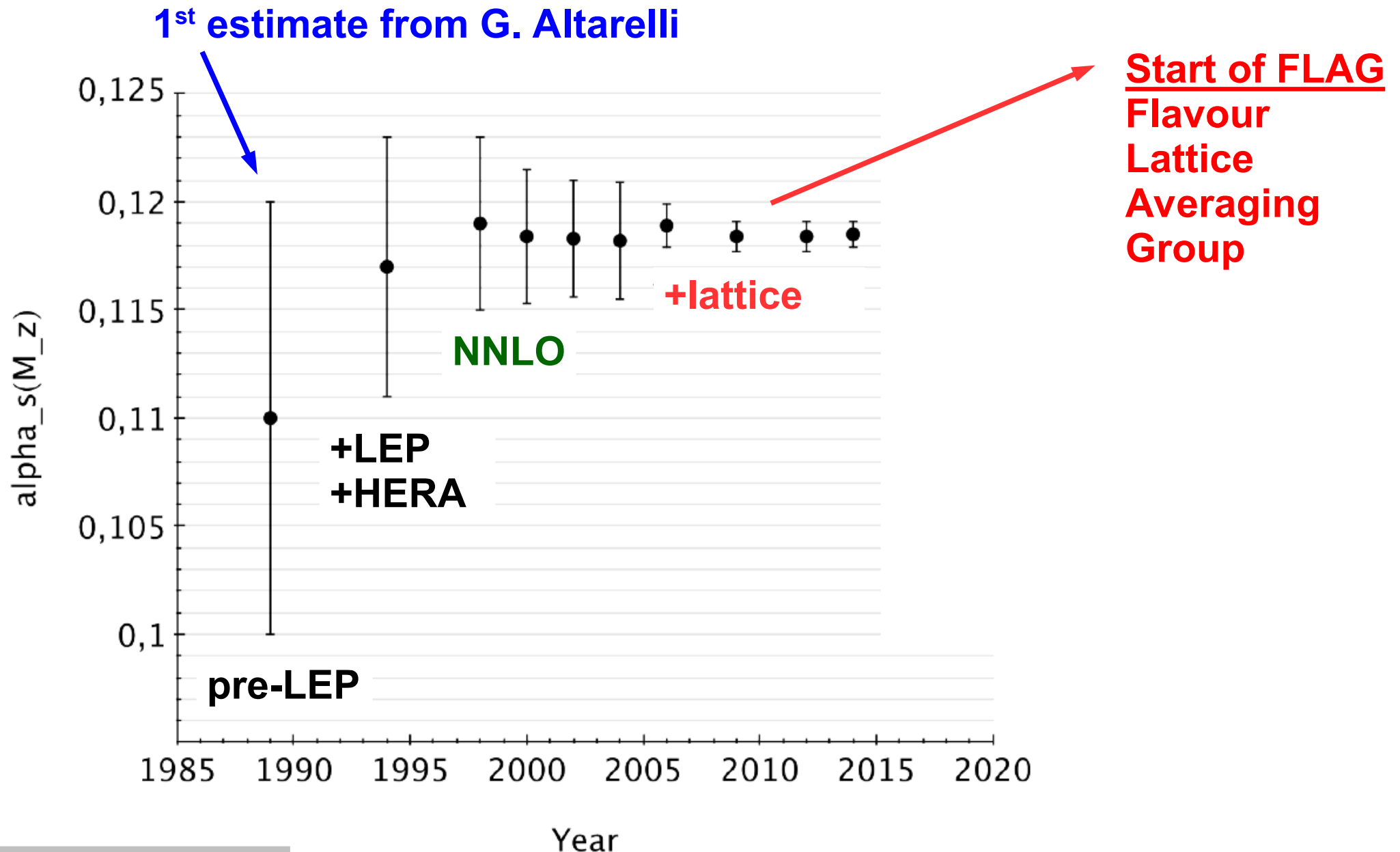
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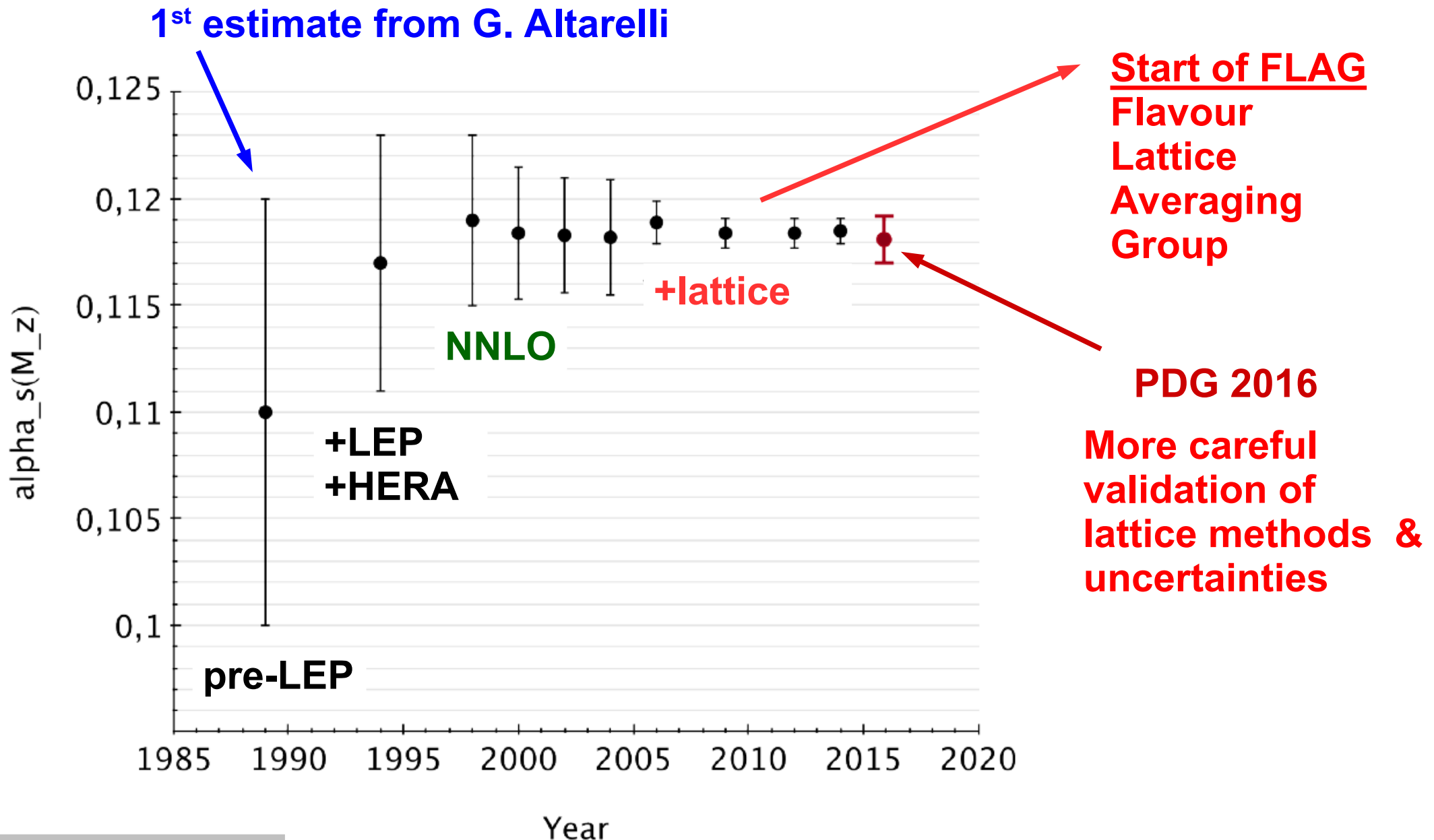
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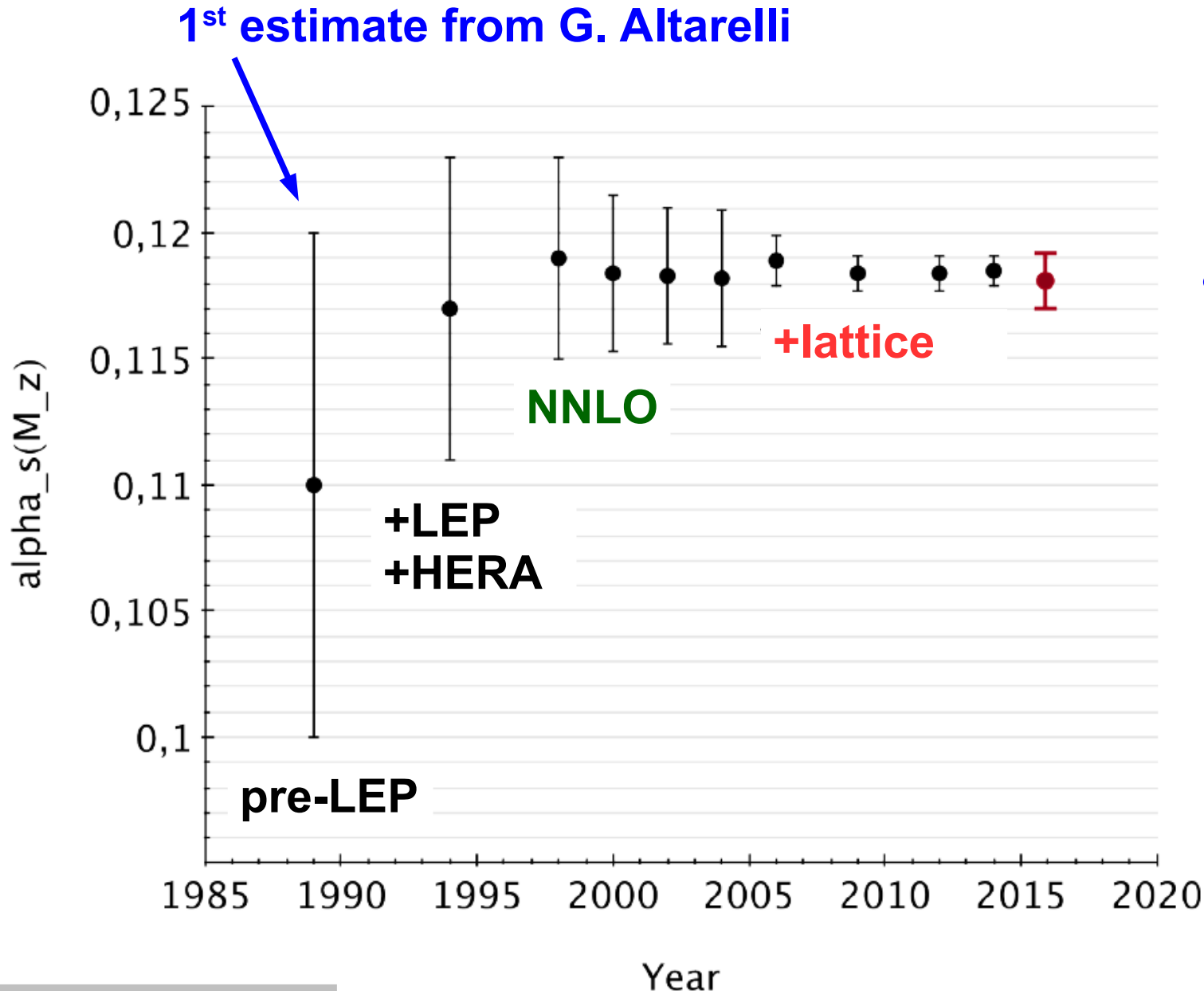
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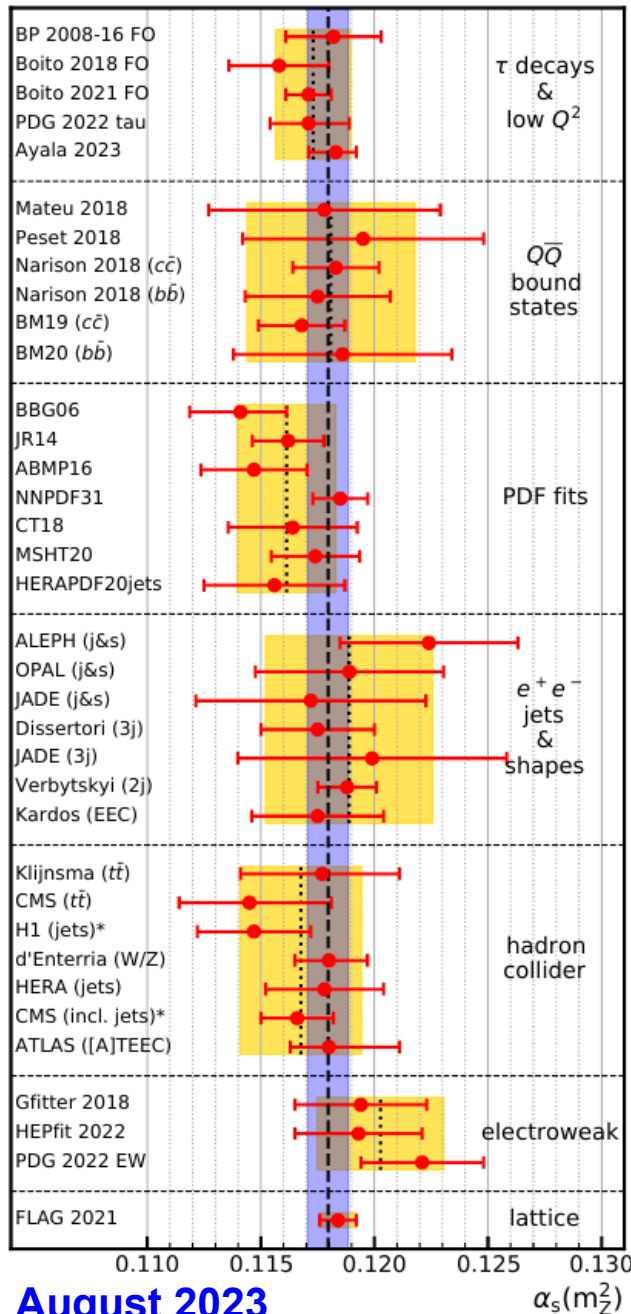
PDG 2023

Significant impact
on uncertainty
from (PDF + α_s) on
Higgs x sections:

Particularly:
tTH & gg-Fusion:
7-13%



PDG α_s averaging in 6 groups



τ hadronic decay widths & spectral functions

heavy quarkonia decays

global fits of proton structure & α_s

event shapes & jet rates in e^+e^-

observables from hh collisions & DIS

electroweak fits

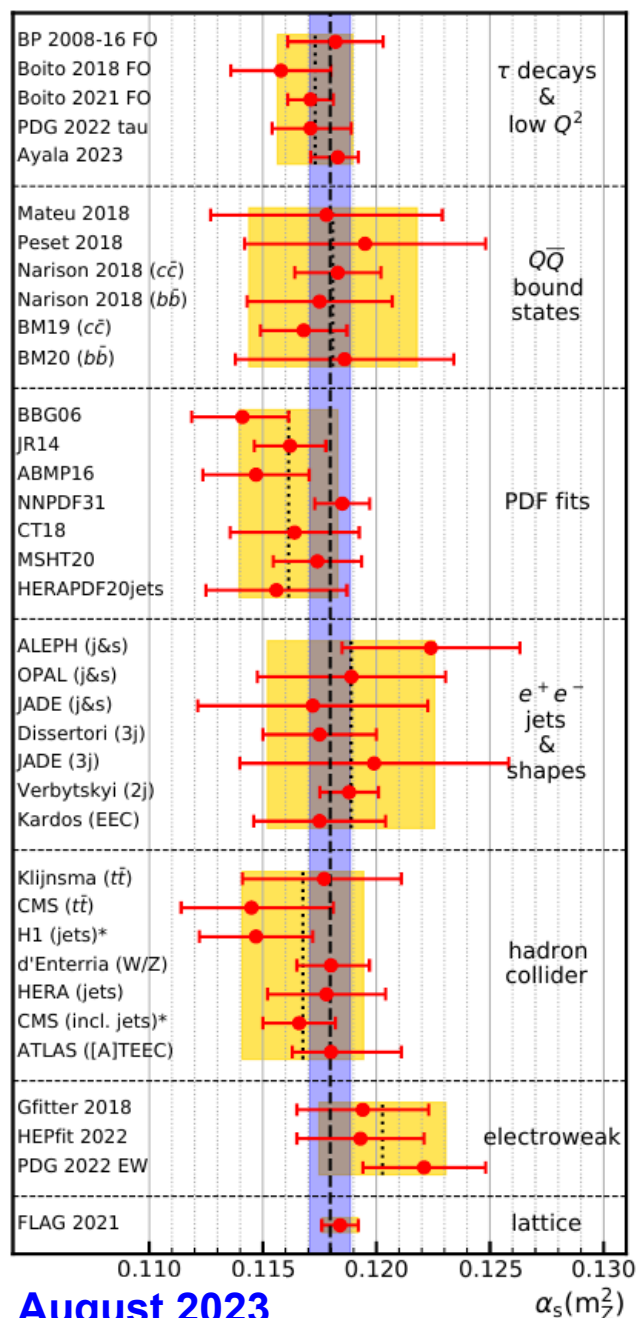
FLAG average from lattice calculations

August 2023

PDG, PRD (2024) 110, 3, 030001.



PDG α_s averaging in 6 groups



τ hadronic decay widths & spectral functions

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partially also
input here

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averages per sub-field	unweighted
τ decays & low Q^2	0.1173 ± 0.0017
$Q\bar{Q}$ bound states	0.1181 ± 0.0037
PDF fits	0.1161 ± 0.0022
e^+e^- jets & shapes	0.1189 ± 0.0037
hadron colliders	0.1168 ± 0.0027
electroweak	0.1203 ± 0.0028
PDG 2023 (without lattice)	0.1175 ± 0.0010

Final average including lattice (FLAG2021):

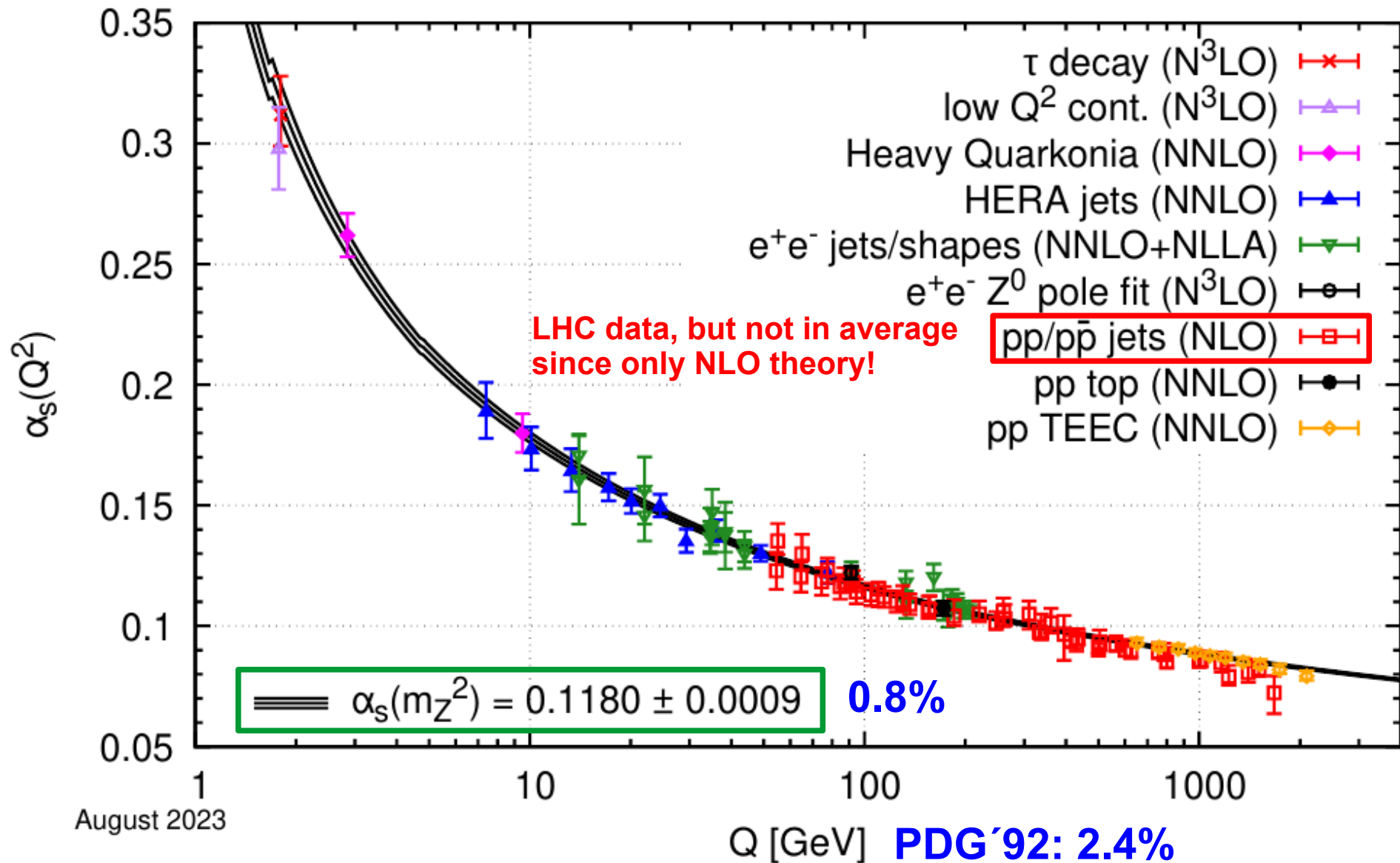
$$\alpha_s(m_Z^2) = 0.1180 \pm 0.0009$$

rel. uncertainty: 0.8%

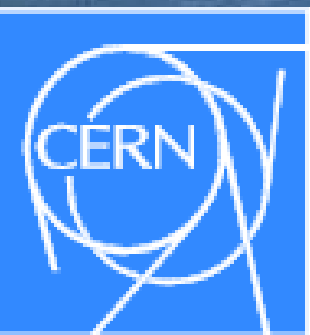
PDG, PRD (2024) 110, 3, 030001.



PDG 2023 α_s running



PDG, PRD (2024) 110, 3, 030001.



New results from LHC

Lake
Geneva



CMS

Airport



LHCb

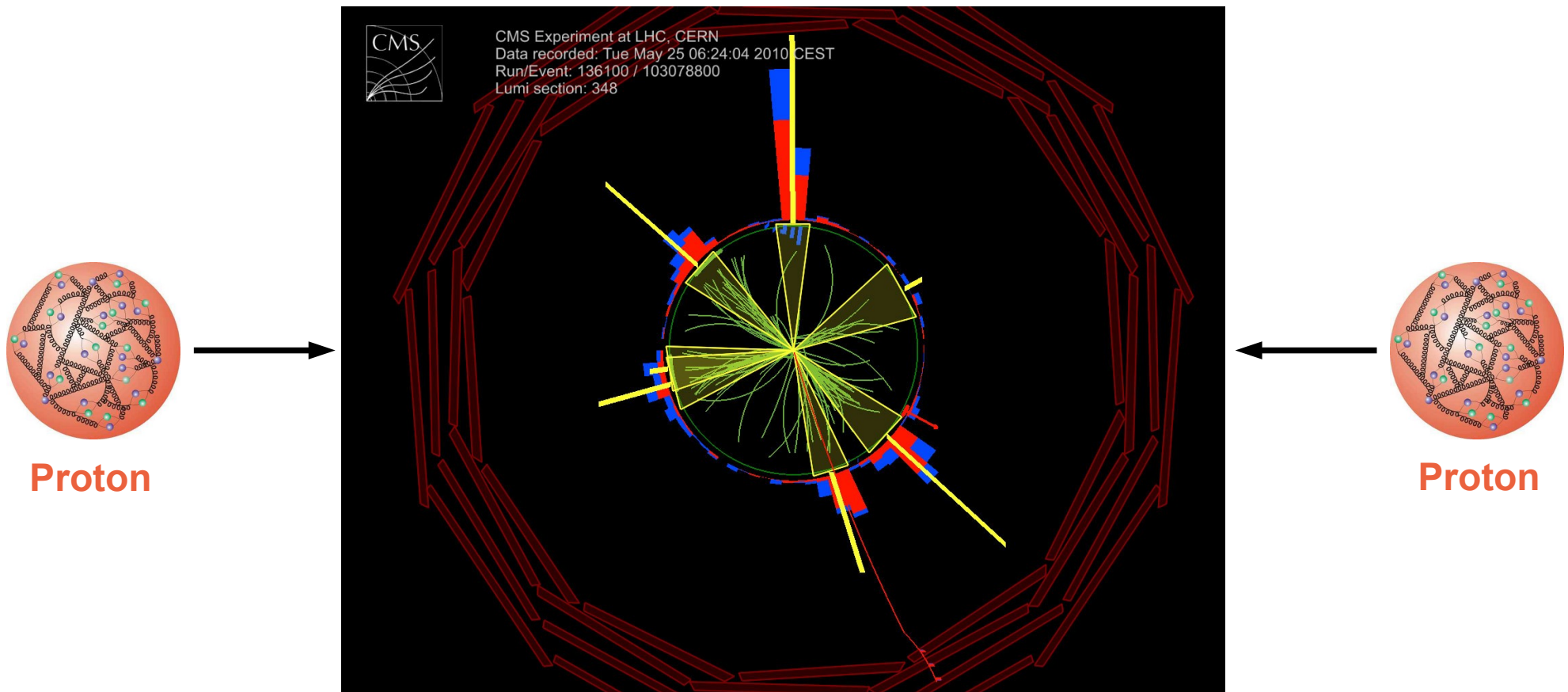


ALICE



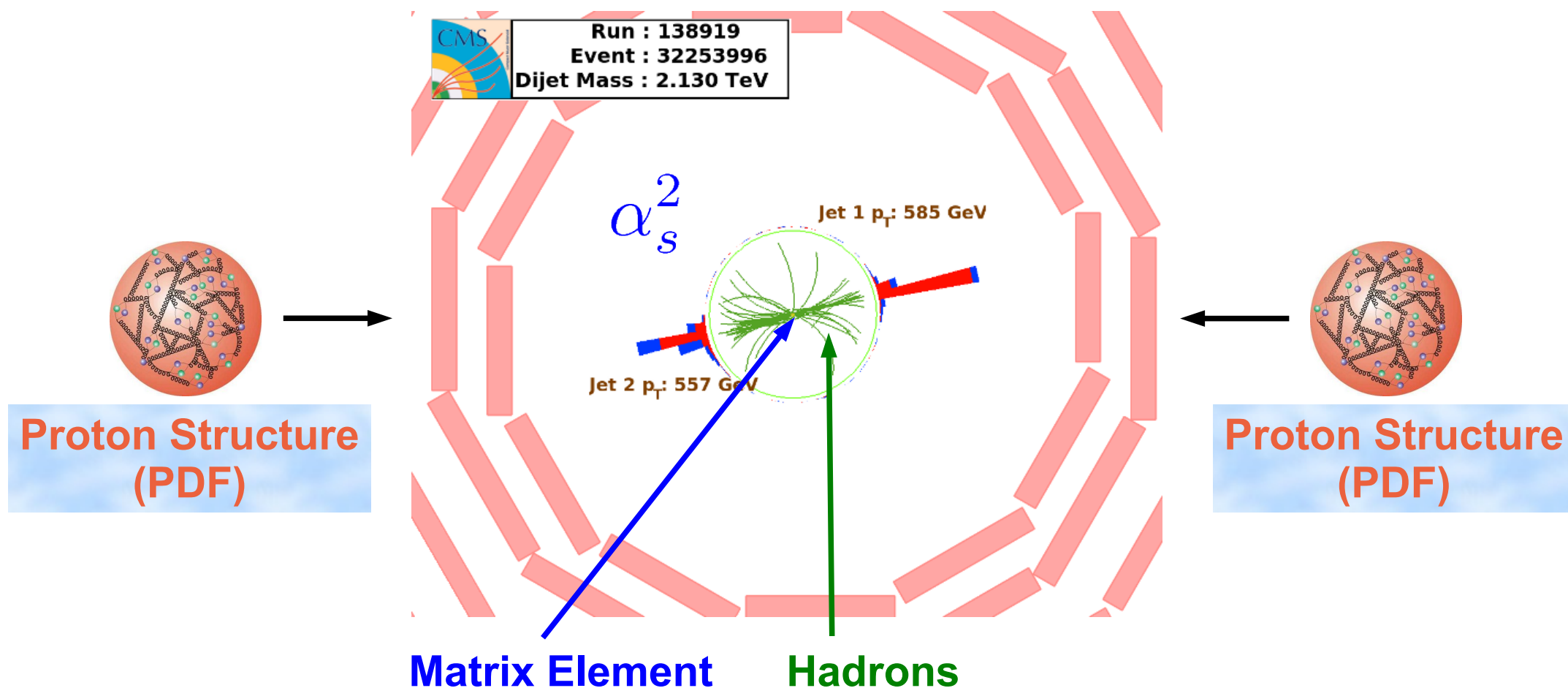
ATLAS

Large transverse momenta



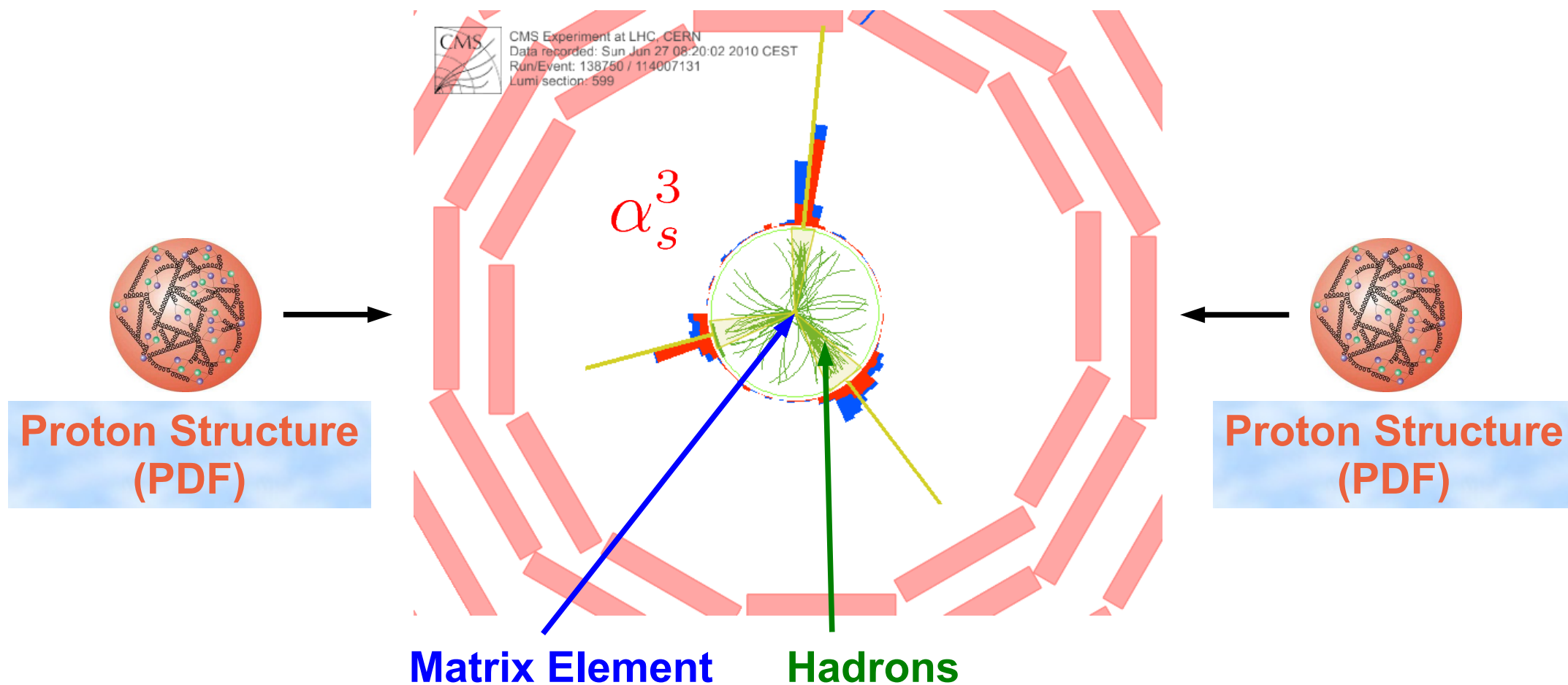
Abundant production of jets:

- Highest reach ever in energy scale Q to determine the strong coupling
- Learn about hard QCD, the proton structure, non-perturbative effects, and electroweak effects at high Q



Abundant production of jets:

➡ Extract $\alpha_s(m_Z)$, the least precisely known fundamental constant!





Jet cross sections $\sim \alpha_s^{2+n}$

- Counting jets or jet events in bins of e.g. momentum and rapidity
- Useful for i.a.:
 - ➔ Determination of $\alpha_s(m_Z)$
 - ➔ Test of running of $\alpha_s(Q)$
 - ➔ Multi-parameter fit of $\alpha_s(m_Z)$ & PDFs
 - ➔ Multi-parameter fit including EFT parameters



Jet cross sections $\sim \alpha_s^{2+n}$

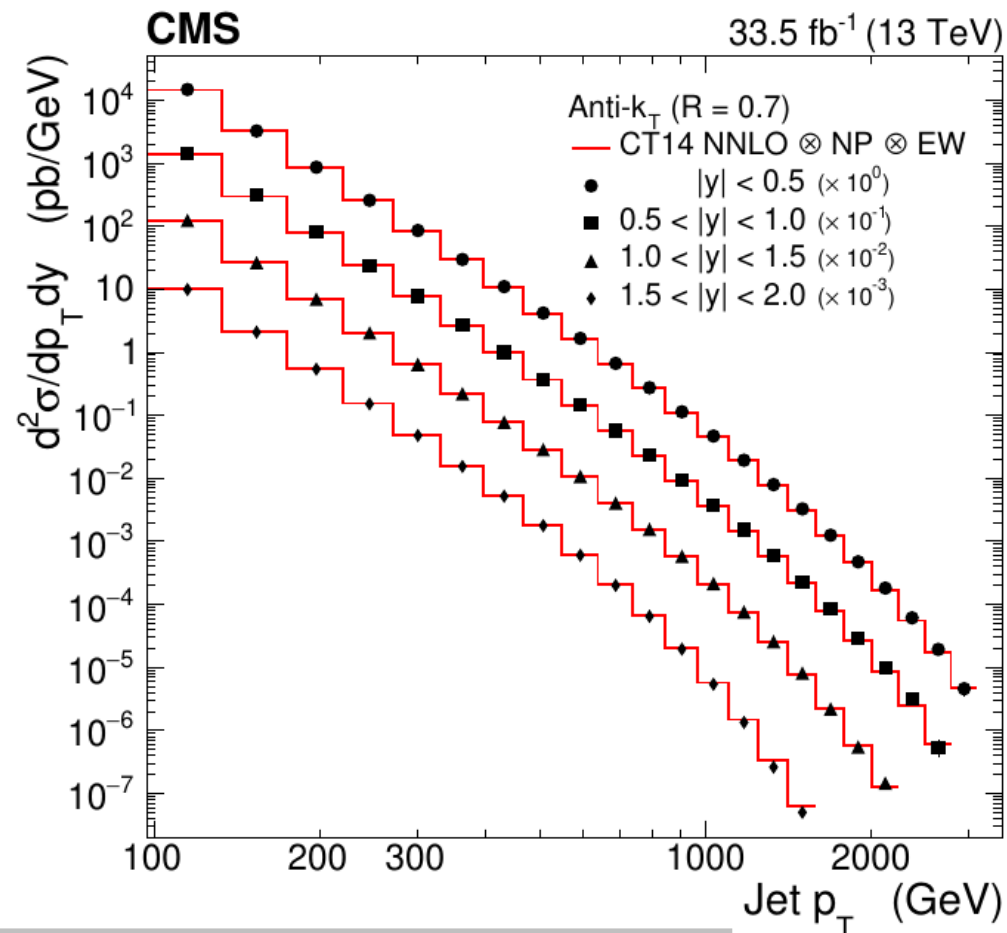
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 - ➔ Multi-parameter fit including EFT parameters
- Subject to many/all systematic uncertainties:
 - ➔ Jet energy calibration (JEC) & resolution (JER)
 - ➔ Luminosity
 - ➔ Missing higher orders
 - ... to name a few



Jet counting in bins of jet transverse momentum and rapidity
Comparison of unfolded measurement with theory:

NNLO x nonperturbative x electroweak correction

(NNLO in leading-color (LC) approximation; subleading effects small)



CMS, JHEP02 (2022) 142 & JHEP12 (2022) 035.

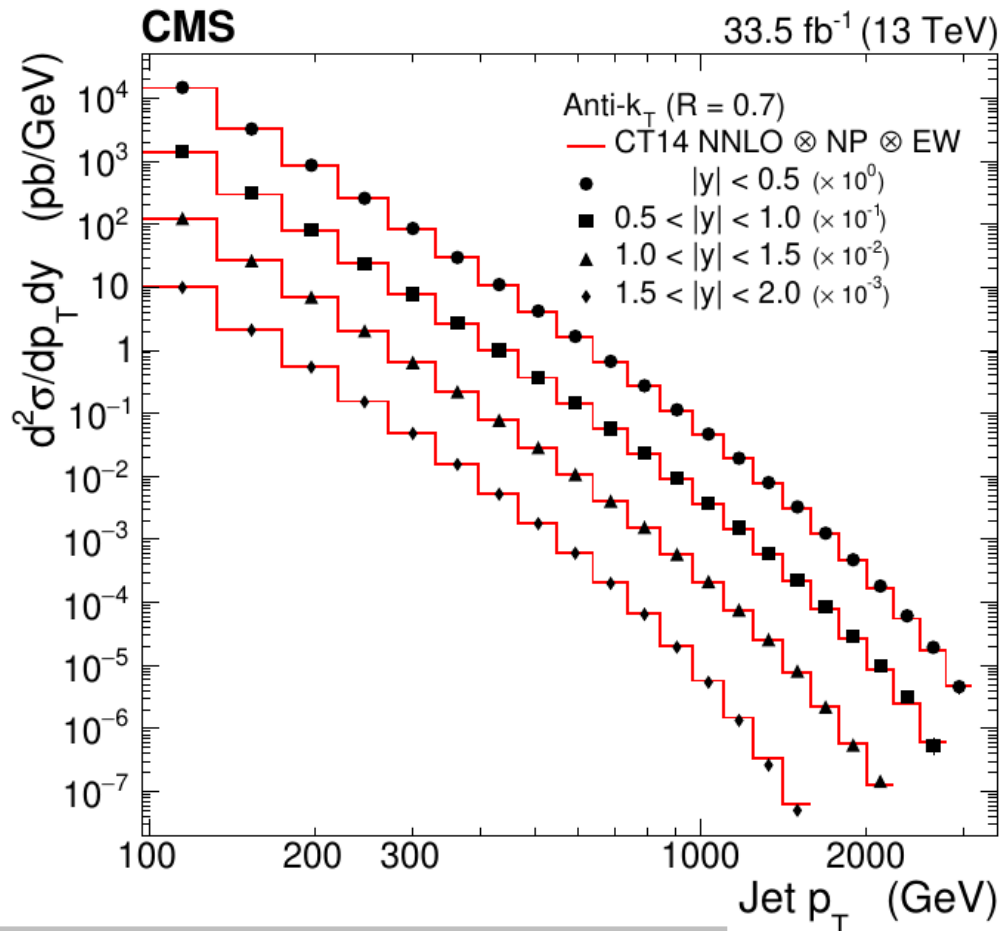


Inclusive jets: α_s & PDFs

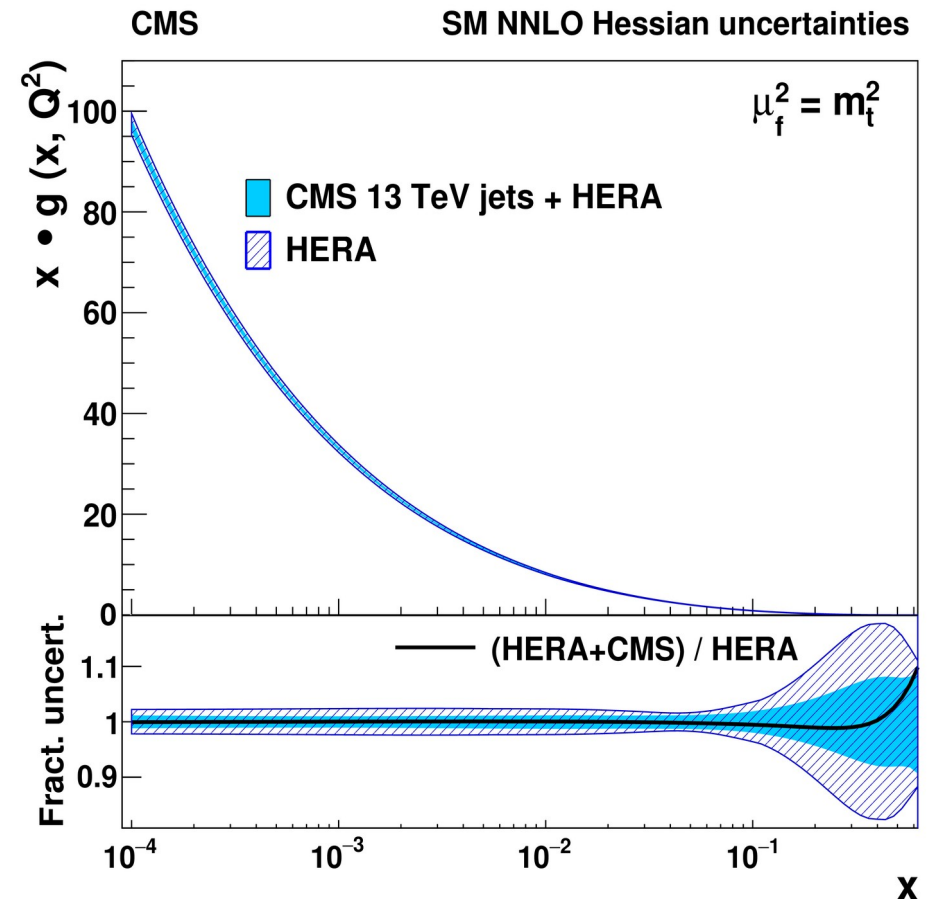
First determination of $\alpha_s(m_Z)$
from jets at NNLO-LC:

$$\alpha_s(m_Z^2) = 0.1166 \pm 0.0016(\text{fitall}) \pm 0.0004(\text{scl})$$

Simultaneous fit with PDFs
→ reduced uncertainties of gluon



CMS, JHEP02 (2022) 142 & JHEP12 (2022) 035.



Jet event counting in bins of dijet mass or average transverse momentum (x) and

rapidity separation

$$y^* = \frac{1}{2} |y_1 - y_2|$$

boost of dijet system

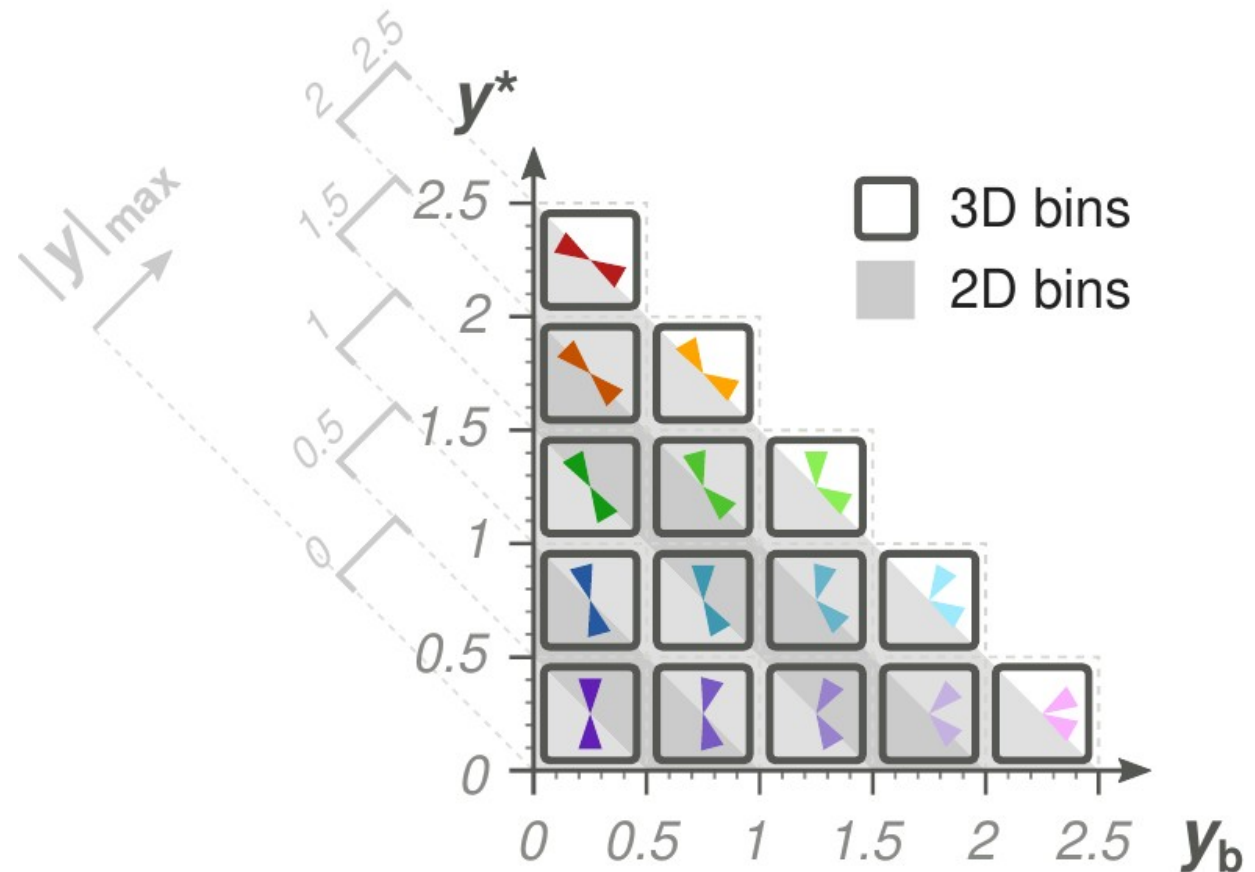
$$y_b = \frac{1}{2} |y_1 + y_2|$$

$$\frac{d^3\sigma}{dy^* dy_b dx} = \frac{1}{\varepsilon \mathcal{L}_{\text{int}}} \frac{N}{\Delta y^* \Delta y_b \Delta x}.$$

Alternativ 2D binning in maximum rapidity of $|y_1|, |y_2|$

CMS, EPJC 85 (2025) 72.

Illustration of dijet event topologies



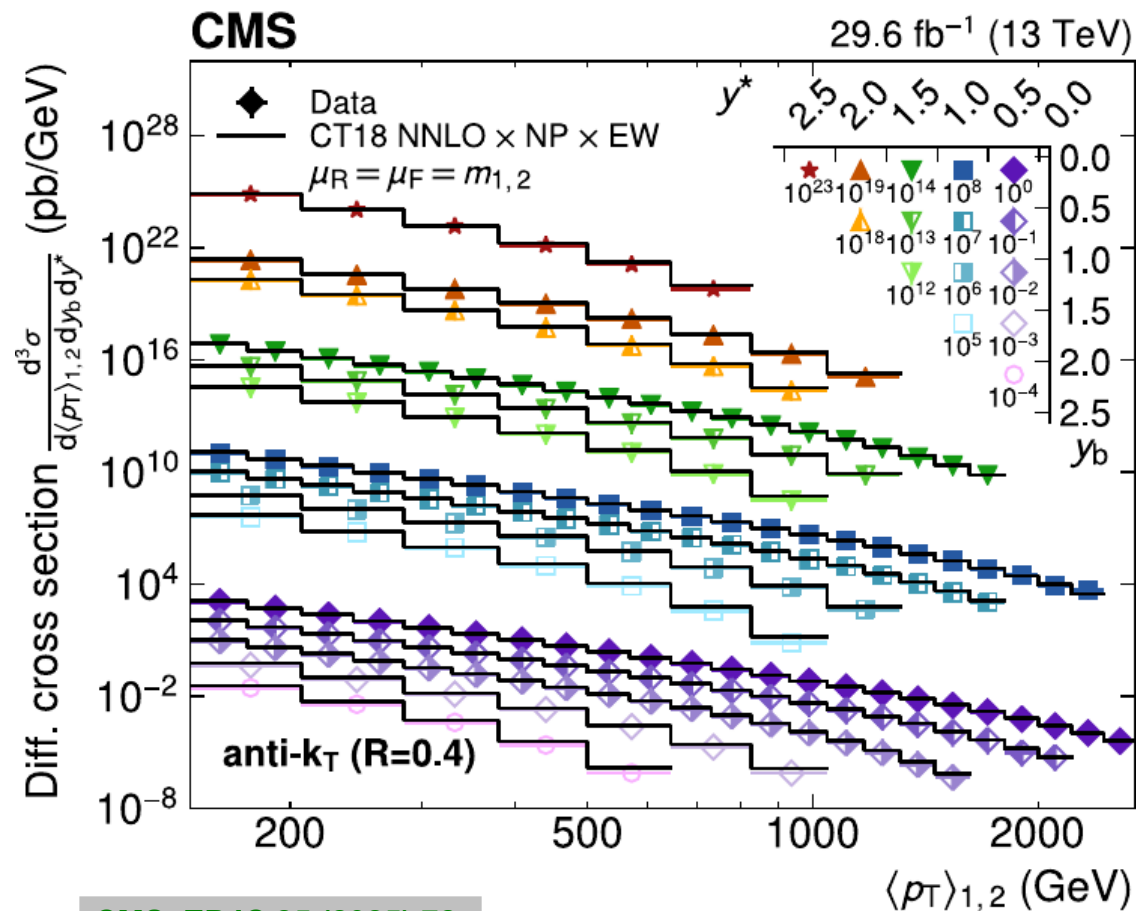
Jet event counting in bins of x, y^*, y_b

Comparison of unfolded measurement with theory:

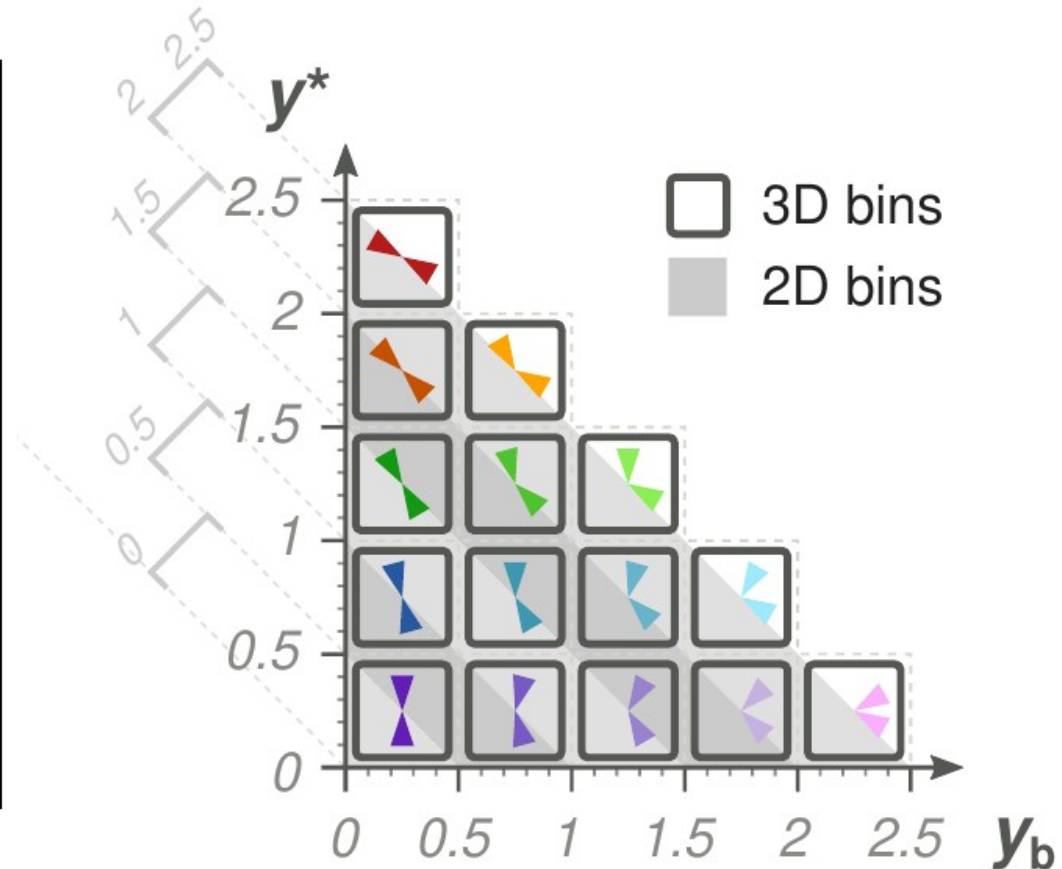
NNLO x nonperturbative x electroweak correction

Illustration of dijet event topologies

(Impact of subleading-color corrections up to few %)



CMS, EPJC 85 (2025) 72.



Determine $\alpha_s(m_Z)$ + PDFs

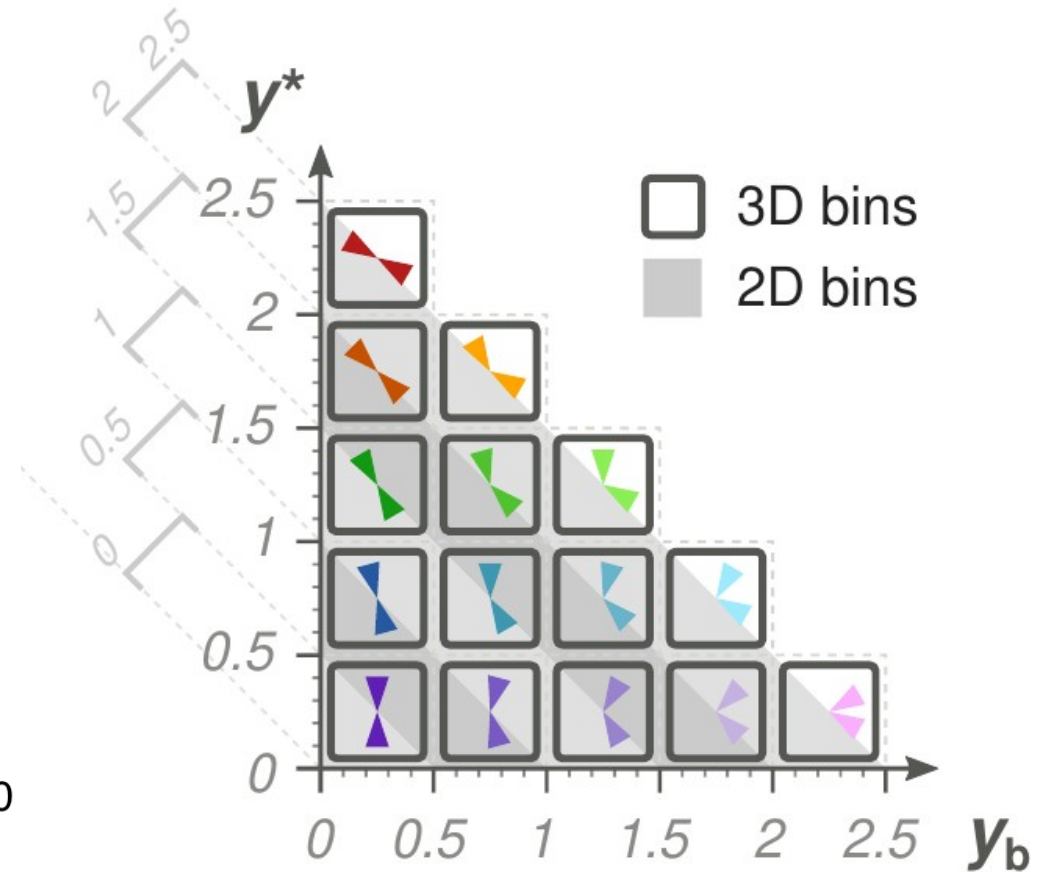
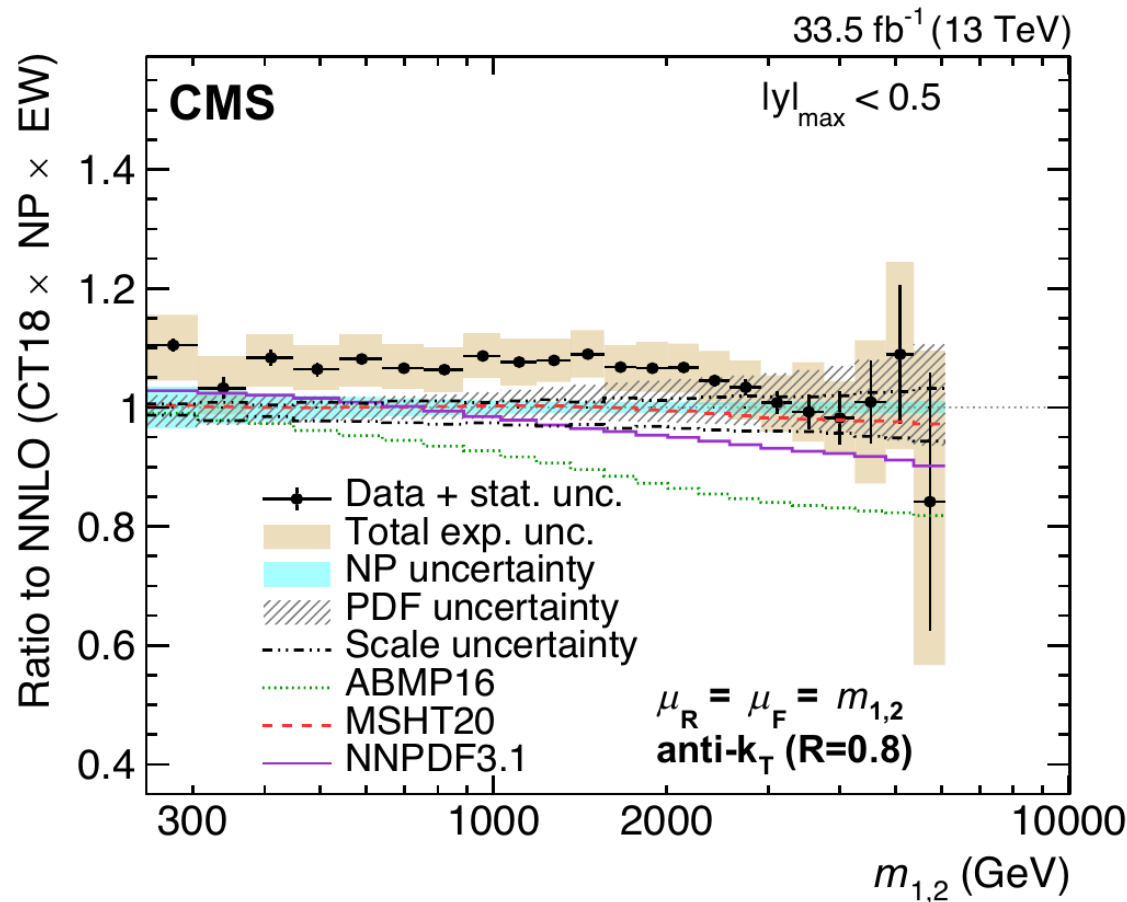
Example of ratio data / theory

2D

$$\alpha_s(m_Z) = 0.1179 \pm 0.0017(\text{fitall}) \pm 0.0008(\text{scl})$$

3D

$$\alpha_s(m_Z) = 0.1181 \pm 0.0019(\text{fitall}) \pm 0.0009(\text{scl})$$





Combining jet datasets $\rightarrow \alpha_s(Q)$

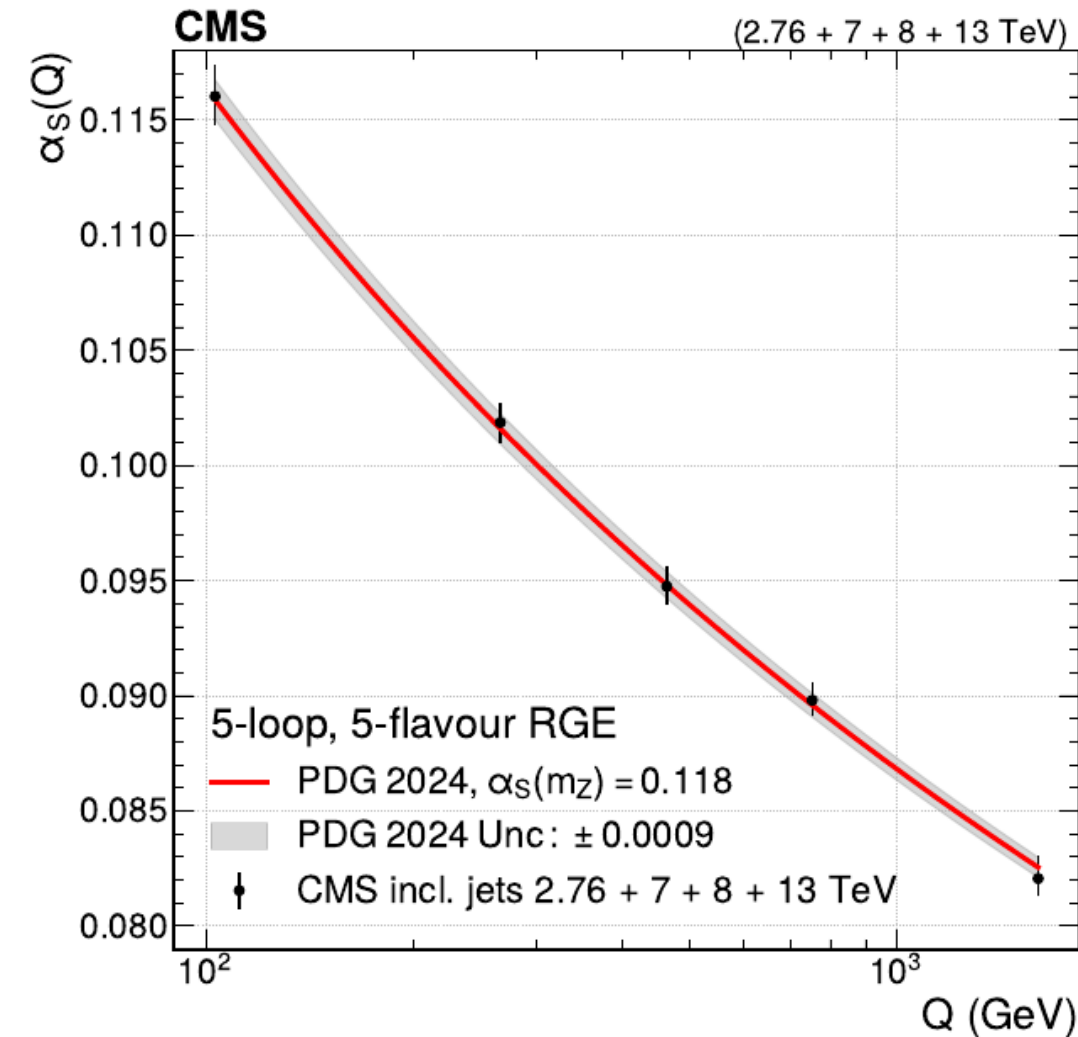
- Up to now: Fit of whole dataset $\rightarrow \alpha_s(m_Z)$ & PDFs (running used implicitly)
- New QCD analysis from CMS
 - Combination of multiple datasets ($\rightarrow$ table)
 - Subdivide into ranges of jet p_T
 - Multi-parameter fit of $\alpha_s(m_Z)$ & PDFs in each range separately
- Uncertainty correlations vs. E_{cms} studied and varied
 - Some correlation in JEC among cms energies, insignificant in JER
 - MHOU (scale variation) fully correlated

$\sqrt{s}$ [TeV]	$\mathcal{L}$ [fb^{-1}]	N_p	p_T [GeV]	$ y $
2.76	0.0054	80	74–592	0.0–3.0
7	5.0	130	114–2116	0.0–2.5
8	20	165	74–1784	0.0–3.0
13	33.5	78	97–3103	0.0–2.0

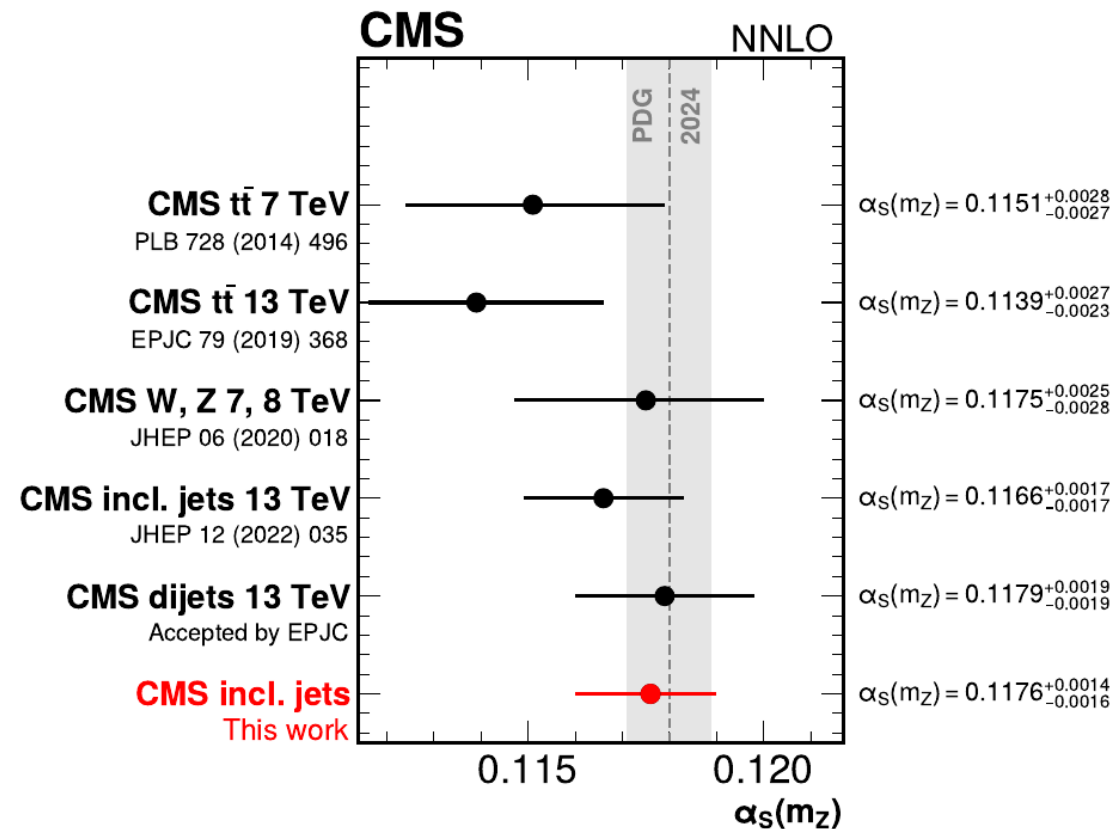
CMS, arXiv:2412.16665.



Running of $\alpha_s(Q)$ in five ranges of jet p_T



Comparison of $\alpha_s(m_Z)$ from all jet data to previous CMS results at NNLO





- New α_s extraction from all dijet data of ATLAS & CMS
 - Combination of multiple datasets ($\rightarrow$ table), in 2nd step also HERA dijets
 - Subdivide into ranges of relevant scale Q
 - Complete NNLO predictions
 - Simultaneous fit of one $\alpha_s(Q)$ per each range (with PDF variations as nuisance parameters at starting scale $\mu_0 = 90$ GeV)

Uncertainty correlations

- Experimental ones assumed negligible
- MHOU (scale variation) fully correlated

Data	$\sqrt{s}$ [TeV]	$d\sigma$	R	$\mathcal{L}$
ATLAS [10]	7	$\frac{d^2\sigma}{dm_{jj}dy^*}$	0.6	$4.5 \text{ fb}^{-1} \pm 1.8 \%$
CMS [12]	7	$\frac{d^2\sigma}{dm_{jj}dy_{\max}}$	0.7	$5.0 \text{ fb}^{-1} \pm 2.2 \%$
CMS [13]	8	$\frac{d^3\sigma}{d\langle p_T \rangle_{1,2} dy^* dy_b}$	0.7	$19.7 \text{ fb}^{-1} \pm 2.6 \%$
ATLAS [11]	13	$\frac{d^2\sigma}{dm_{jj}dy^*}$	0.4	$3.2 \text{ fb}^{-1} \pm 2.1 \%$
CMS [14]	13	$\frac{d^2\sigma}{dm_{jj}dy_{\max}}$	0.8	$33.5 \text{ fb}^{-1} \pm 1.2 \%$
CMS [14]	13	$\frac{d^3\sigma}{dm_{jj}dy^* dy_b}$	0.8	$29.6 \text{ fb}^{-1} \pm 1.2 \%$

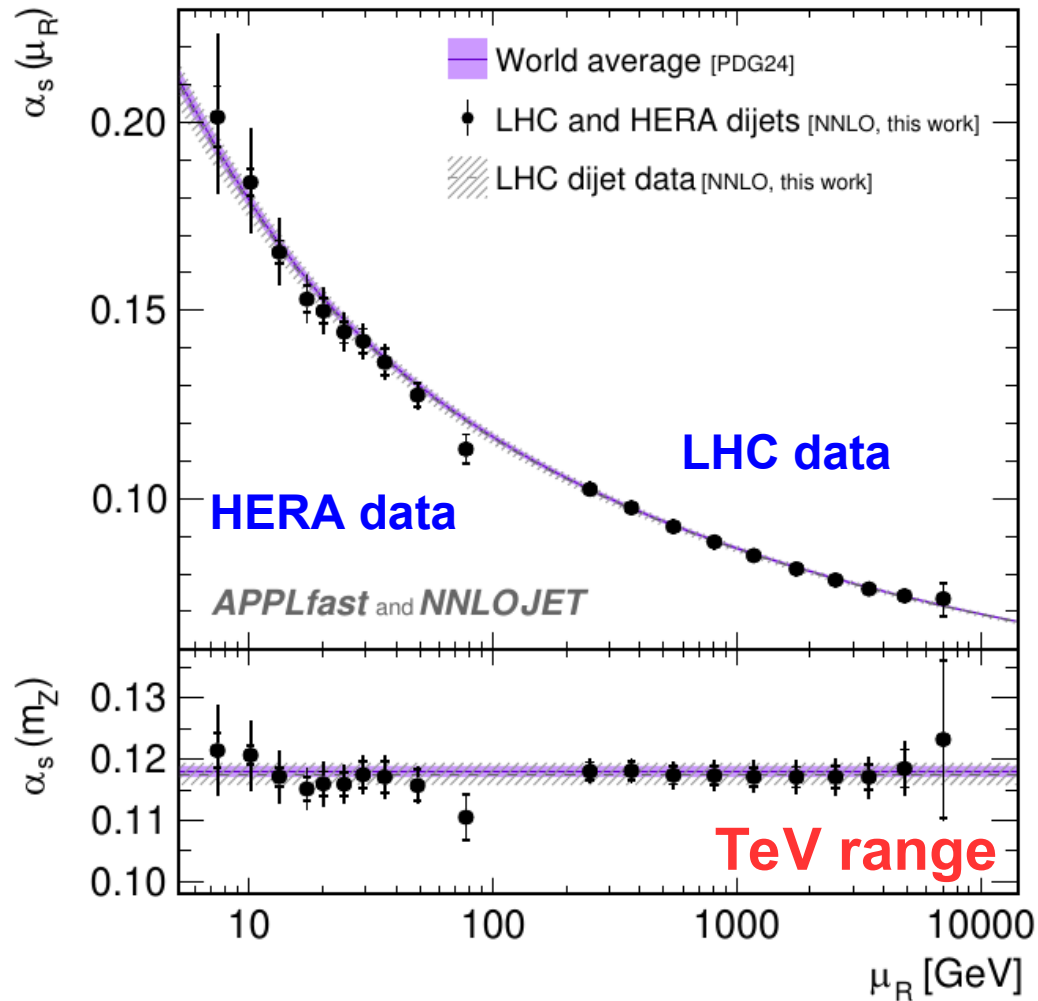
Ahmadova, KR, et al., arXiv:2412.21165.



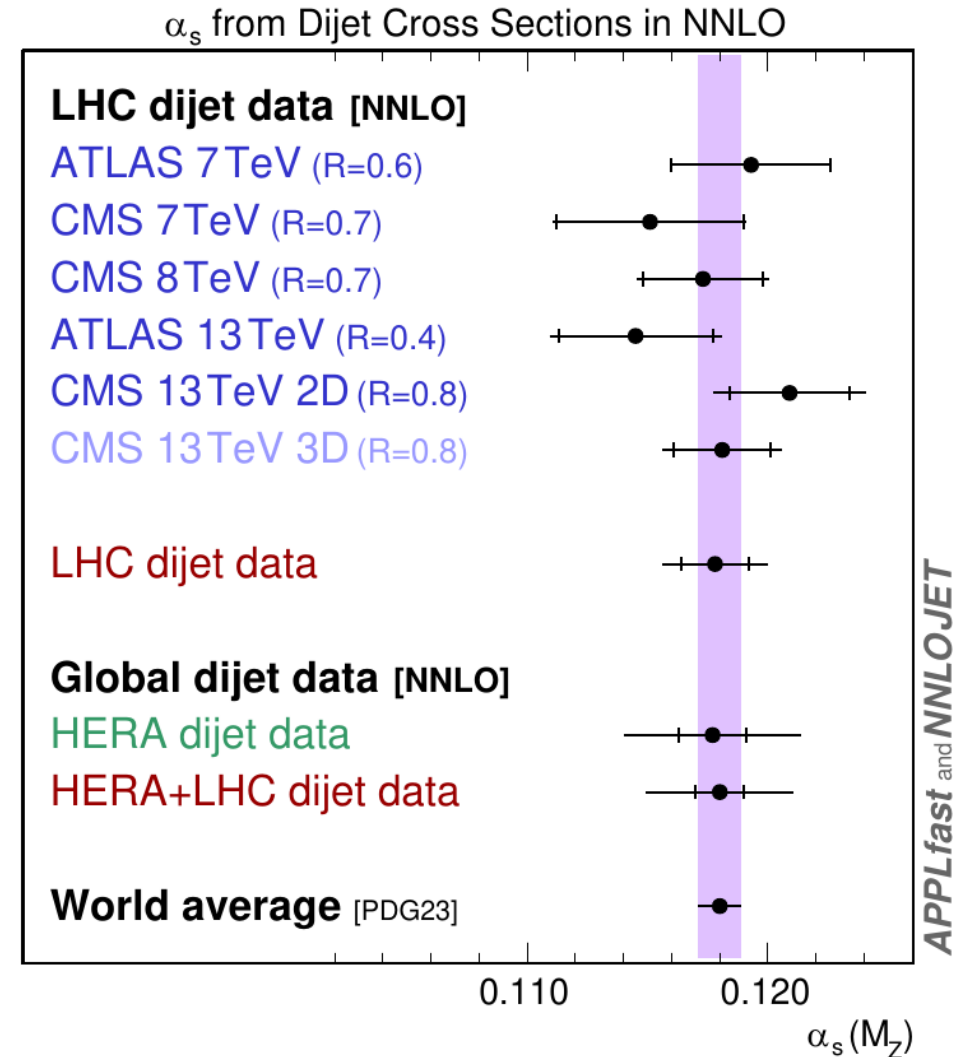
Combining dijet datasets $\rightarrow \alpha_s(Q)$ **ETP**

From LHC dijet data $\alpha_s(m_Z^2) = 0.1178 \pm 0.0014(\text{fitall}) \pm 0.0017(\text{scl})$

$\alpha_s(Q)$ from HERA & LHC dijet data



α_s fit results per dataset



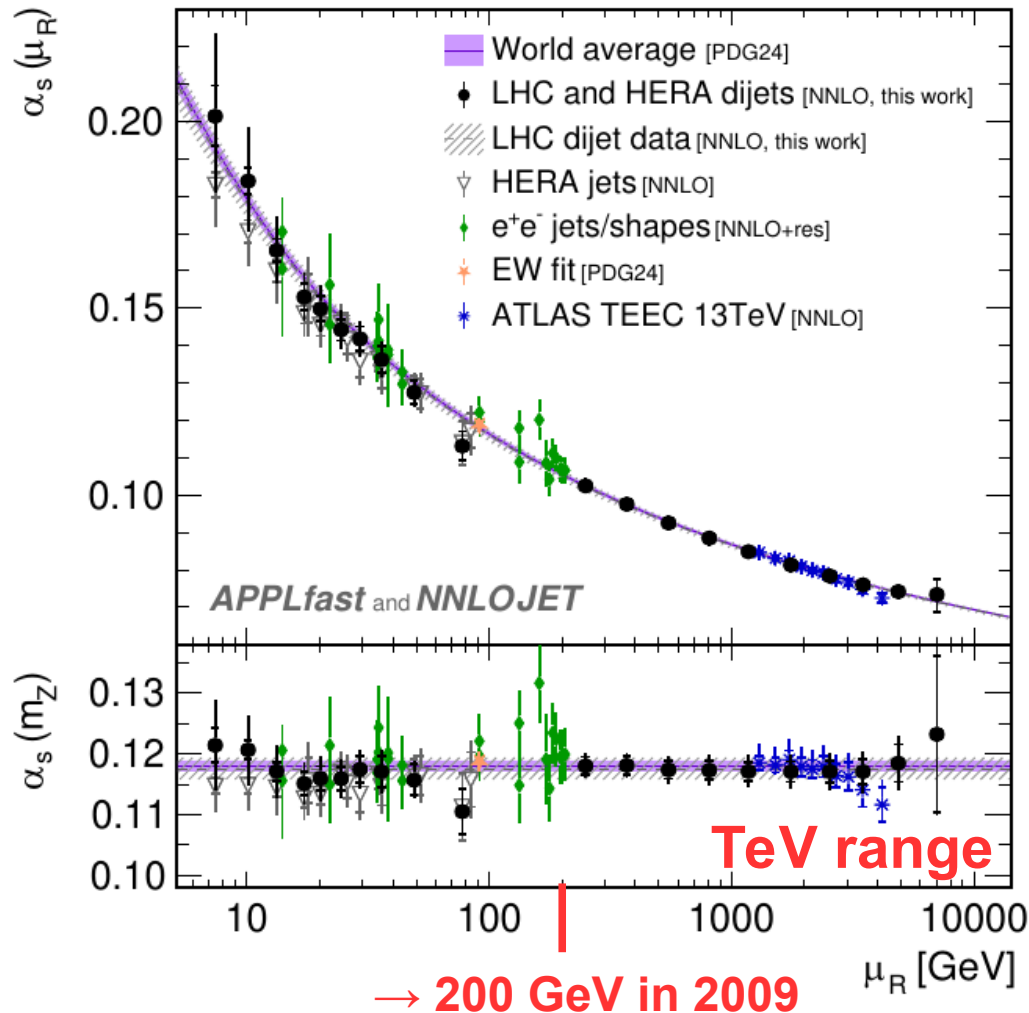
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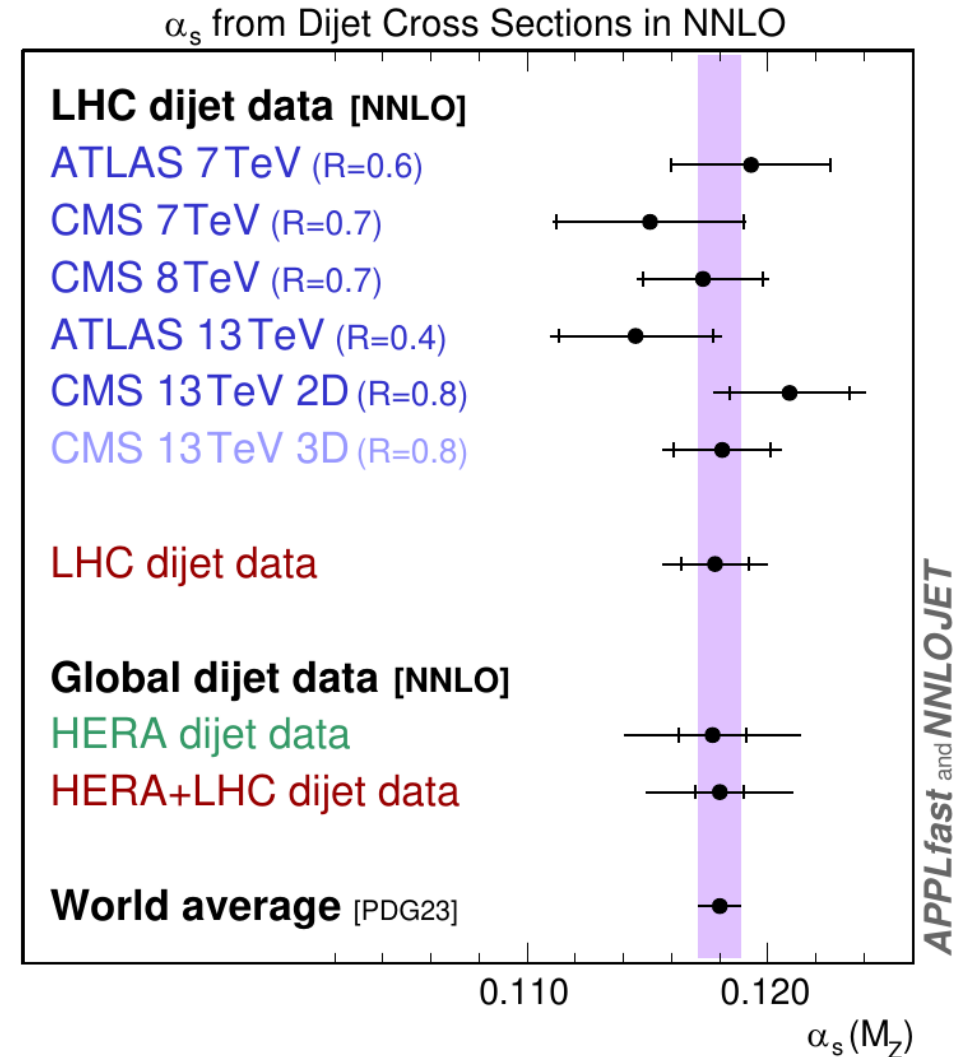
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Comparison to selected other data



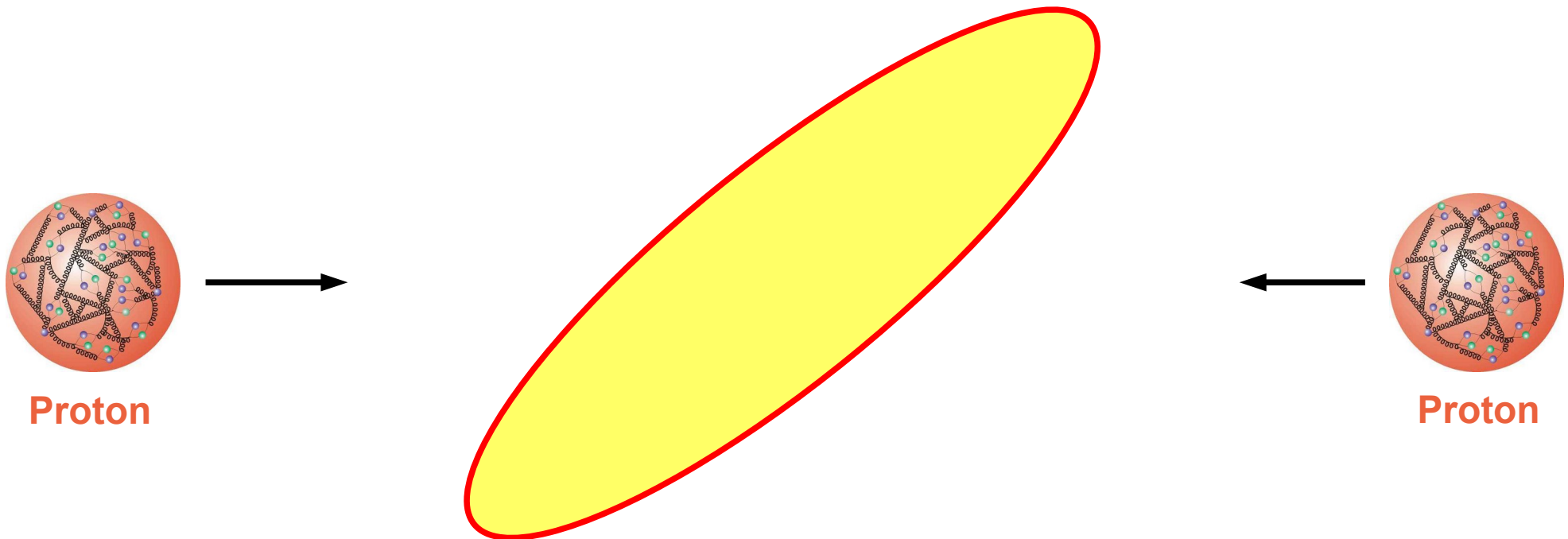
α_s fit results per dataset



Ahmadova, KR, et al., arXiv:2412.21165.



Energy flow





- Analysing the energy flow within an event in bins of some momentum scale and a suitable observable
- Often so-called “event shapes” like thrust or energy-energy correlations
- Go back to definitions suggested for e^+e^- collisions; for pp only transverse momenta used
- Useful for i.a.:
 - ➔ Determination of $\alpha_s(m_Z)$ & test of running of $\alpha_s(Q)$
 - ➔ MC generator comparison & tuning
 - ➔ Search for new physics



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- Independent of luminosity
- Reduced sensitivity to other systematic effects



- Analysing the energy flow within an event in bins of some momentum scale and a suitable observable
- Often so-called “event shapes” like thrust or energy-energy correlations
- Go back to definitions suggested for e^+e^- collisions; for pp only transverse momenta used
- Useful for i.a.:
 - ➔ Determination of $\alpha_s(m_Z)$ & test of running of $\alpha_s(Q)$
 - ➔ MC generator comparison & tuning
 - ➔ Search for new physics
- Independent of luminosity
- Reduced sensitivity to other systematic effects
- Often multi-scale problem → more complicated scale dependence
- Can contain transition region from perturbative to nonperturbative QCD



- Example of an event shape: energy-energy correlation (EEC)

- Goes back to definition in e^+e^-

Basham, Brown, Ellis, Love, PRL41 (1978) 1585;
PRD 19 (1979) 2018.

- Here specialised to pp collisions → only transverse momenta (TEEC)

- Measures E_T weighted azimuthal differences:

$$\frac{1}{\sigma} \frac{d\Sigma}{d \cos \phi} = \frac{1}{N} \sum_{A=1}^N \sum_{ij} \frac{E_{Ti}^A E_{Tj}^A}{\left(\sum_k E_{Tk}^A\right)^2} \delta(\cos \phi - \cos \phi_{ij})$$

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Transverse jet
energy weights

Measure of pairwise
azimuthal distance
between jets i, j

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Normalisation to
no. of events N

Normalisation to
 E_T sum per event

Sum over all
events in sample

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Normalisation to
no. of events N

Normalisation to
 E_T sum per event

Sum over all
events in sample

Additional binning in
e.g. transverse energy
sum $H_T \rightarrow$ scale Q

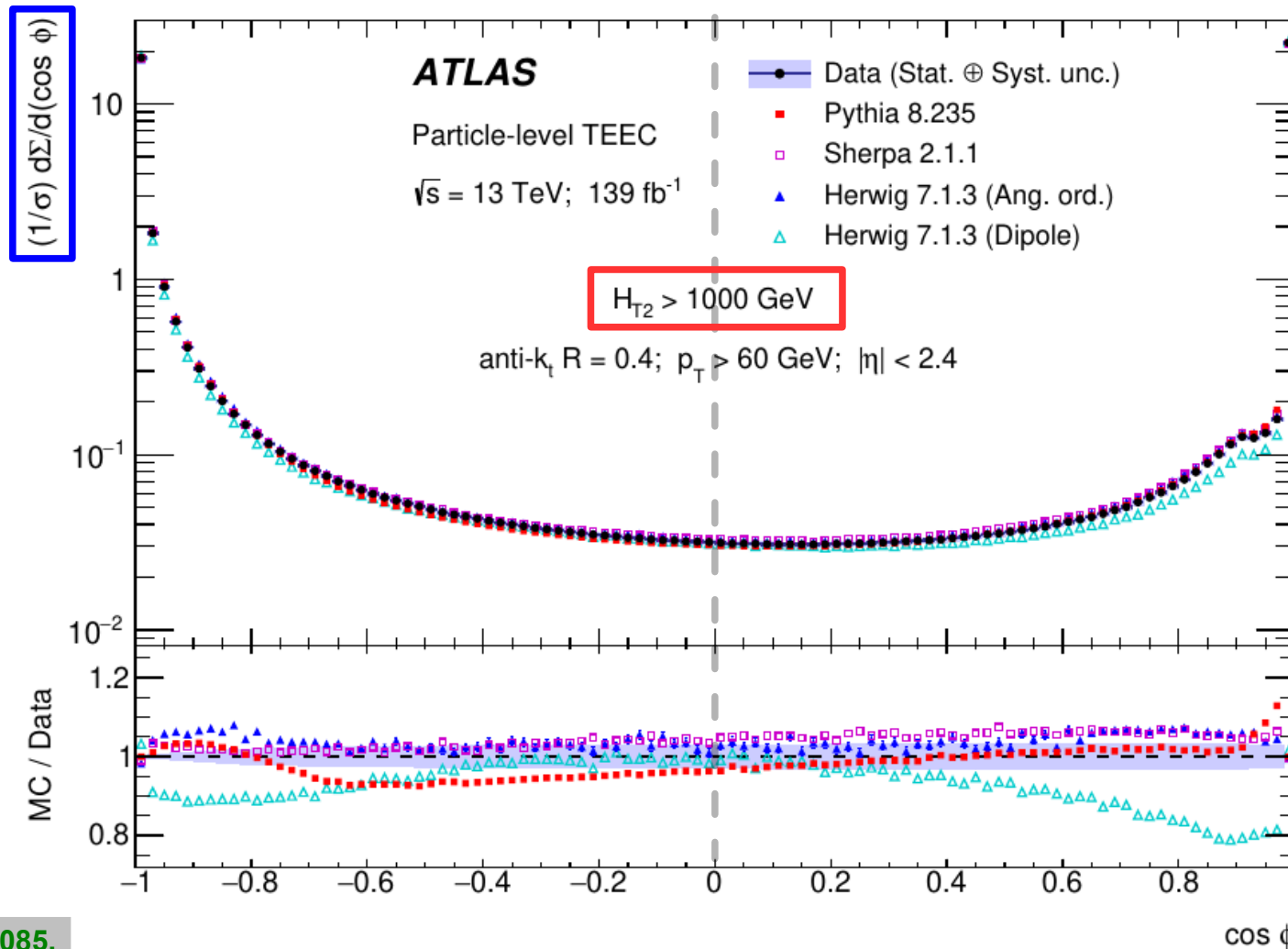


Transverse energy-energy correlation

$$\frac{1}{\sigma} \frac{d\Sigma}{d \cos \phi} = \frac{1}{N} \sum_{A=1}^N \sum_{ij} \frac{E_{Ti}^A E_{Tj}^A}{\left(\sum_k E_{Tk}^A\right)^2} \delta(\cos \phi - \cos \phi_{ij})$$

Event shape: $\text{TEEC} \propto \alpha_s$

Normalised



Multiple
bins in H_T



Transverse energy-energy correlation

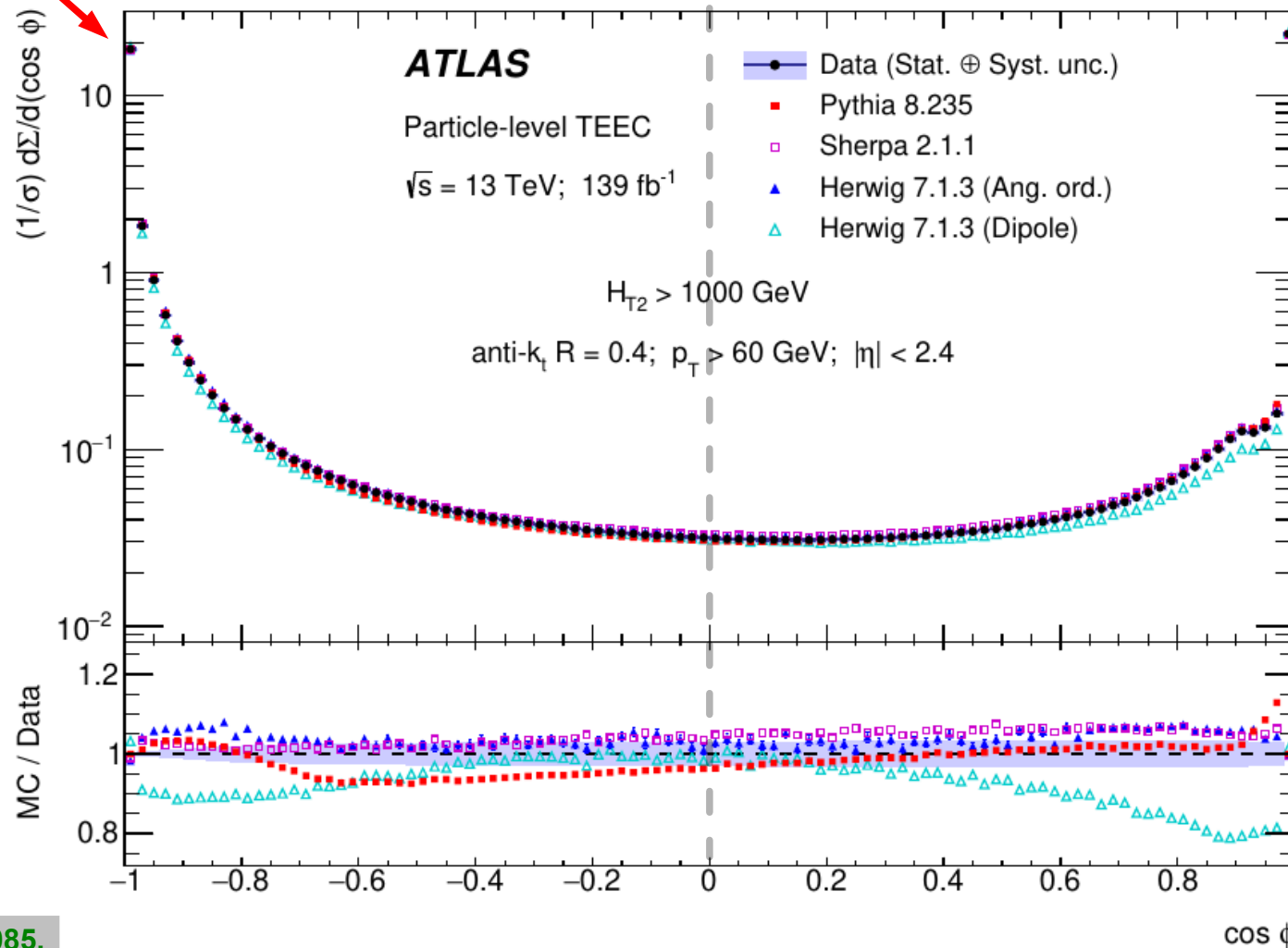
$$\frac{1}{\sigma} \frac{d\Sigma}{d \cos \phi} = \frac{1}{N} \sum_{A=1}^N \sum_{ij} \frac{E_{Ti}^A E_{Tj}^A}{\left(\sum_k E_{Tk}^A\right)^2} \delta(\cos \phi - \cos \phi_{ij})$$

Event shape: $\text{TEEC} \propto \alpha_s$

2 → 2 back-to-back jets

autocorrelation $i=j$ 2 → 2

Normalised



Multiple bins in H_T



Transverse energy-energy correlation

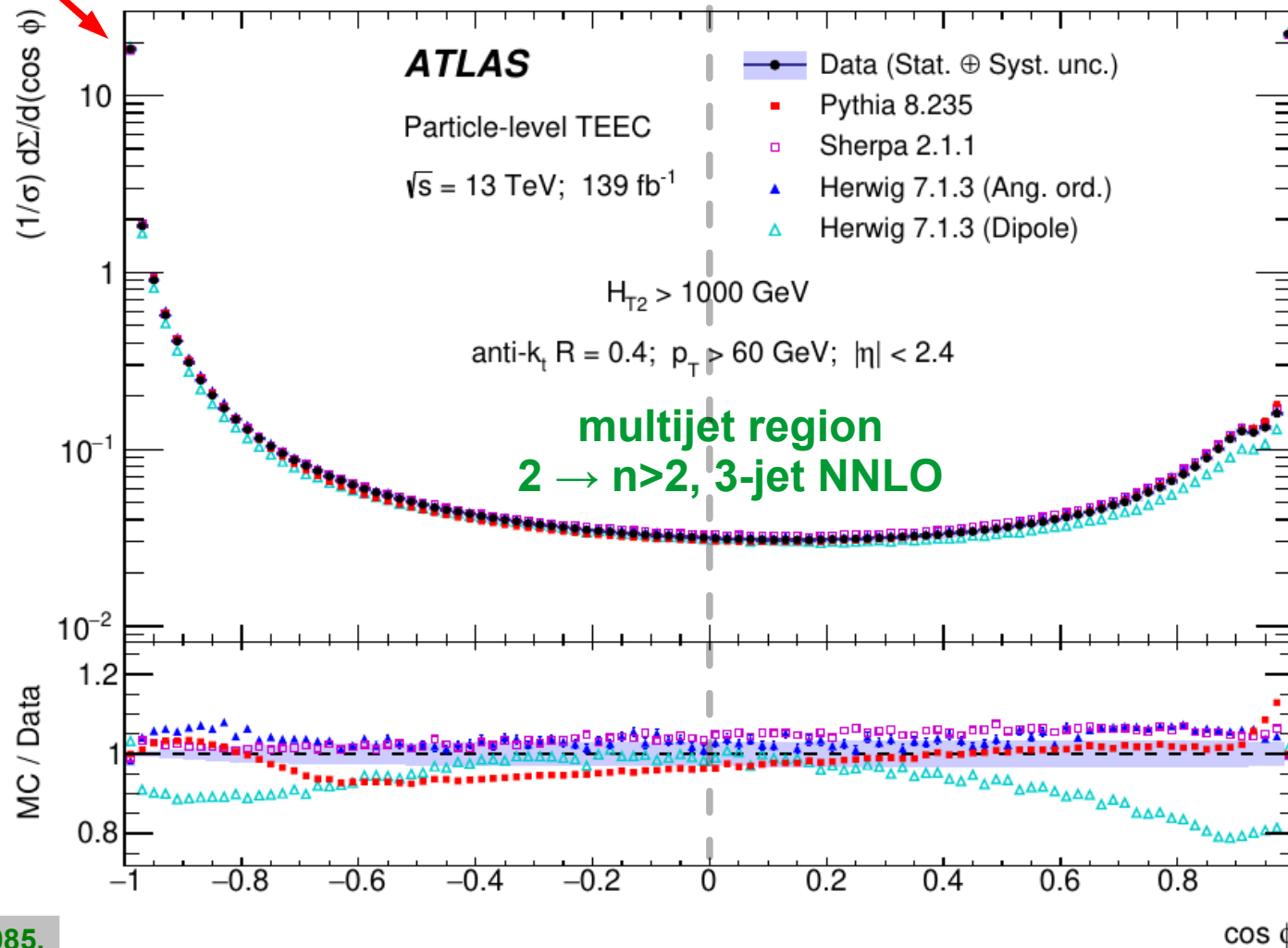
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Multiple bins in H_T

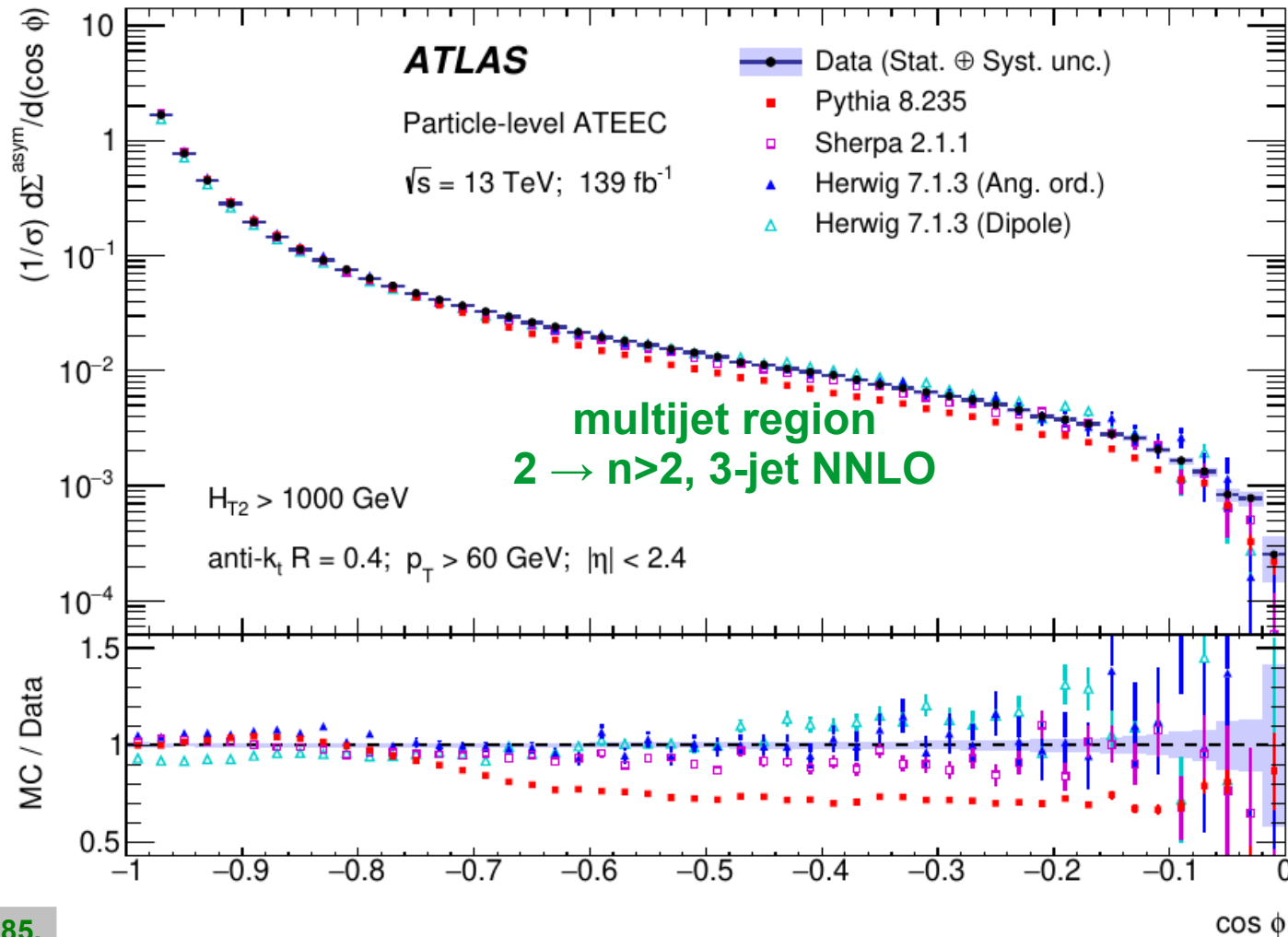


$$\frac{1}{\sigma} \frac{d\Sigma^{\text{asym}}}{d\cos\phi} = \frac{1}{\sigma} \frac{d\Sigma}{d\cos\phi} \bigg|_{\phi} - \frac{1}{\sigma} \frac{d\Sigma}{d\cos\phi} \bigg|_{\pi-\phi}$$

Event shape: $\text{ATEEC} \propto \alpha_s$

Asymmetry

Normalised



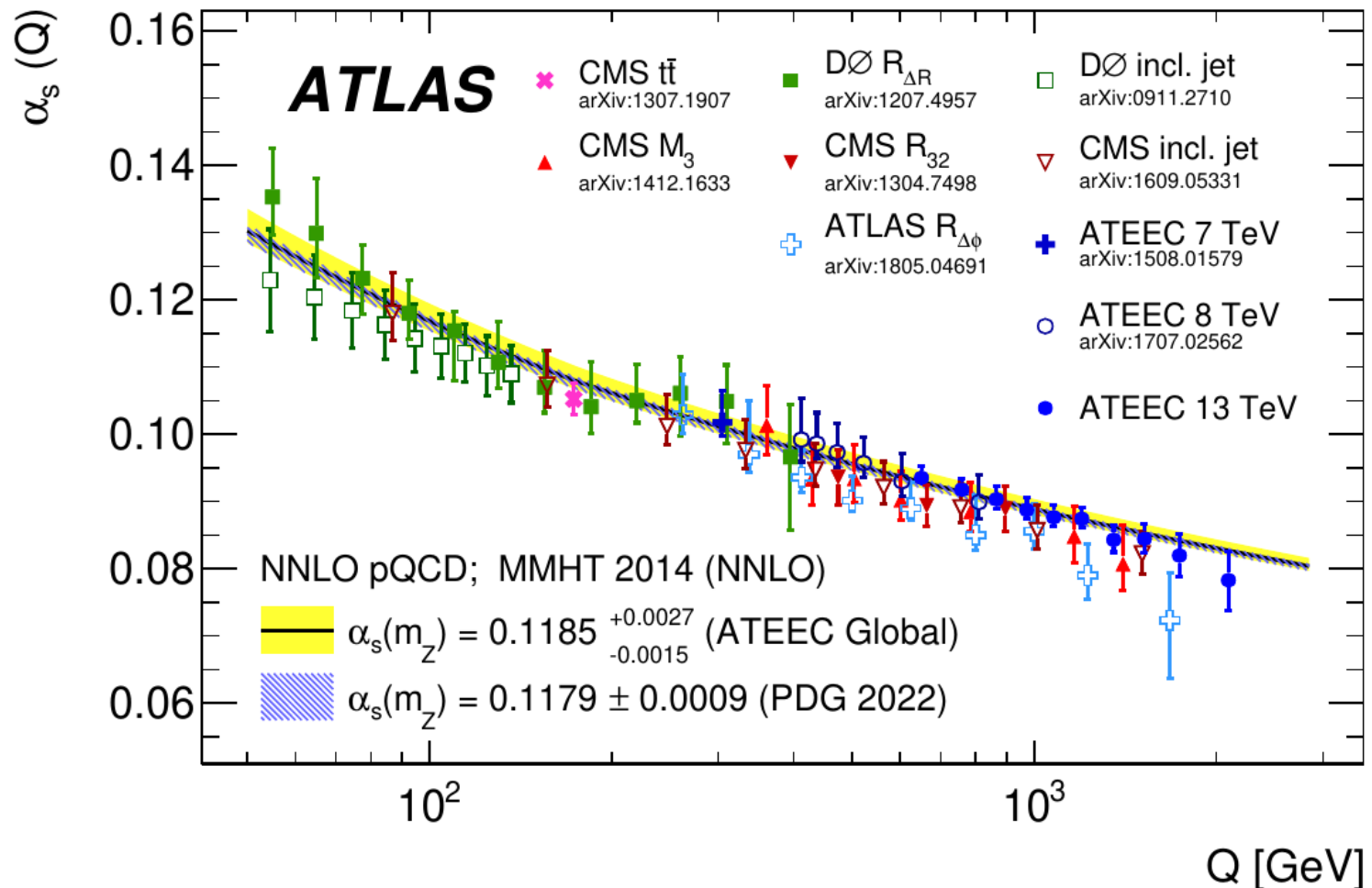
Multiple bins in H_T

TEEC $\alpha_s(m_Z) = 0.1175 \pm 0.0006 \text{ (exp.)}^{+0.0034}_{-0.0017} \text{ (theo.)}$

ATEEC $\alpha_s(m_Z) = 0.1185 \pm 0.0009 \text{ (exp.)}^{+0.0025}_{-0.0012} \text{ (theo.)}$

Remark:
Reduction in MHO
uncertainty smaller
than maybe anticipated

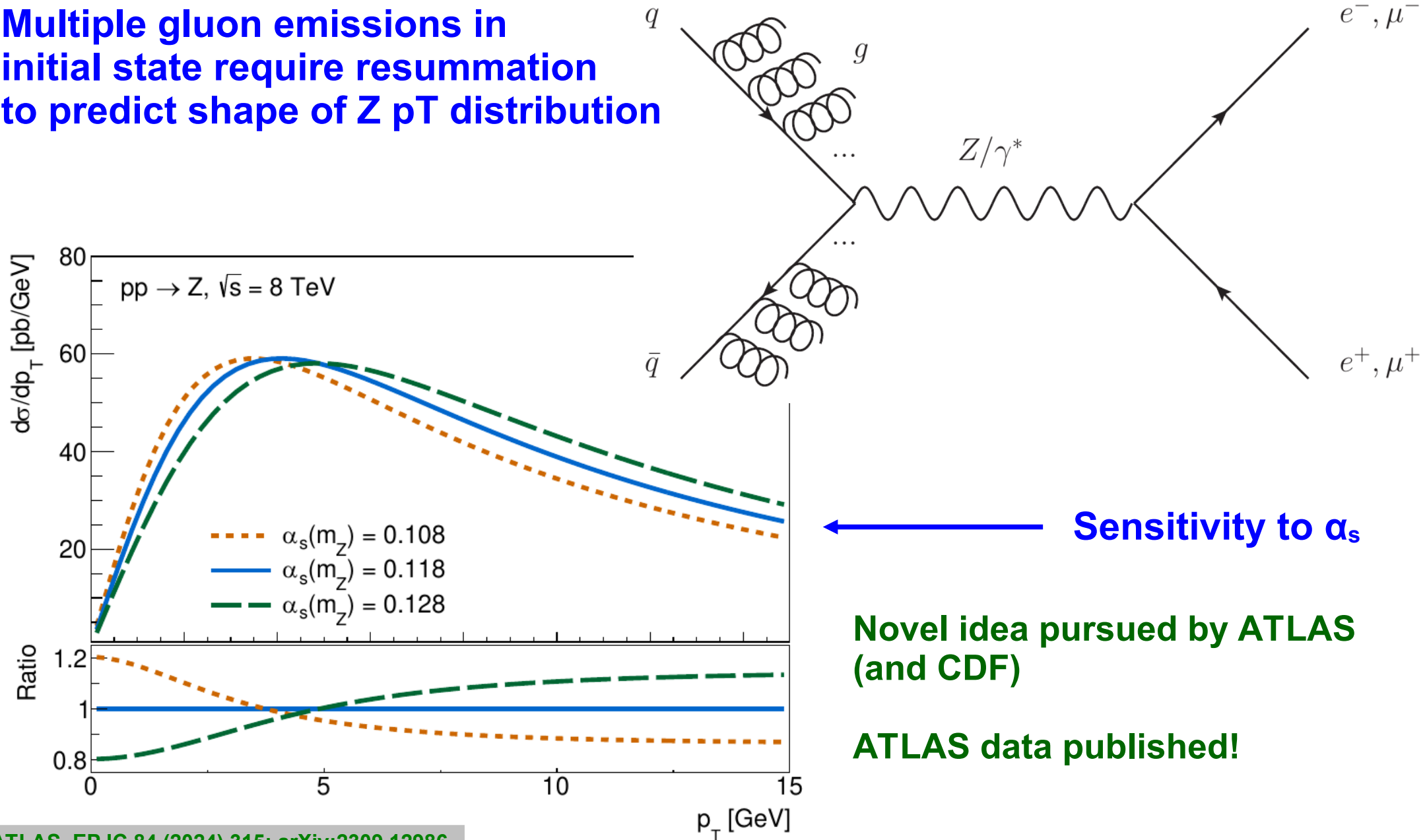
First “running”
points from pp
observable at NNLO





Sudakov peak of $DY Z p_T$

Multiple gluon emissions in initial state require resummation to predict shape of $Z p_T$ distribution

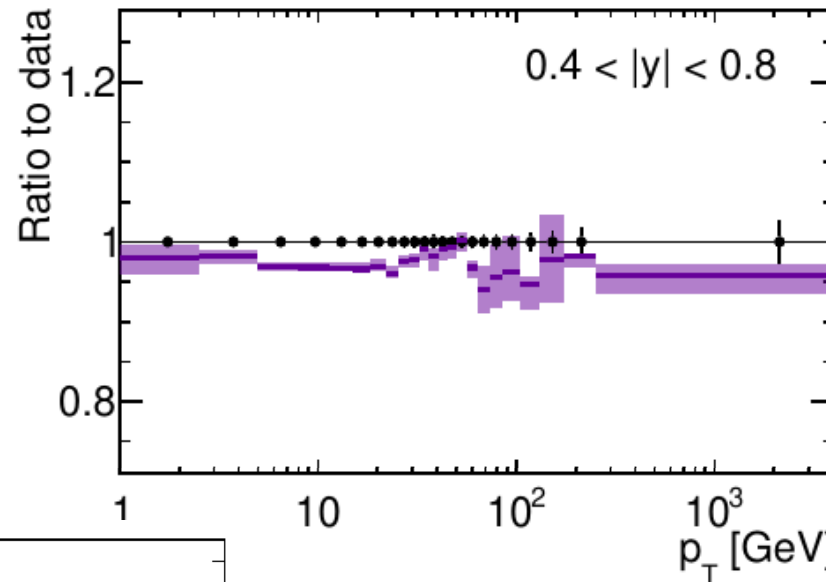


ATLAS, EPJC 84 (2024) 315; arXiv:2309.12986.



Sudakov peak of $DY Z p_T$

Ratio to data of DYTurbo
resummed+matched fixed-order
prediction



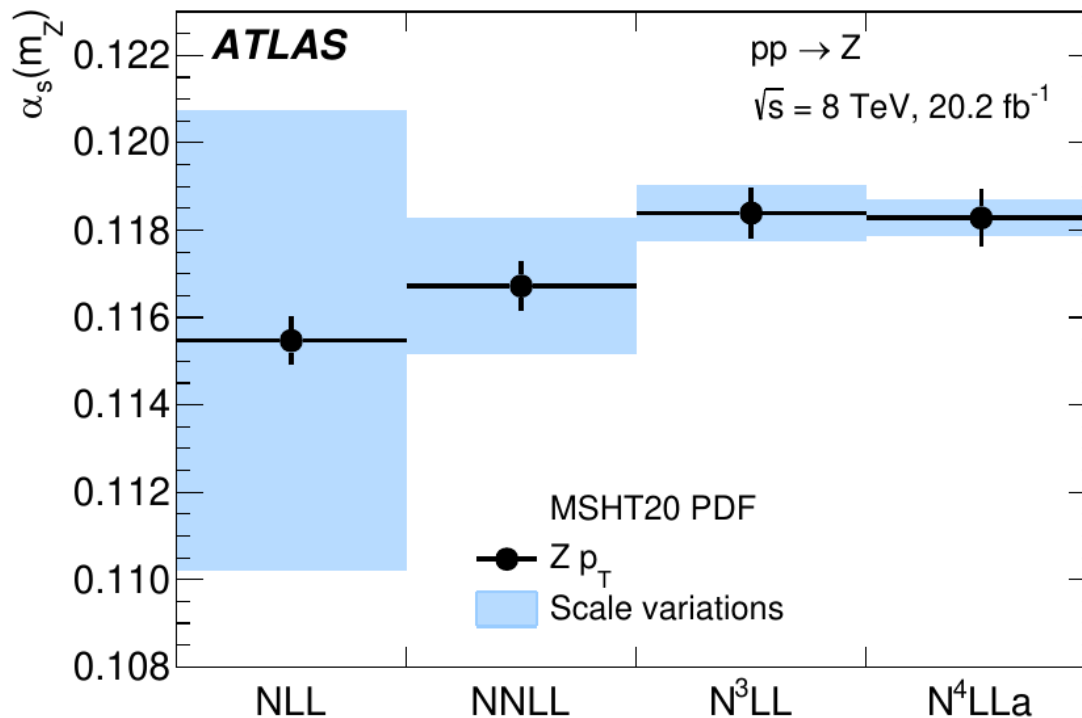
ATLAS

$pp \rightarrow Z$

$\sqrt{s} = 8 \text{ TeV}, 20.2 \text{ fb}^{-1}$

$\frac{d\sigma}{dp_T}$

DYTurbo



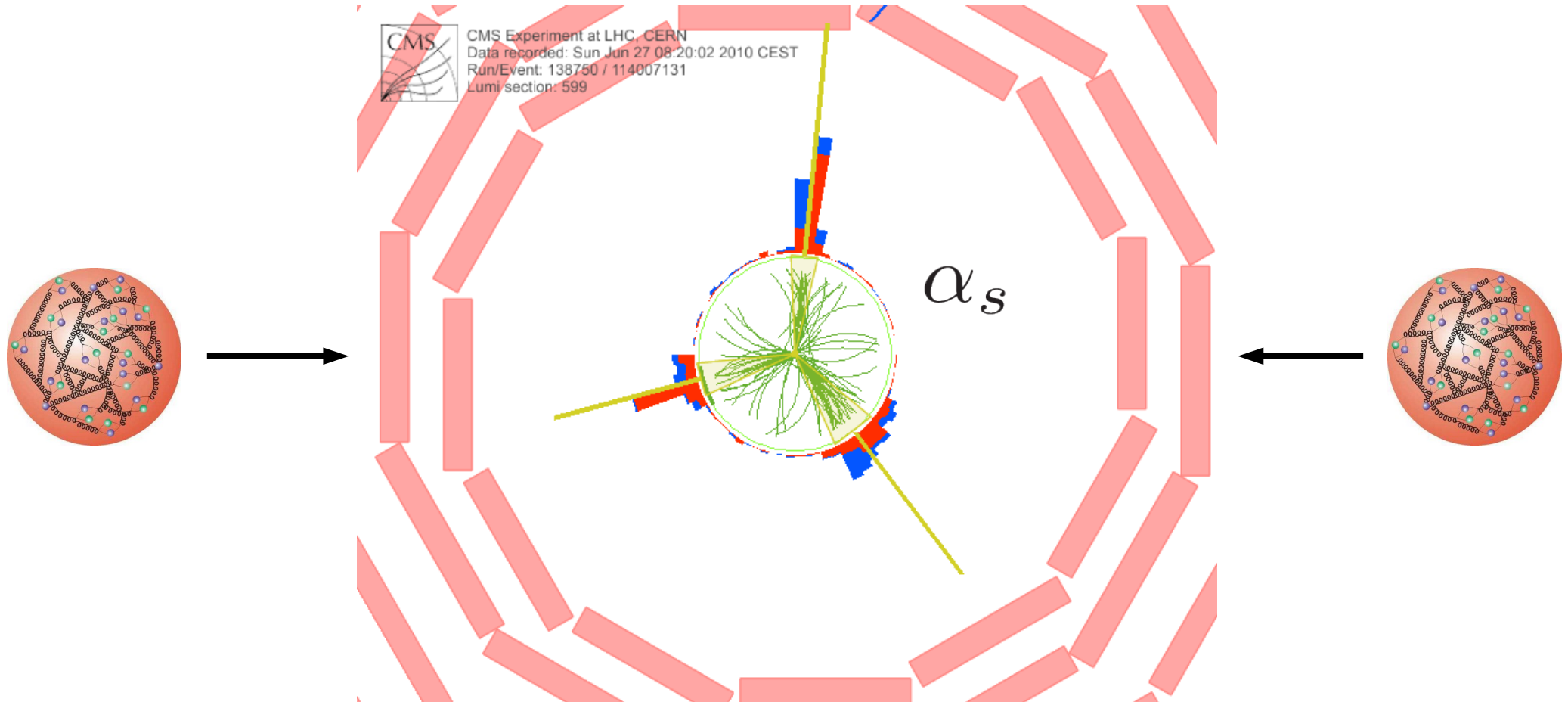
Series of α_s extractions with
increasing precision

$$\alpha_s(M_Z) = 0.1183 \pm 0.0009$$

Most precise single determination (except
Lattice), but so far not accepted by journal.
Stimulated discussion notably on theory
uncertainty.



Higher multiplicity





- **Aim to reduce or eliminate impact of systematic effects**
- **Useful for i.a.:**
 - ➔ **Determination of $\alpha_s(m_Z)$ & test of running of $\alpha_s(Q)$**
 - ➔ **MC generator comparison & tuning**
 - ➔ **Investigation of jet size R dependence**
 - ➔ **Search for new physics**



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 - ➔ Search for new physics
- Often independent of luminosity
- Reduced or eliminated sensitivity to systematic effects



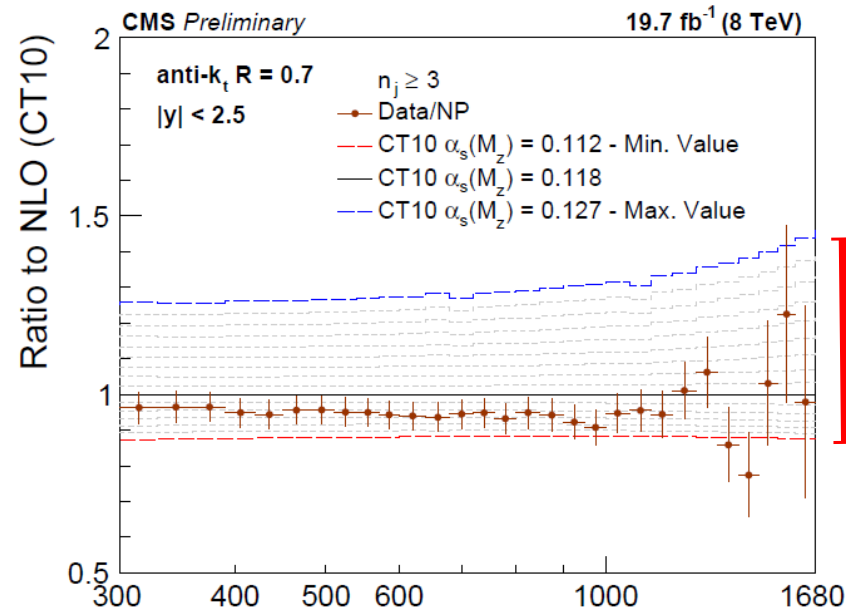
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 - ➔ Investigation of jet size R dependence
 - ➔ Search for new physics
- Often independent of luminosity
- Reduced or eliminated sensitivity to systematic effects
- Correlations between numerator & denominator
- Often multi-scale problem → more complicated scale dependence
- Also reduced sensitivity to desired effect



Sensitivity vs. systematic effects

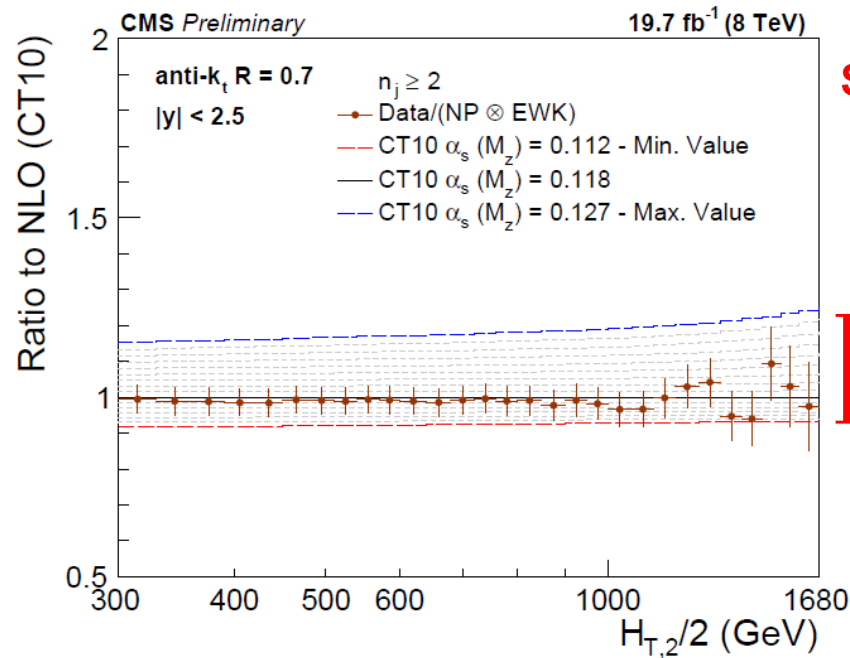
Inclusive 3-jet
cross section

$$\sigma_{3j} \propto \alpha_s^3$$



Inclusive 2-jet
cross section

$$\sigma_{2j} \propto \alpha_s^2$$



Sensitivity



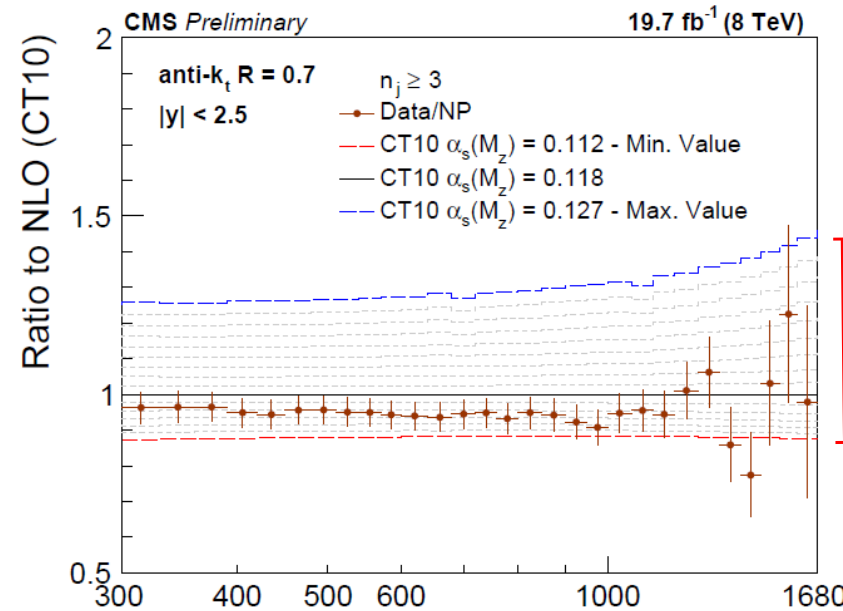
Sensitivity vs. systematic effects

Inclusive 3-jet cross section

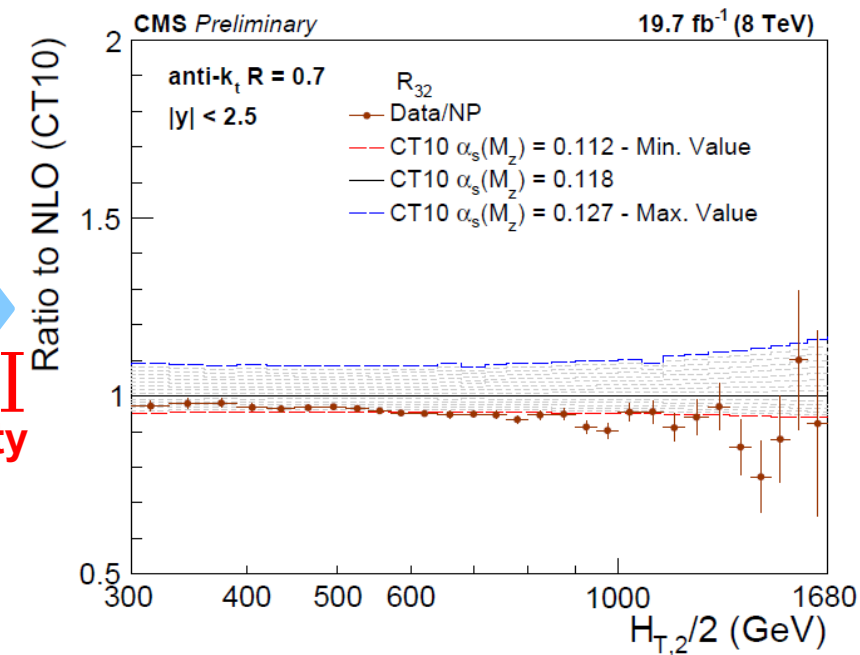
$$\sigma_{3j} \propto \alpha_s^3$$

Inclusive 3-jet to inclusive 2-jet cross section ratio

$$R_{3/2} \propto \alpha_s$$

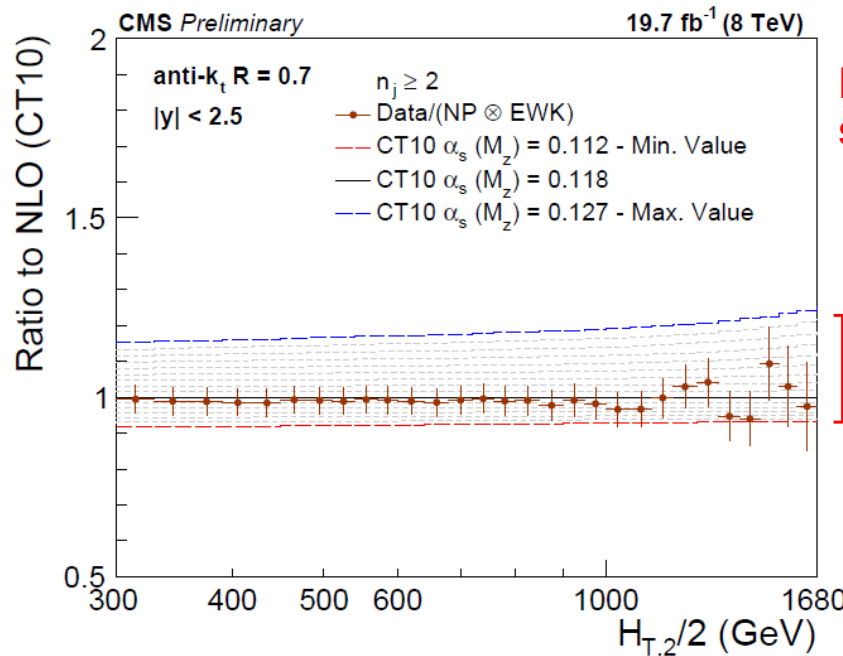


Reduced sensitivity



Inclusive 2-jet cross section

$$\sigma_{2j} \propto \alpha_s^2$$





Sensitivity vs. systematic effects

Inclusive 3-jet cross section

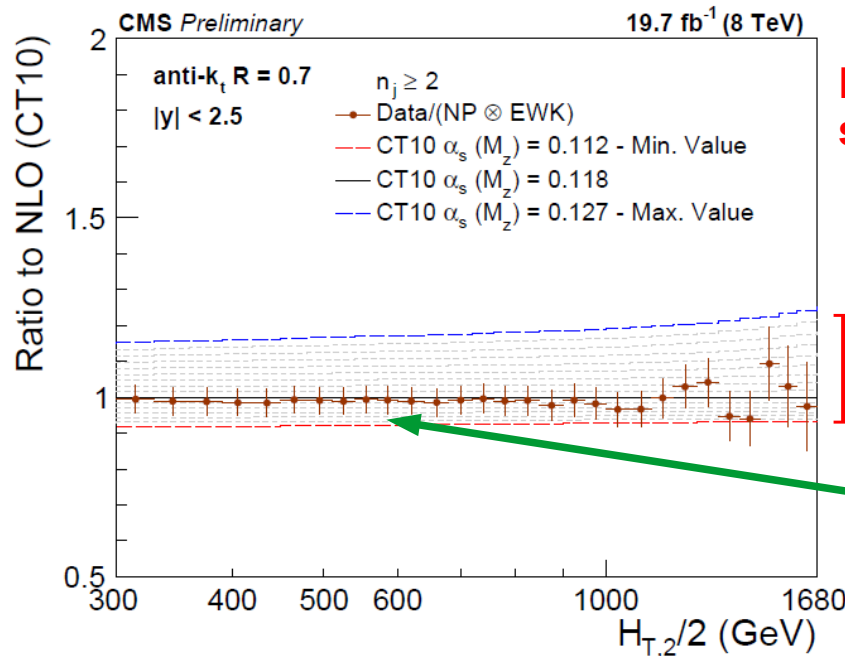
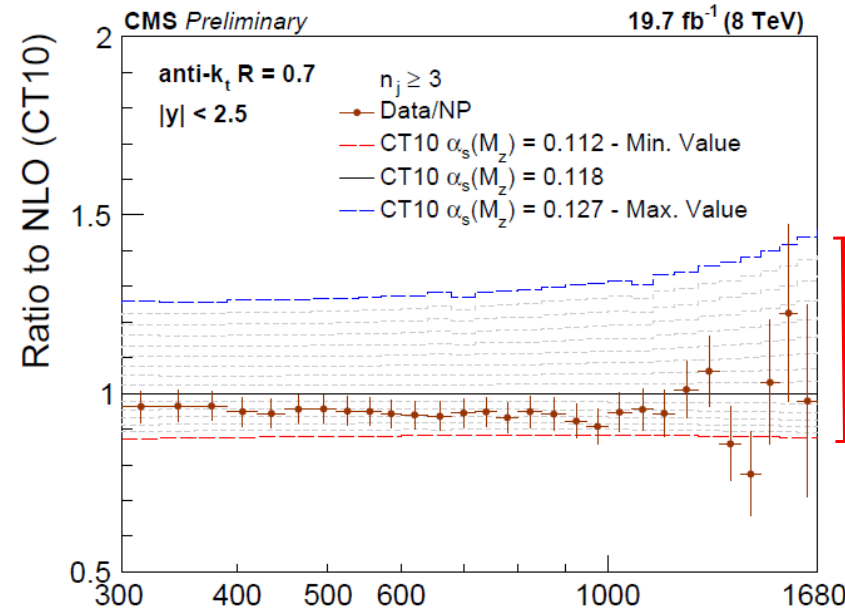
$$\sigma_{3j} \propto \alpha_s^3$$

Inclusive 3-jet to inclusive 2-jet cross section ratio

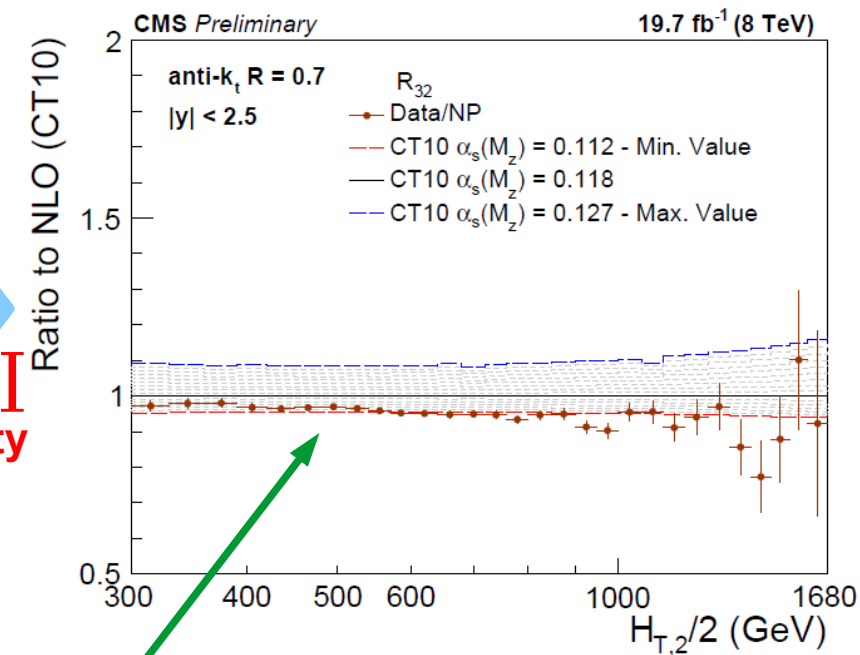
$$R_{3/2} \propto \alpha_s$$

Inclusive 2-jet cross section

$$\sigma_{2j} \propto \alpha_s^2$$



Reduced sensitivity

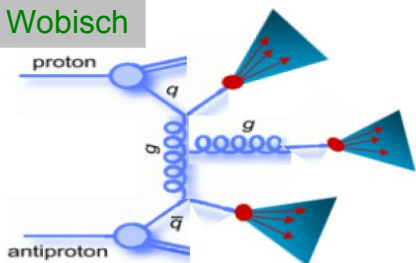


Much reduced systematic uncertainty



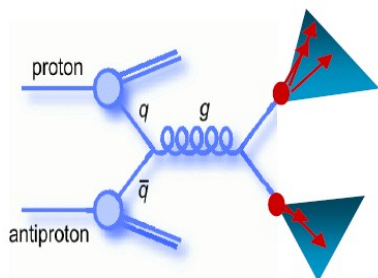
3- to 2-jet ratios

M. Wobisch



$R_{3/2}$

$\longleftrightarrow \alpha_s$

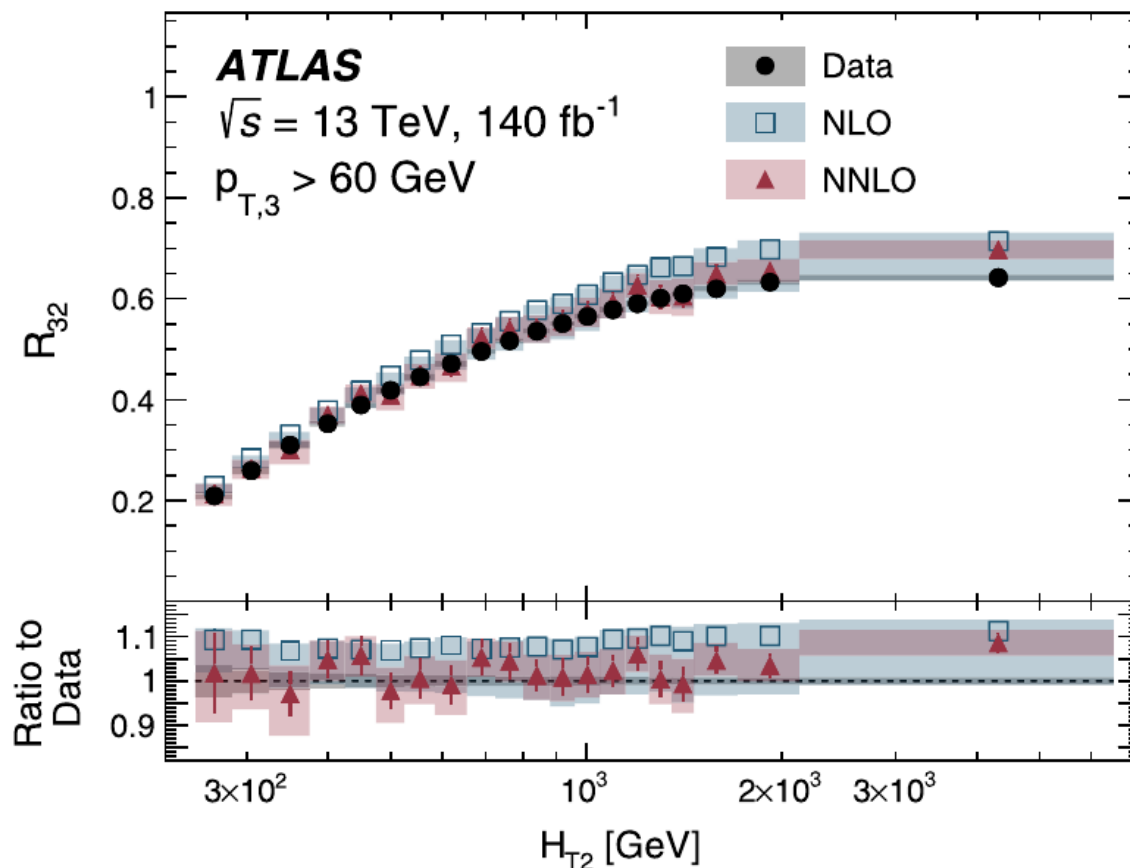


$$\frac{\sigma_{3+\text{jet}}}{\sigma_{2+\text{jet}}} \propto \alpha_s^1$$

Nice measurement from ATLAS:
 $R_{3/2}$, $R_{4/2}$, $R_{4/3}$, $R_{5/4}$

First comparison of $R_{3/2}$ vs. NNLO

But no α_s fit (yet).



ATLAS, PRD 110 (2024) 072019.



- Ratio observable differential in jet p_T , where:
 - ➔ Numerator counts no. of neighbouring jets with minimal p_T within azimuthal distance $2\pi/3 < \Delta\Phi < 7\pi/8 \rightarrow$ enforces ≥ 3 -jet configuration
 - ➔ Denominator counts all jets in p_T bin

$$R_{\Delta\phi}(p_T) = \frac{\sum_{i=1}^{N_{\text{jet}}(p_T)} N_{\text{nbr}}^{(i)}(\Delta\phi, p_{T\text{min}}^{\text{nbr}})}{N_{\text{jet}}(p_T)}$$

$$R_{\Delta\phi}(p_T) \propto \alpha_s$$

CMS, EPJC 84 (2024) 842.

- ➔ Requires 3-jet NNLO
- ➔ Similar observables previously used by
 - ➔ **D0:** $R_{\Delta R}(p_T)$ D0, PLB 718 (2012) 56.
 - ➔ **ATLAS:** $R_{\Delta\phi}(H_T)$ ATLAS, PRD 98 (2018) 092044.



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 - ➔ **Denominator counts all jets in p_T bin**

$$R_{\Delta\phi}(p_T) = \frac{\sum_{i=1}^{N_{\text{jet}}(p_T)} N_{\text{nbr}}^{(i)}(\Delta\phi, p_{T\text{min}}^{\text{nbr}})}{N_{\text{jet}}(p_T)} \quad R_{\Delta\phi}(p_T) \propto \alpha_s$$

CMS, EPJC 84 (2024) 842.

- ➔ **Nice feature: Equivalent definition with 2D quantity $N(p_T, n)$ counting neighbouring jets**

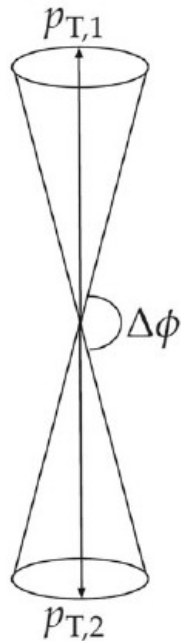
$$R_{\Delta\phi}(p_T) = \frac{\sum_n n N(p_T, n)}{\sum_n N(p_T, n)}$$

- ➔ **enables unfolding accounting for all correlations**

Example configurations (on request)

$$R_{\Delta\phi}(p_T) = \frac{\sum_{i=1}^{N_{\text{jet}}(p_T)} N_{\text{nbr}}^{(i)}(\Delta\phi, p_{T\text{min}}^{\text{nbr}})}{N_{\text{jet}}(p_T)}$$

2-jet topology



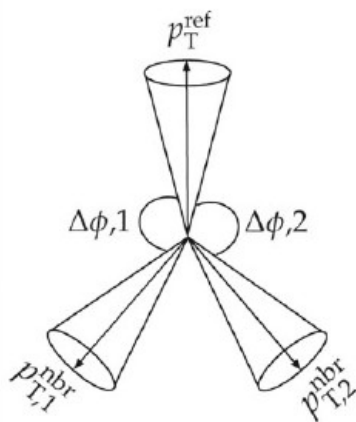
$R_{\Delta\phi}(p_T)$ entries

$\Delta\phi \approx \pi$

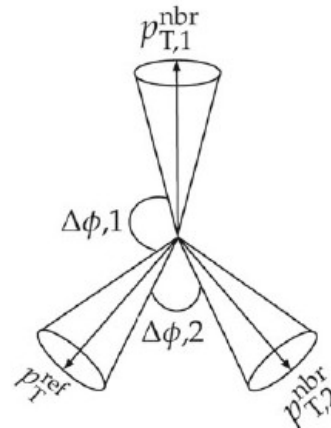
Numerator: 0

Denominator: 2

3-jet topology (all jets with $p_T > 100$ GeV)



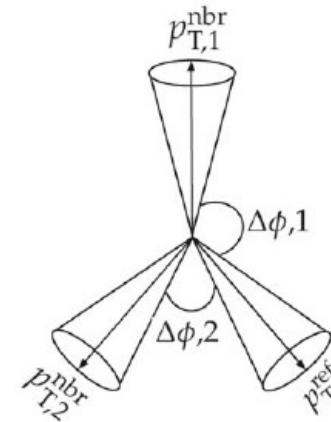
Numerator: 2
 $2\pi/3 < \Delta\phi,1 < 7\pi/8$
 $2\pi/3 < \Delta\phi,2 < 7\pi/8$



$R_{\Delta\phi}(p_T)$ entries

Numerator: 1
 $2\pi/3 < \Delta\phi,1 < 7\pi/8$
 $\Delta\phi,2 < 2\pi/3$

Denominator: 3



Numerator: 1
 $2\pi/3 < \Delta\phi,1 < 7\pi/8$
 $\Delta\phi,2 < 2\pi/3$

CMS, EPJC 84 (2024) 842.



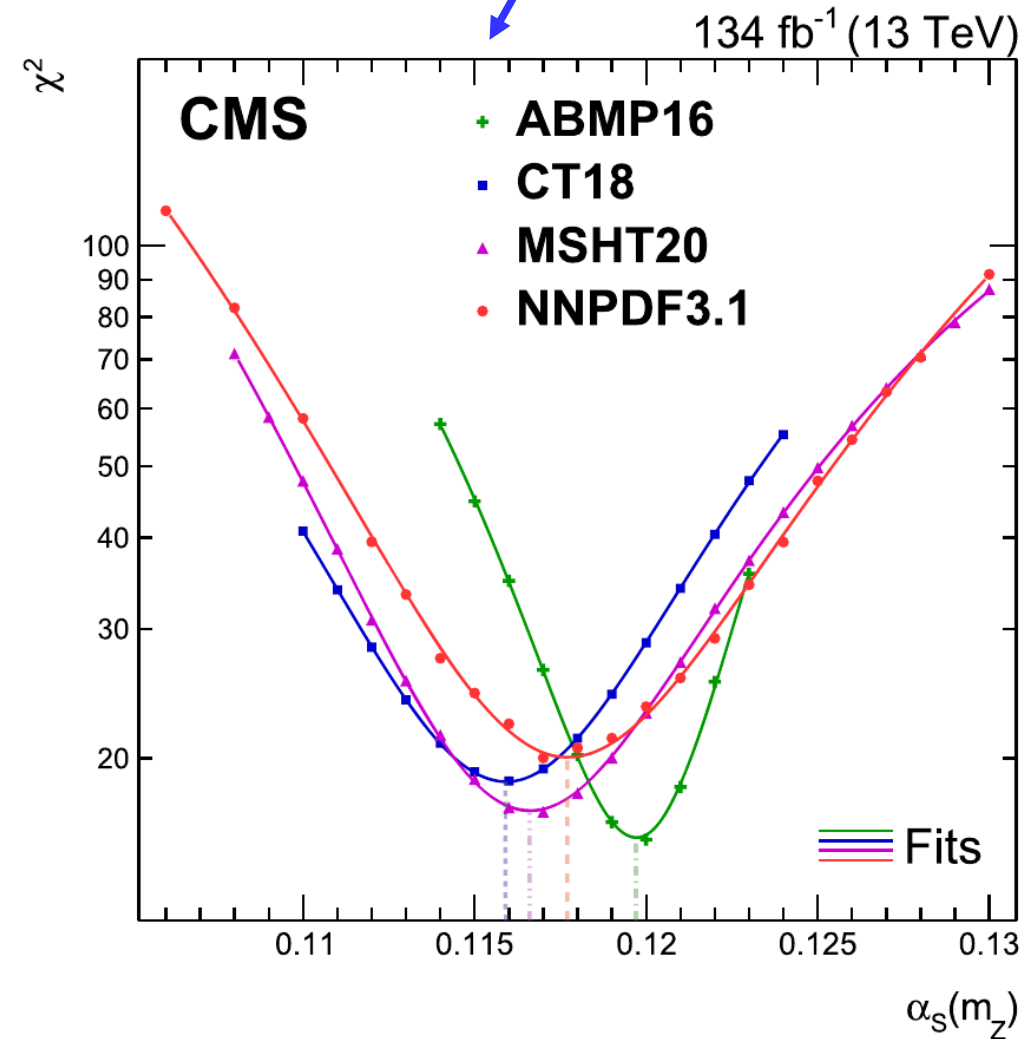
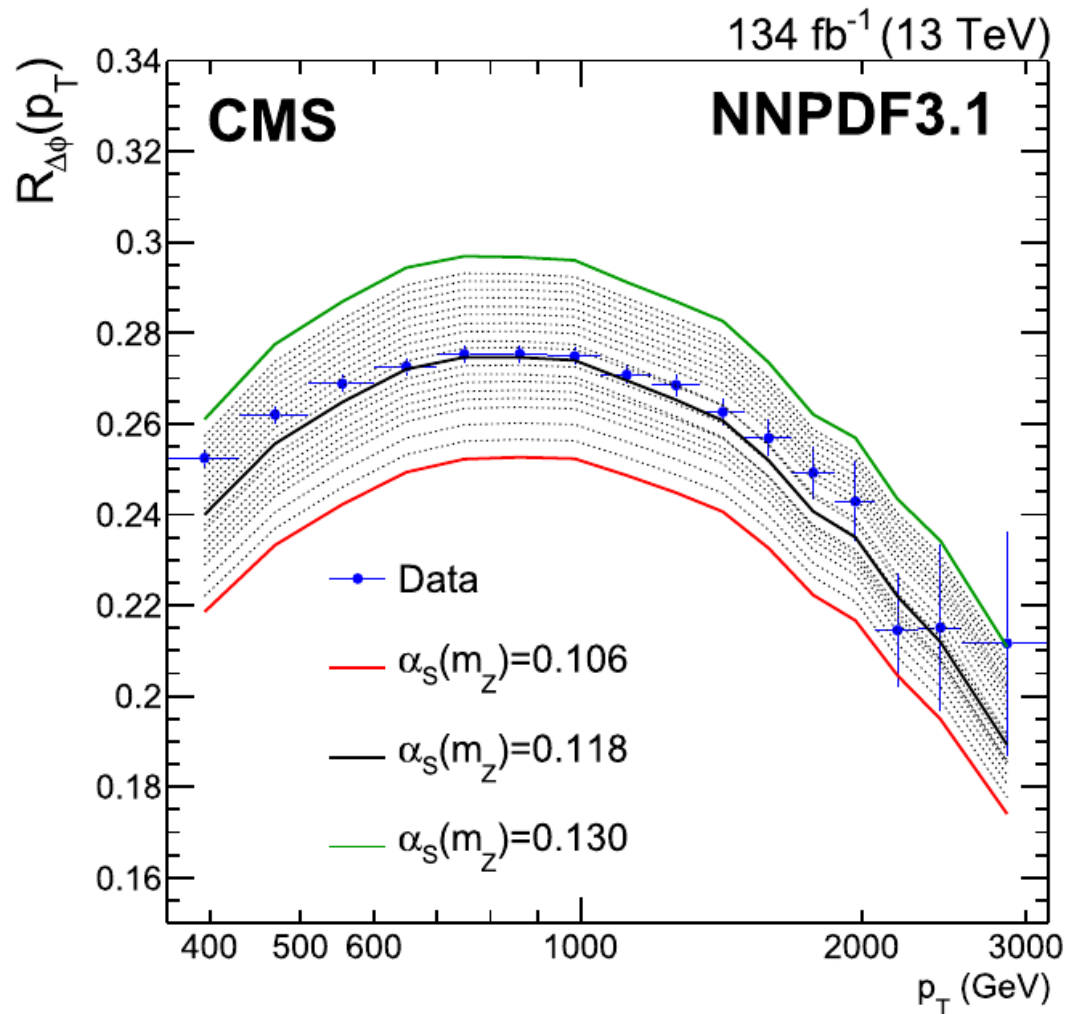
Multijet azimuthal correlation

Good sensitivity, but

So far only NLO → huge MHO uncertainty

Significant PDF dependence

$$\alpha_s(m_Z) = 0.1177 \pm 0.0028(\text{all}) \quad \boxed{+0.0114(\text{scl}) \\ -0.0068}$$





N-point E-E correlators in jets

Jet substructure variable representing correlations of energy flow inside jets

➔ **2-point energy correlators**

$$E2C = \frac{d\sigma}{dx_L} = \sum_{i,j}^n \int d\sigma \underbrace{\frac{E_i E_j}{E_{\text{jet}}^2}}_{\text{weight}} \underbrace{\delta(x_L - \Delta R_{ij})}_{\text{distance}}$$

➔ **multiple entries, e.g. n=3 → 9 pairs inside jet**

Chen et al., arXiv:2004.11381;
Lee et al., arXiv:2205.0314;
Chen et al., arXiv:2307.07510.



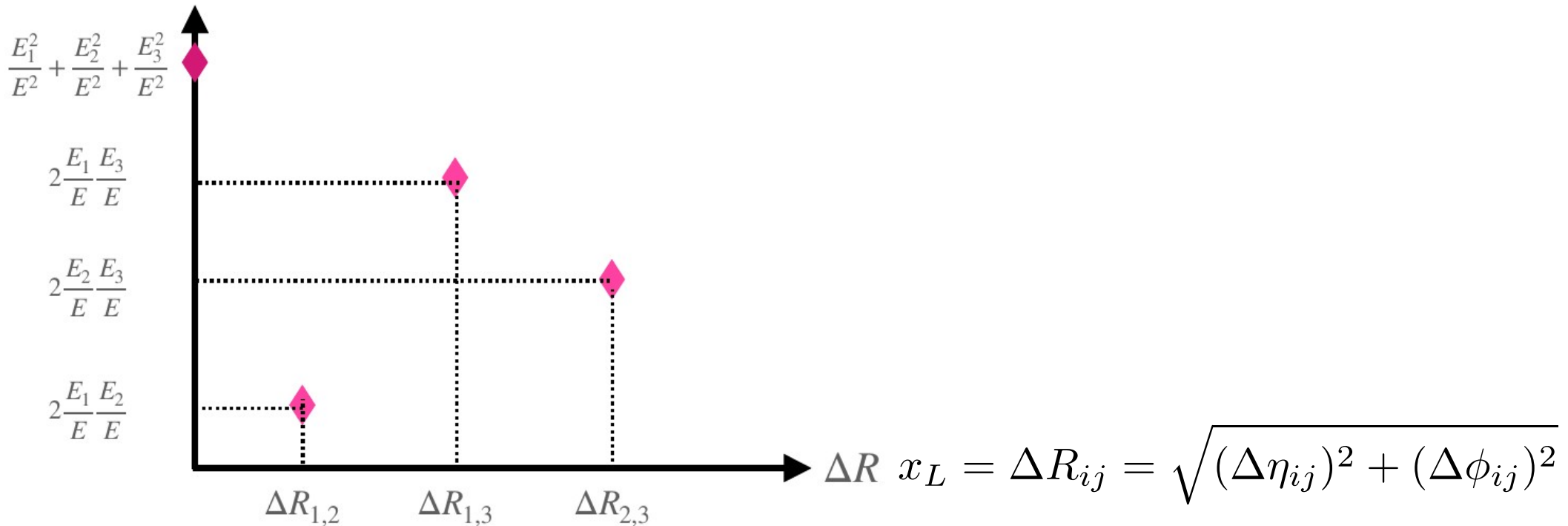
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➔ multiple entries, e.g. n=3 → 9 pairs inside jet



Yulei Ye

3 * 3 = 9 pairs

Chen et al., arXiv:2004.11381;
Lee et al., arXiv:2205.0314;
Chen et al., arXiv:2307.07510.



N-point E-E correlators in jets

Jet substructure variable representing correlations of energy flow inside jets

➔ 3-point energy correlators

$$E3C = \sum_{i,j,k}^n \int d\sigma \underbrace{\frac{E_i E_j E_k}{E_{\text{jet}}^3}}_{\text{weight}} \delta \left(x_L - \underbrace{\max \{ \Delta R_{ij}, \Delta R_{jk}, \Delta R_{ki} \}}_{\text{distance}} \right)$$

➔ e.g. n=3 → 27 triplets inside jet

$$x_L = \max \Delta R_{ij}, \Delta R_{jk}, \Delta R_{ki}$$

Chen et al., arXiv:2004.11381;
Lee et al., arXiv:2205.0314;
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N-point E-E correlators in jets

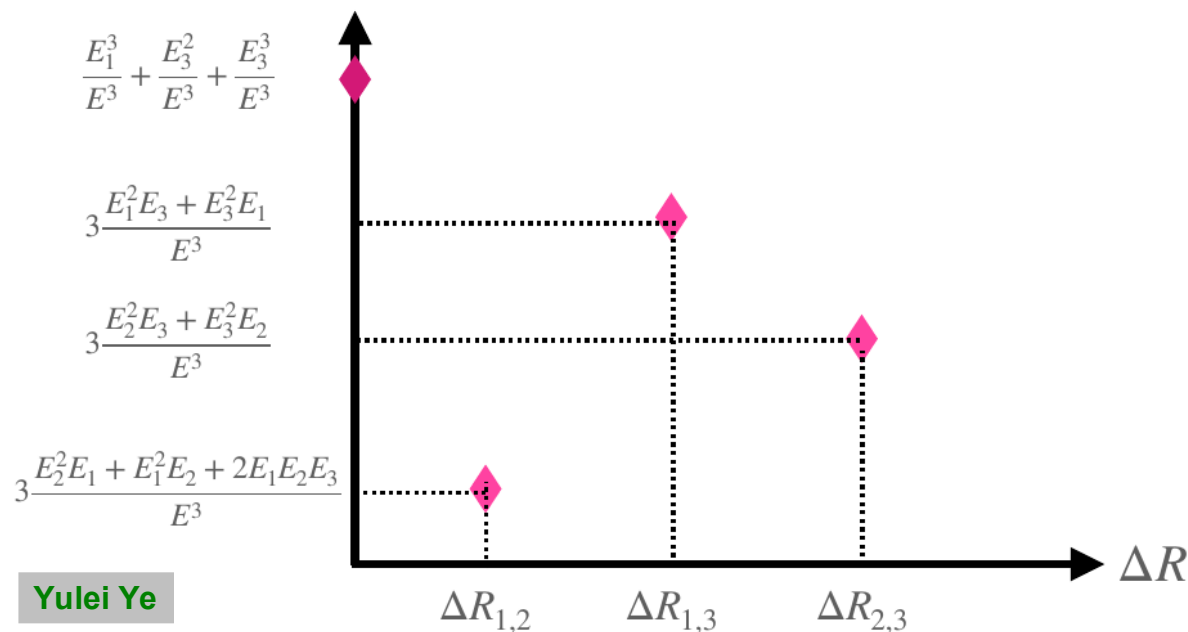
Jet substructure variable representing correlations of energy flow inside jets

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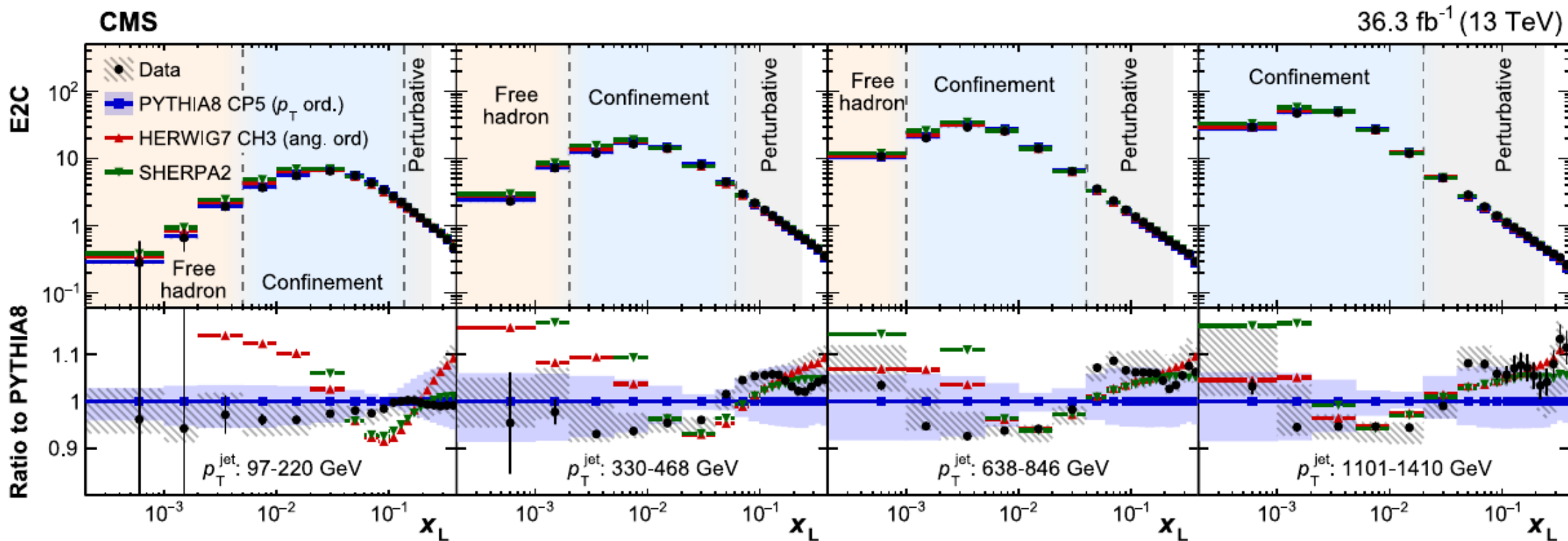
Chen et al., arXiv:2004.11381;
Lee et al., arXiv:2205.0314;
Chen et al., arXiv:2307.07510.



- 2016 data comprising 36.3 fb^{-1}
 - Series of single jet triggers with lowest p_T threshold 60 GeV
 - Dijet event selection with:
 - ➔ Anti-kT jets $R=0.4$
 - ➔ $p_T > 97 \text{ GeV}$, $|\eta| < 2.1$ → good momentum resolution
 - ➔ $|\Delta\Phi| > 2$ → back-to-back jets
 - ➔ Only two leading jets used
 - All particle candidates with $p_T > 1 \text{ GeV}$ used for E2C & E3C
 - Iterative unfolding with early stopping (D'Agostini)
 - ➔ Jet matching efficiency 99%
 - ➔ For particle candidates more problematic
- largest uncertainty on EEC from MC modelling!

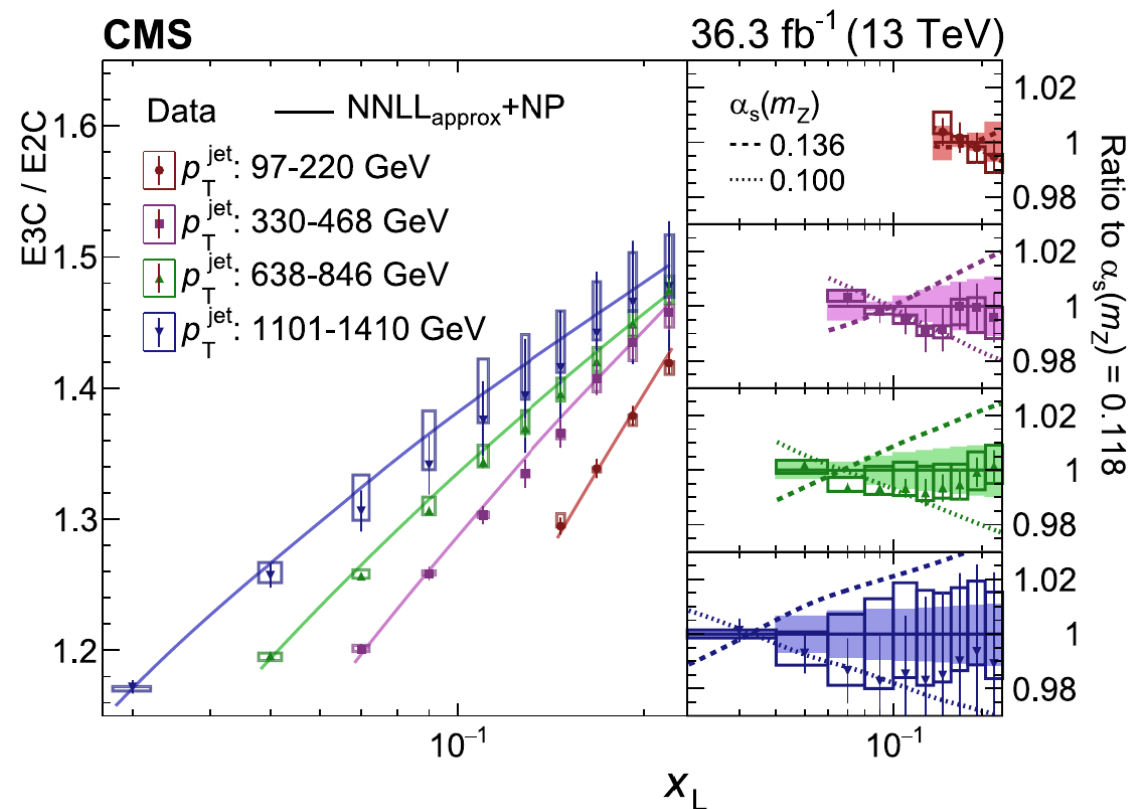
As ΔR goes smaller: perturbative region $\rightarrow$ confinement $\rightarrow$ free hadron

Transition depends on jet p_T bin, dashed lines move to the left with p_T up

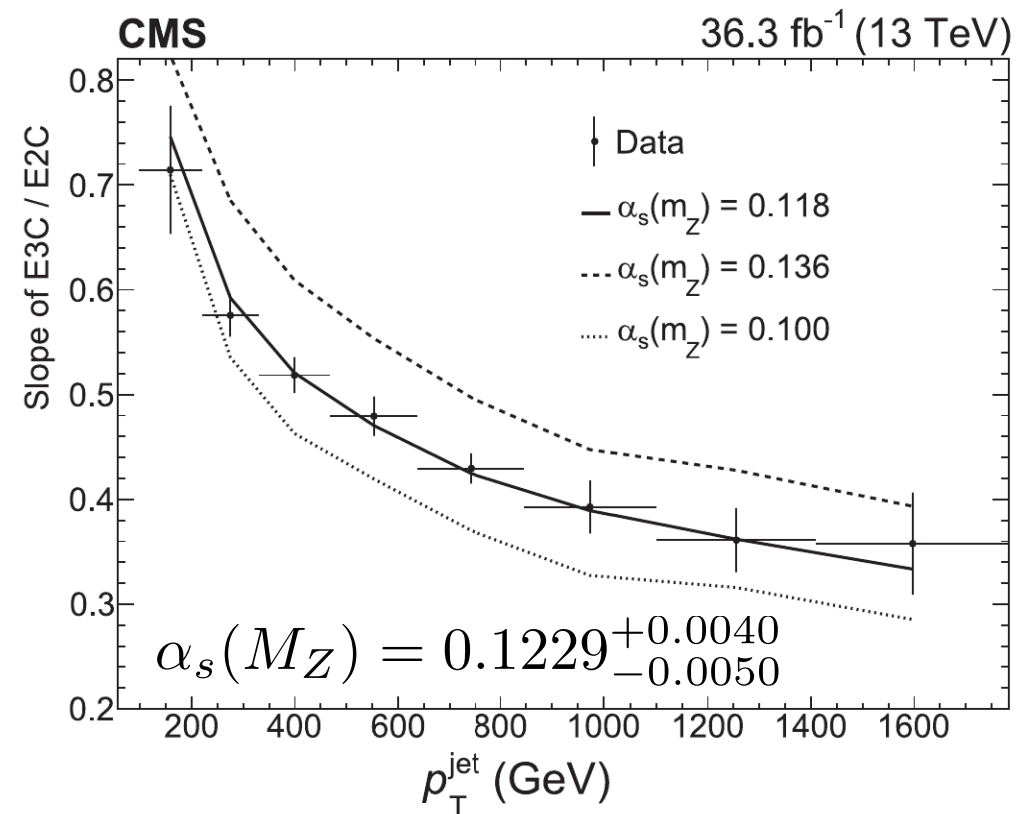
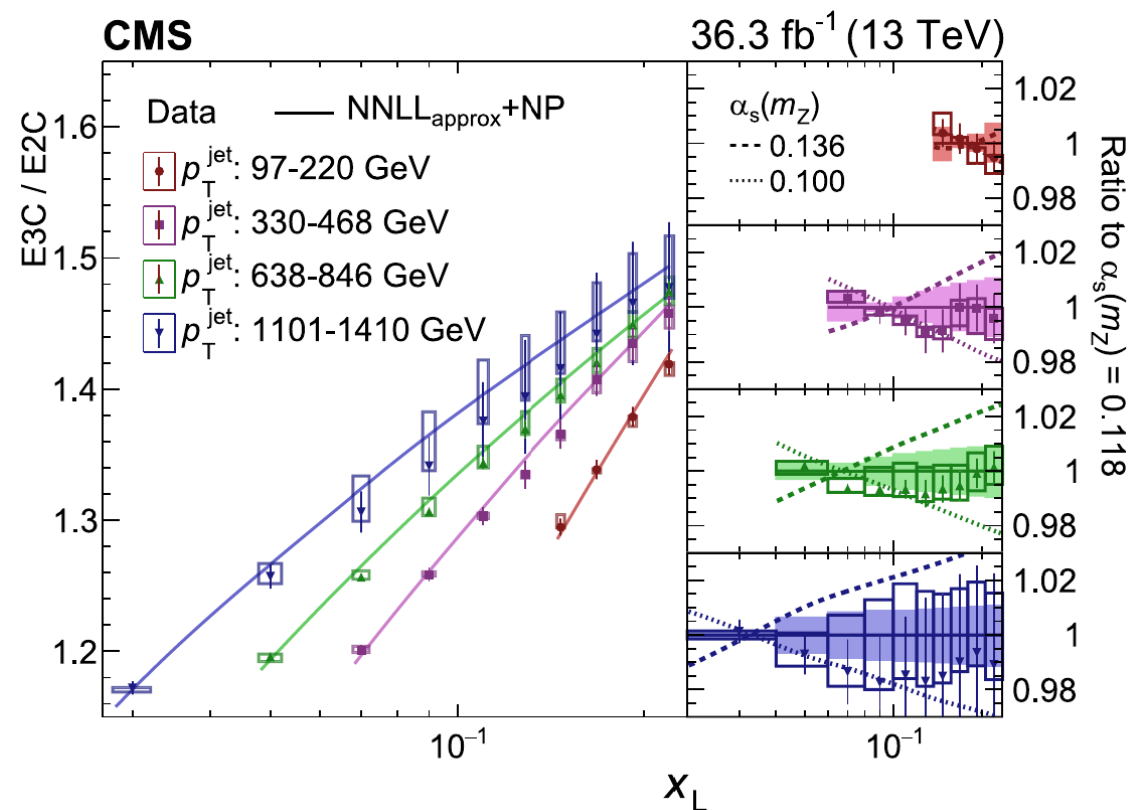


CMS, PRL 133 (2024) 071903.

- Focus on perturbative region:
 - Comparison to $\text{NLO} + \text{NNLL}_{\text{approx}}$
 - Use ratio $\text{E2C} / \text{E3C} \rightarrow \propto \alpha_s(Q) \ln R + \mathcal{O}(\alpha_s^2)$
 - Effects of e.g. gluon vs. quark jet expected to cancel $\rightarrow \alpha_s(Q)$ (NLO)



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Summary & Outlook

- LHC at 7, 8, and 13 TeV enabled to test $\alpha_s(Q)$ up to $Q \sim 7$ TeV
- LHC results reached $\Delta\alpha_s(M_Z) \sim 0.5\%$ experimentally
- LHC theory uncertainty still leads to $\Delta\alpha_s(M_Z) \sim 1.5\%$ in total (mostly)
- Still more theory understanding required
- Novel ideas like energy-energy correlators in jets (or Z p_T) very promising, need some more experience

Thank you for your attention!

**Thank you very much to the
organisers for the invitation to
this very special place**

