

Physics at the TeV Scale (aka Introduction to the Terascale)

Roger Wolf 20. April 2015

INSTITUTE OF EXPERIMENTAL PARTICLE PHYSICS (IEKP) – PHYSICS FACULTY

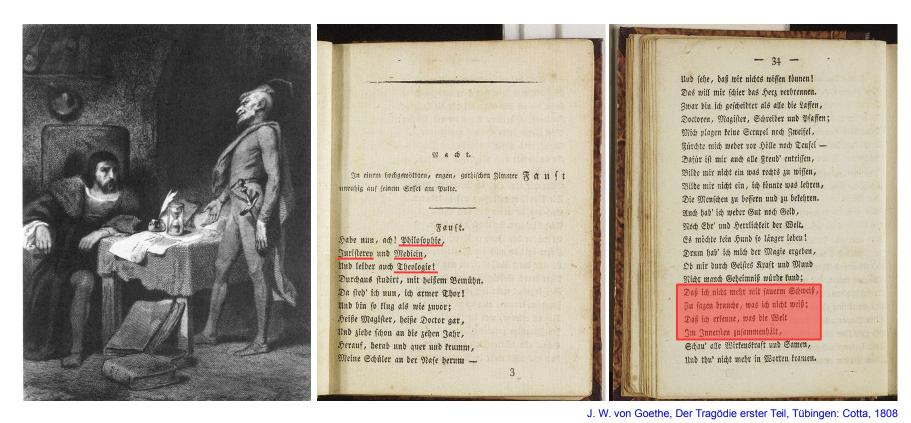


KIT – University of the State of Baden-Wuerttemberg and National Research Center of the Helmholtz Association

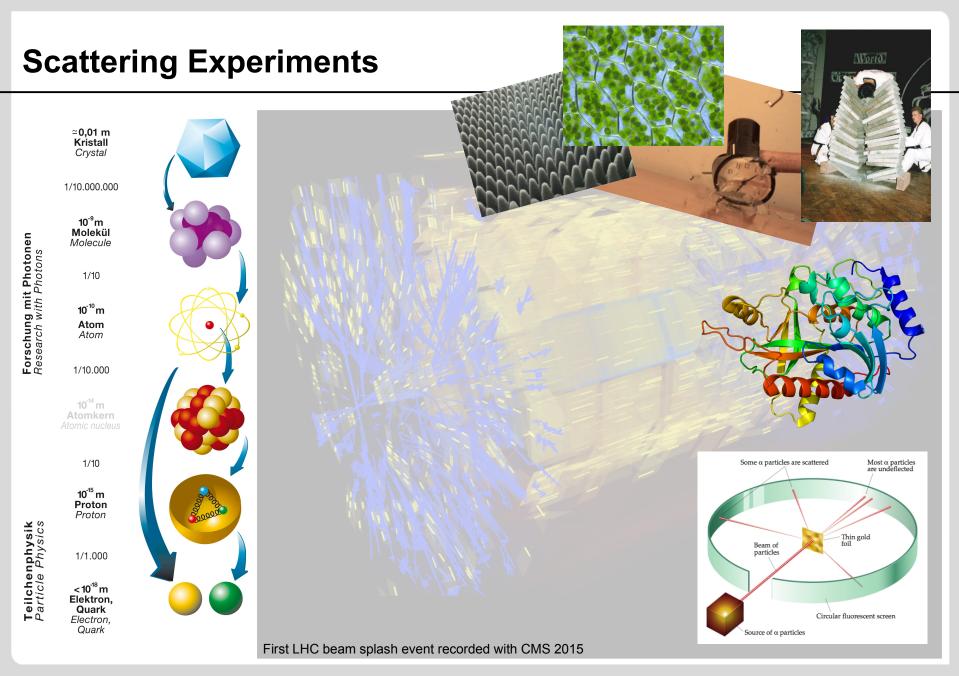
www.kit.edu

Nobel Aims

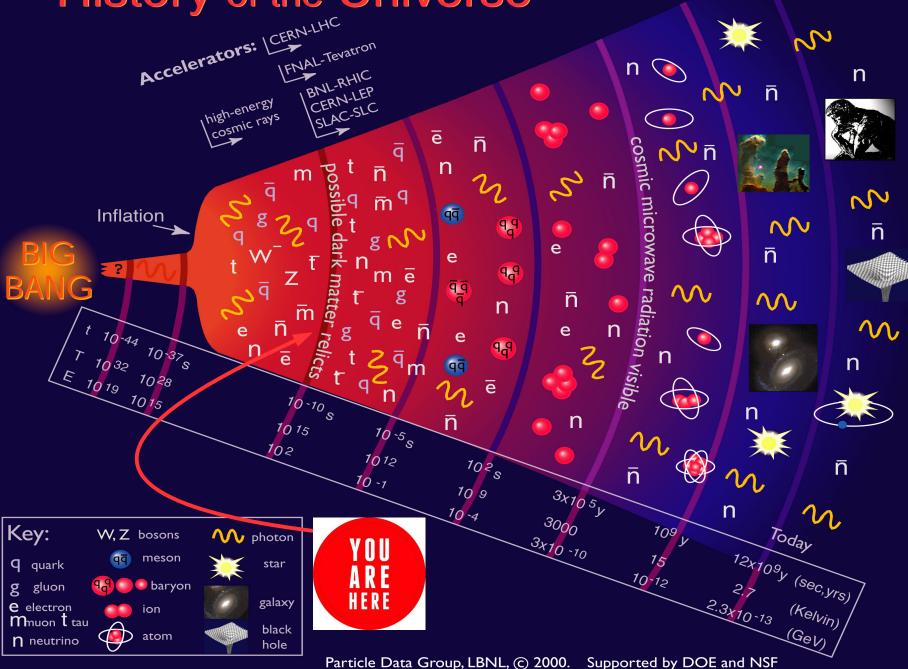


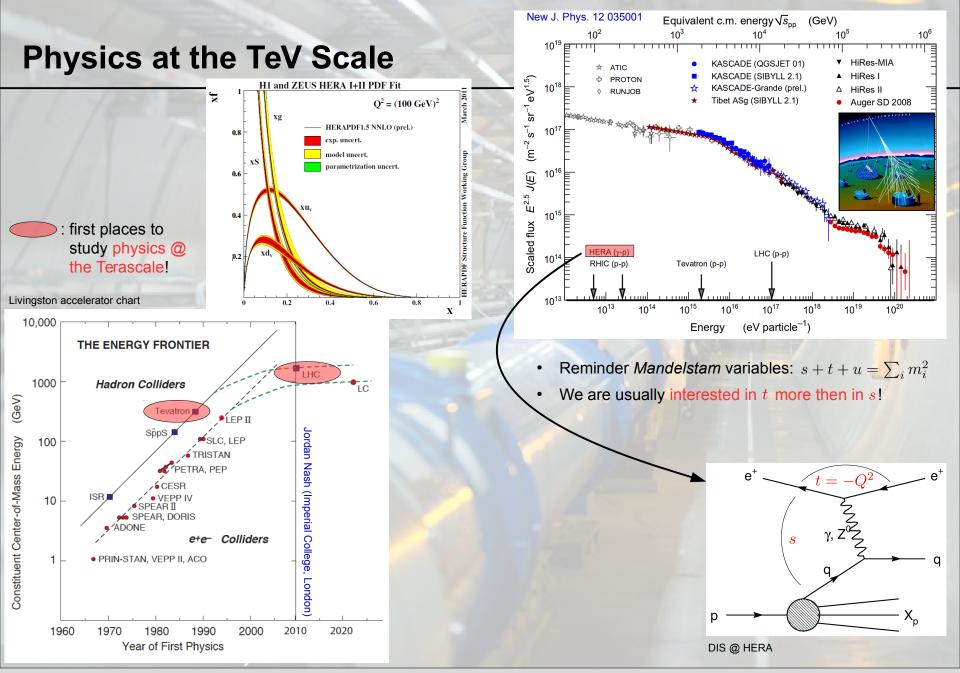


- 207 years later.
- Same questions.
- More success oriented ansatz.

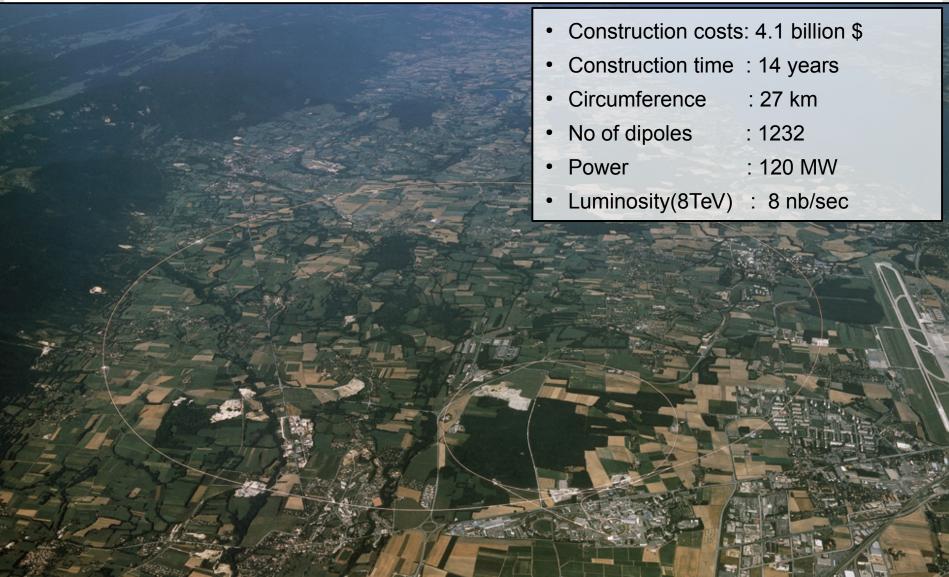


History of the **Universe**











Energy radiated off per rotation cycle:

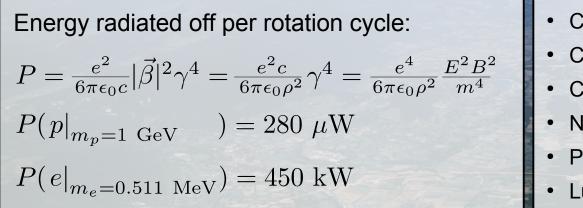
$$P = \frac{e^2}{6\pi\epsilon_0 c} |\vec{\beta}|^2 \gamma^4 = \frac{e^2 c}{6\pi\epsilon_0 \rho^2} \gamma^4 = \frac{e^4}{6\pi\epsilon_0 \rho^2} \frac{E^2 B^2}{m^4}$$
$$P(p|_{m_p=1 \text{ GeV}}) = 280 \ \mu\text{W}$$

$$P(e|_{m_e=0.511 \text{ MeV}}) = 450 \text{ kW}$$

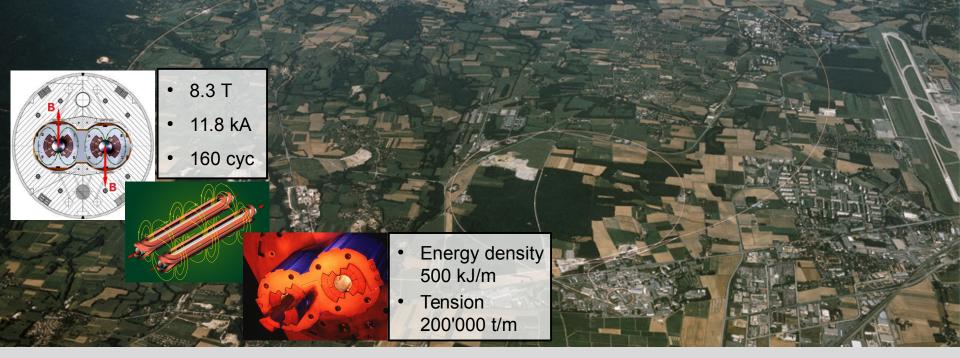
- Construction costs: 4.1 billion \$
- Construction time : 14 years
- Circumference : 27 km
- No of dipoles : 1232
- Power : 120 MW
- Luminosity(8TeV) : 8 nb/sec



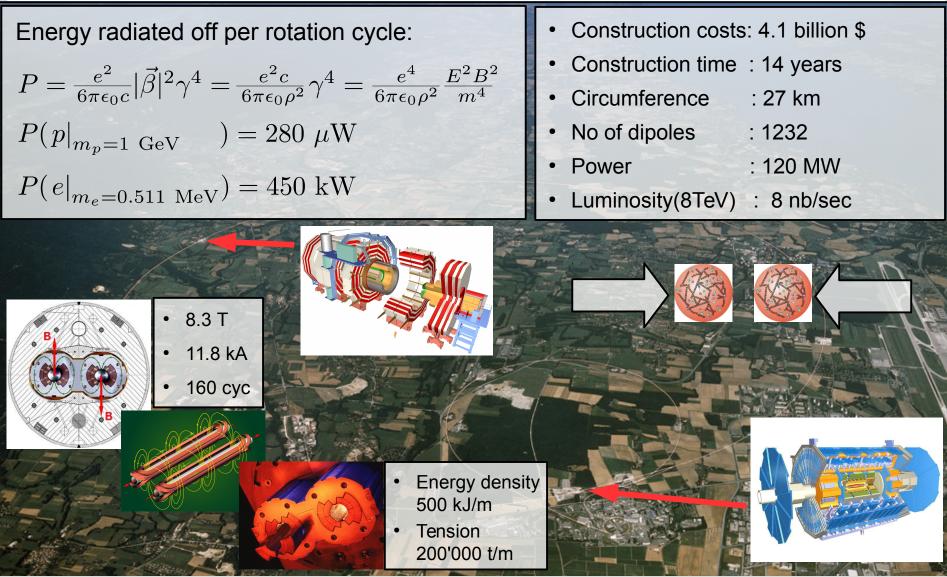




- Construction costs: 4.1 billion \$
- Construction time : 14 years
- Circumference : 27 km
- No of dipoles : 1232
- Power : 120 MW
- Luminosity(8TeV) : 8 nb/sec

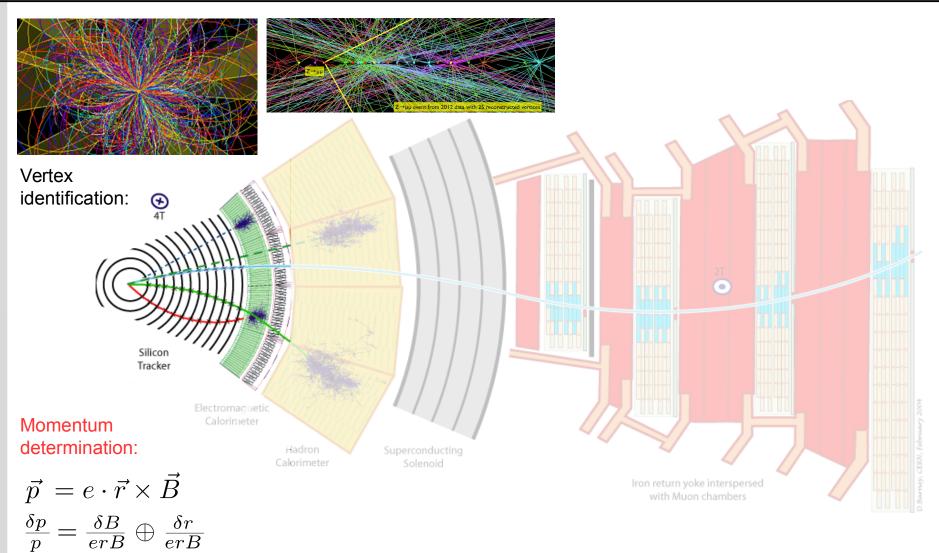






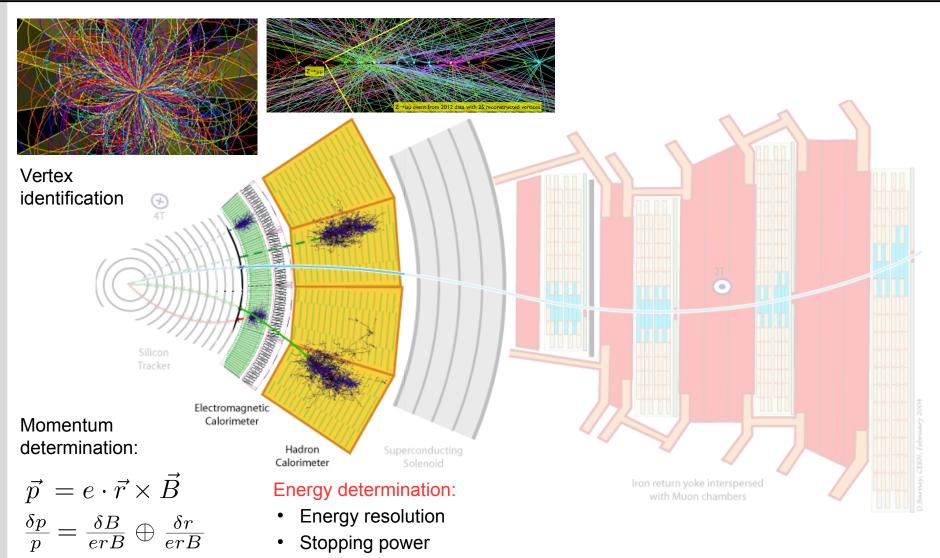
Key Demands on Experiments





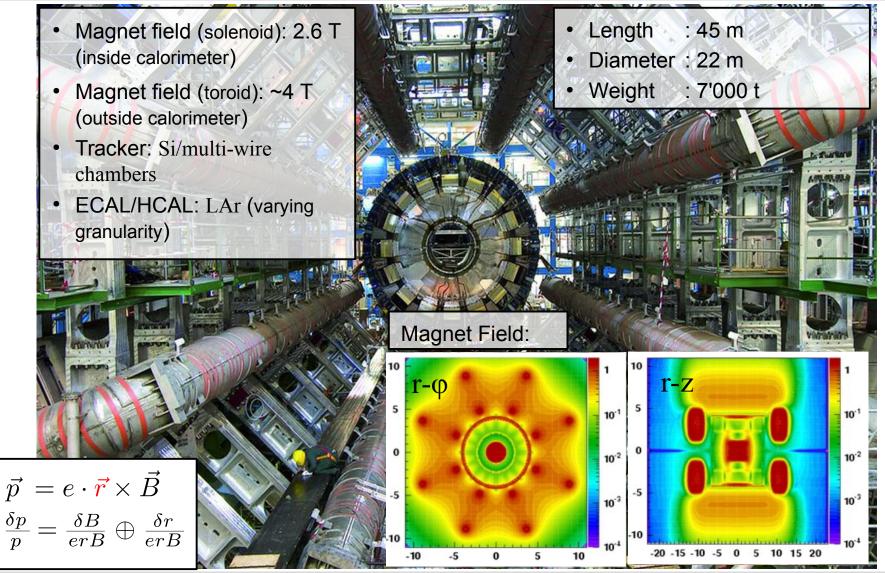
Key Demands on Experiments





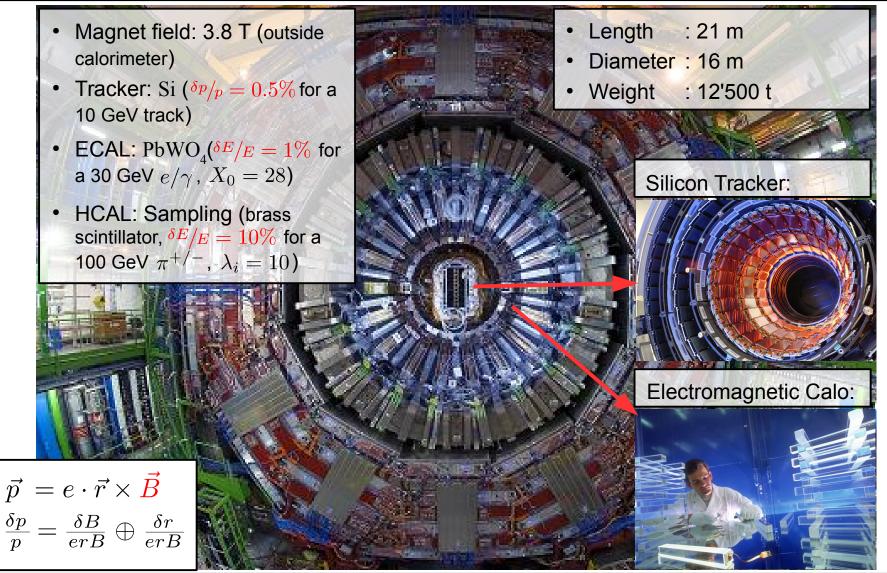
The Large Scale Solution (ATLAS)





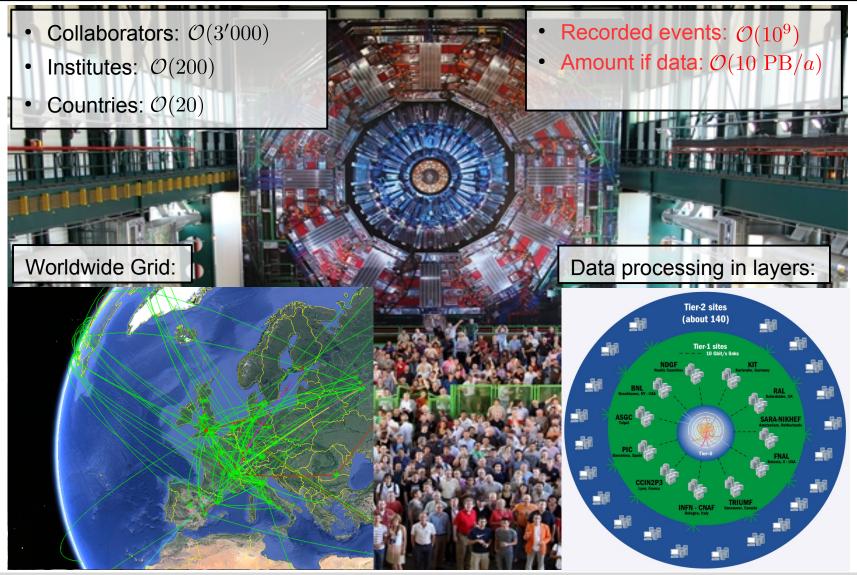
The Compact Solution (CMS)



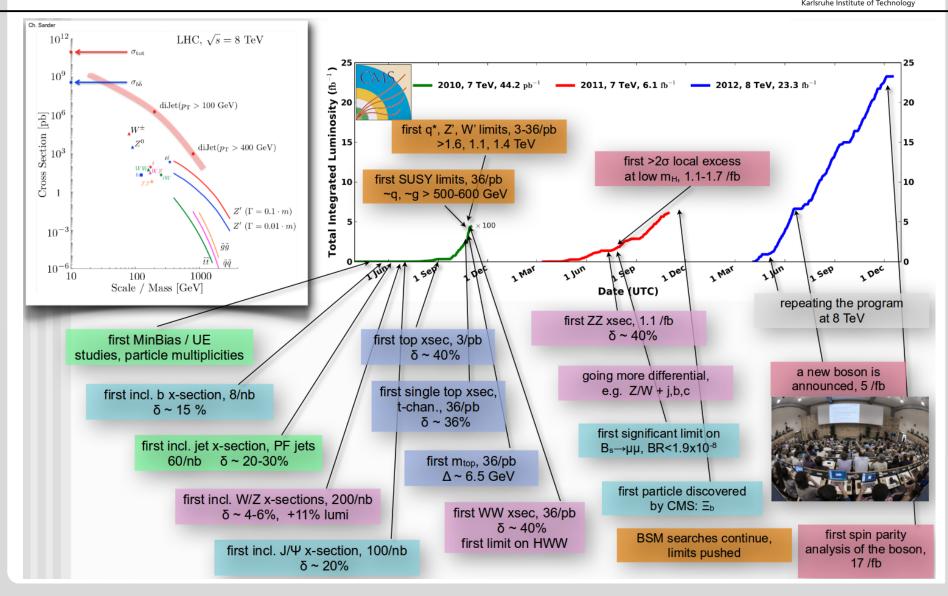


Worldwide Distribution of Data



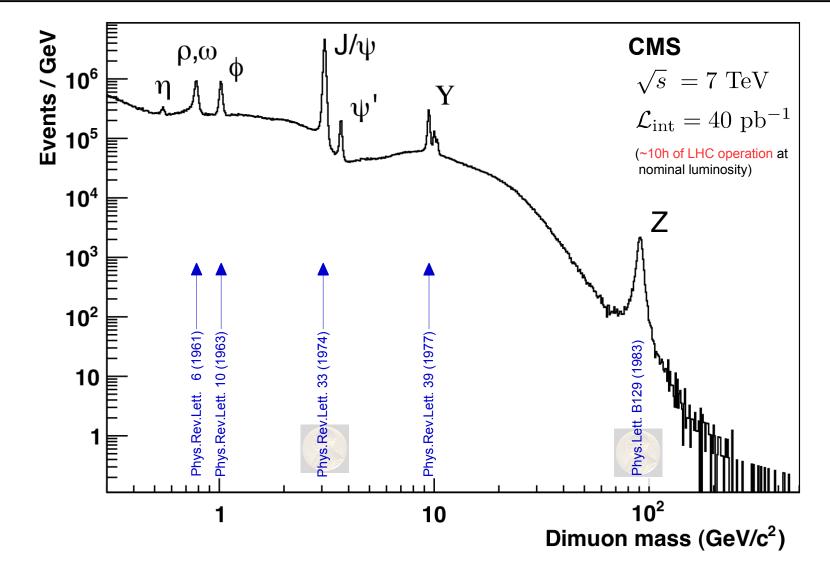


LHC History (represented in physics measurements)

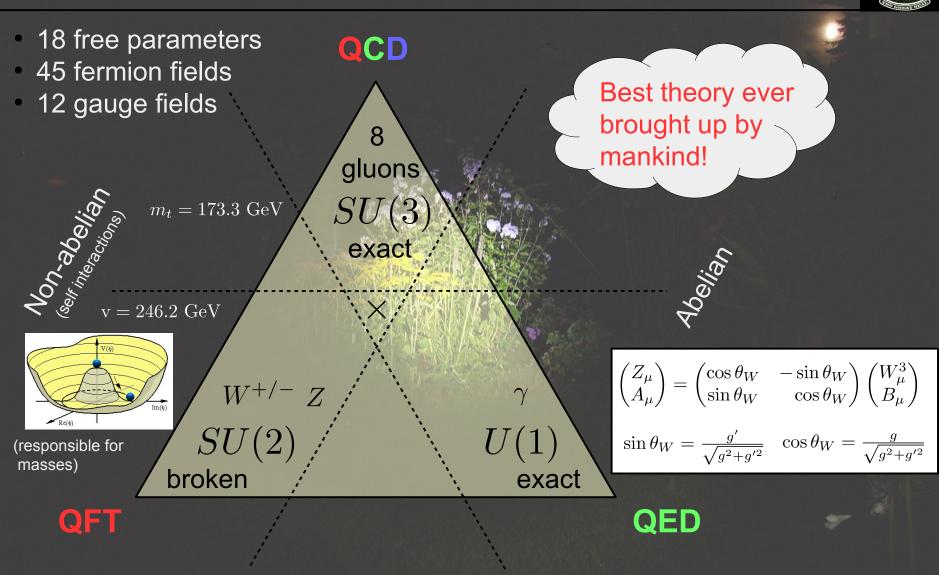


LHC Repeating History (in fast forward)



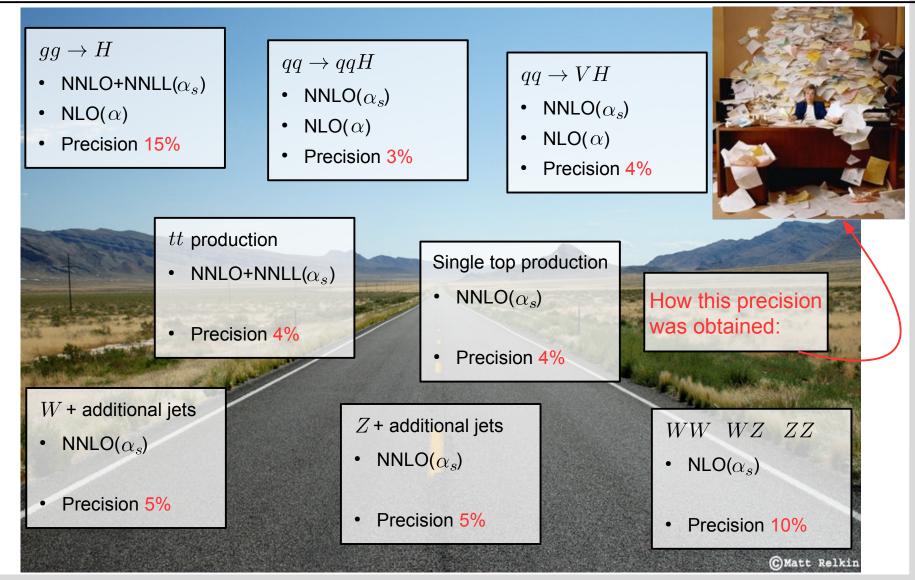


The Standard Model of Particle Physics (SM)



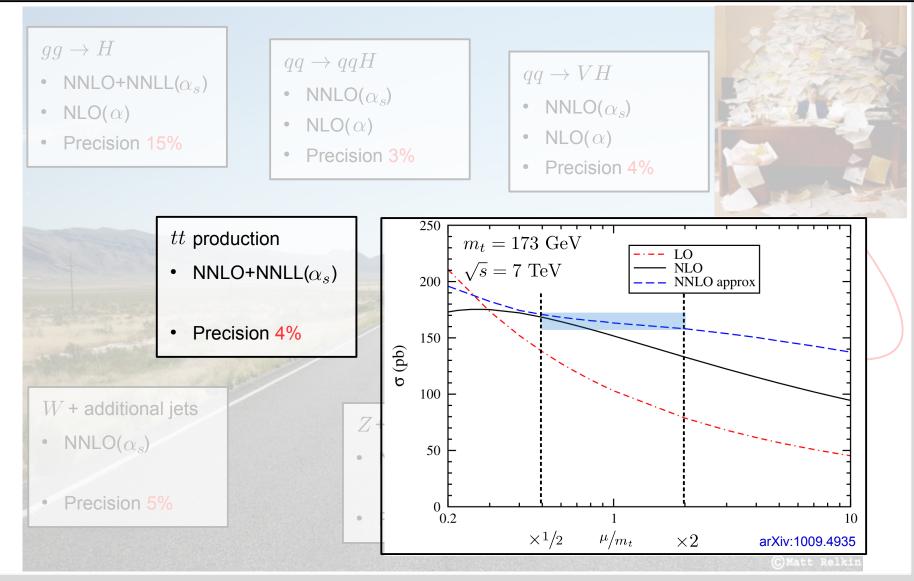
A Long Road of Theory Developments





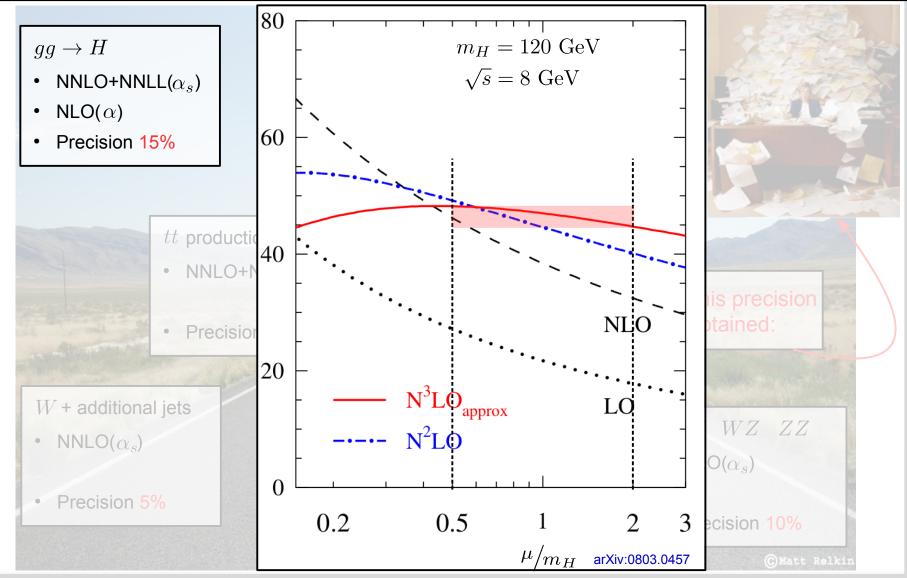
A Long Road of Theory Developments





A Long Road of Theory Developments



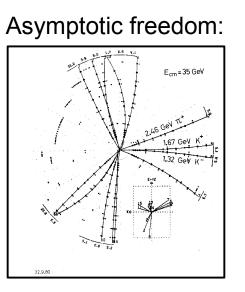


The Strong Sector of the SM

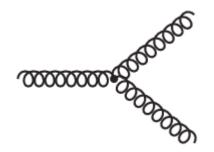


SU(3) exact gauge symmetry:

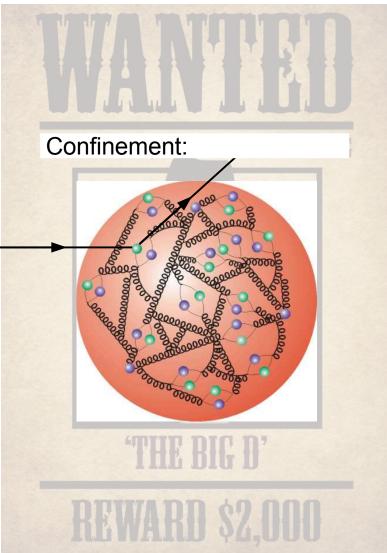
 \rightarrow 8 massless gluons as gauge bosons.

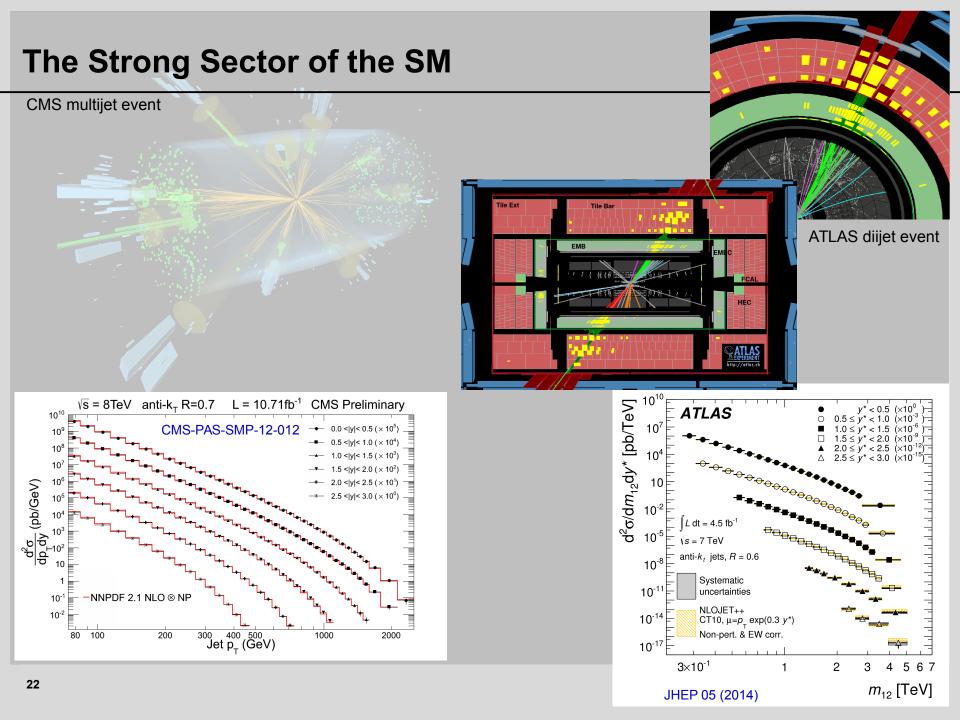


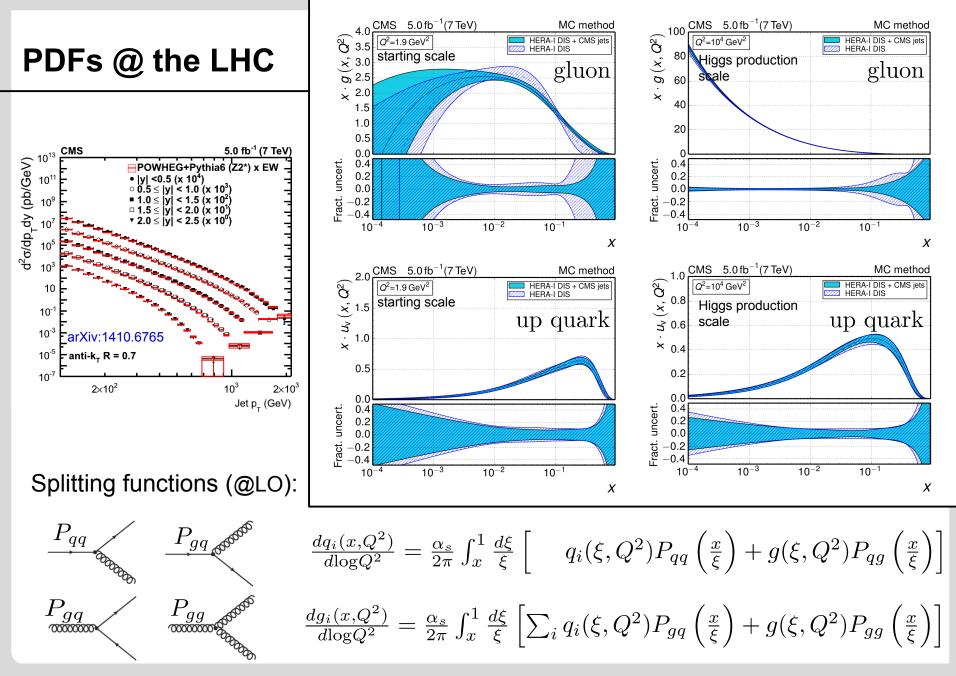
Non-Abelian gauge structure:



 \rightarrow leads to gluon-gluon self-couplings.





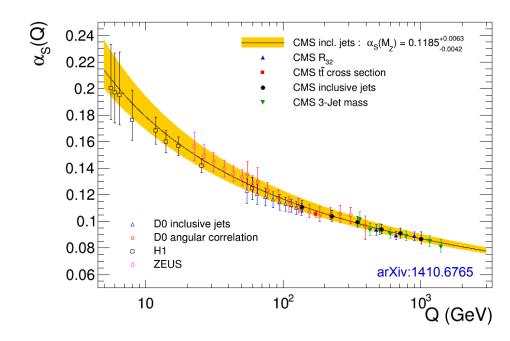


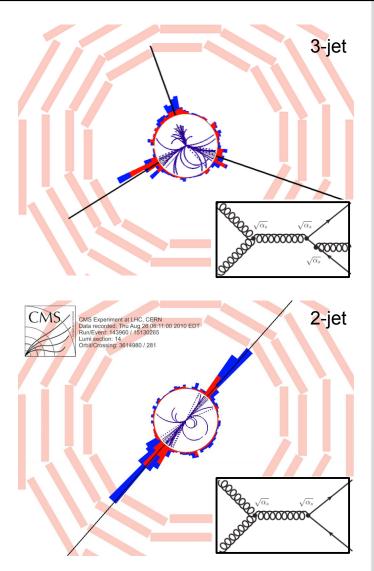
The Running of α_s



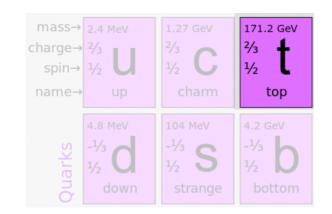
Running $\alpha_s(\mu^2)$ (e.g. from 1-loop RGE): $\alpha_s(\mu^2) = \frac{g_s^2(\mu^2)}{4\pi} = \frac{1}{\beta_0 \ln(\mu^2/\Lambda_{\rm QCD}^2)}$

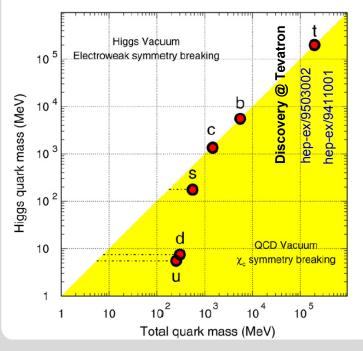
Measure $\alpha_s(\mu^2)$ from ratio of 3-jet over 2-jet events at given scale:



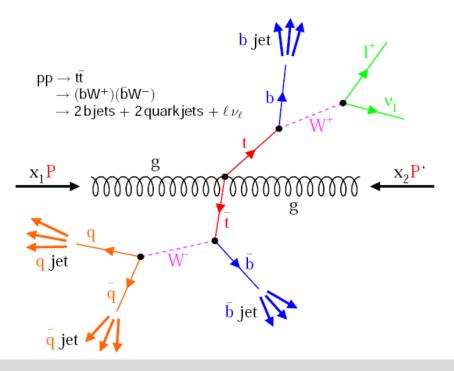






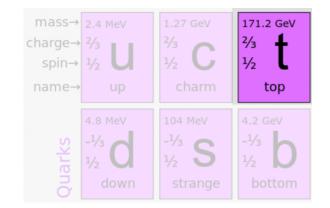


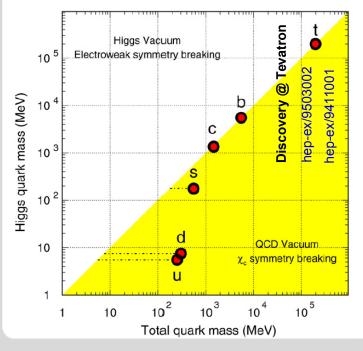
- Weak decay nearly exclusively via $t \rightarrow b$ ($BR \approx 0.998$).
- Lifetime $t = 5 \times 10^{-25}$ s (hadronization time scales $\delta t \approx \frac{\hbar}{\Lambda_{QCD}} \approx 3 \times 10^{-21}$ s).



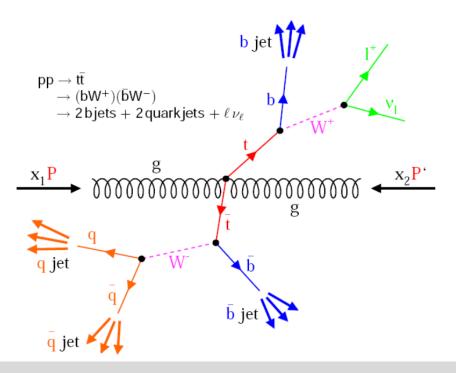
Physics of the Top Quark





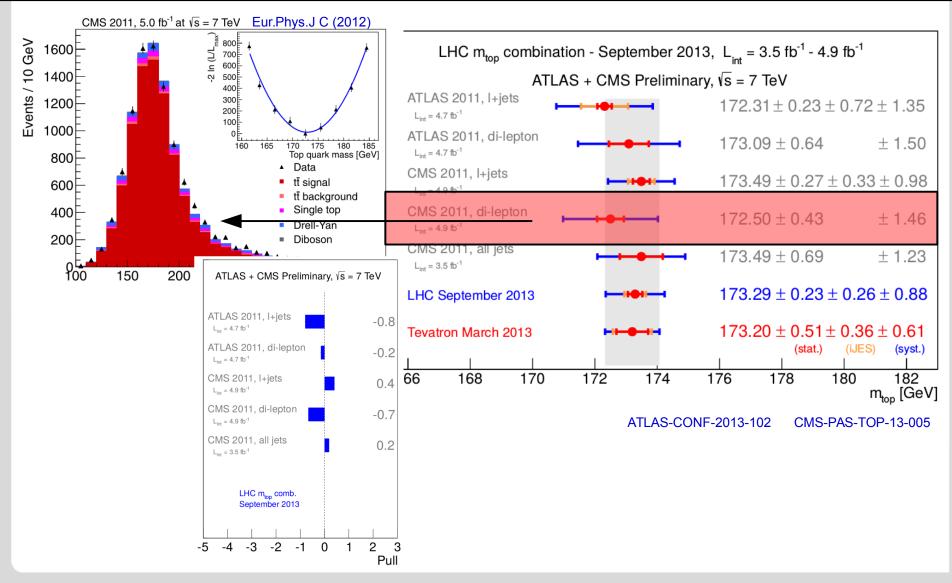


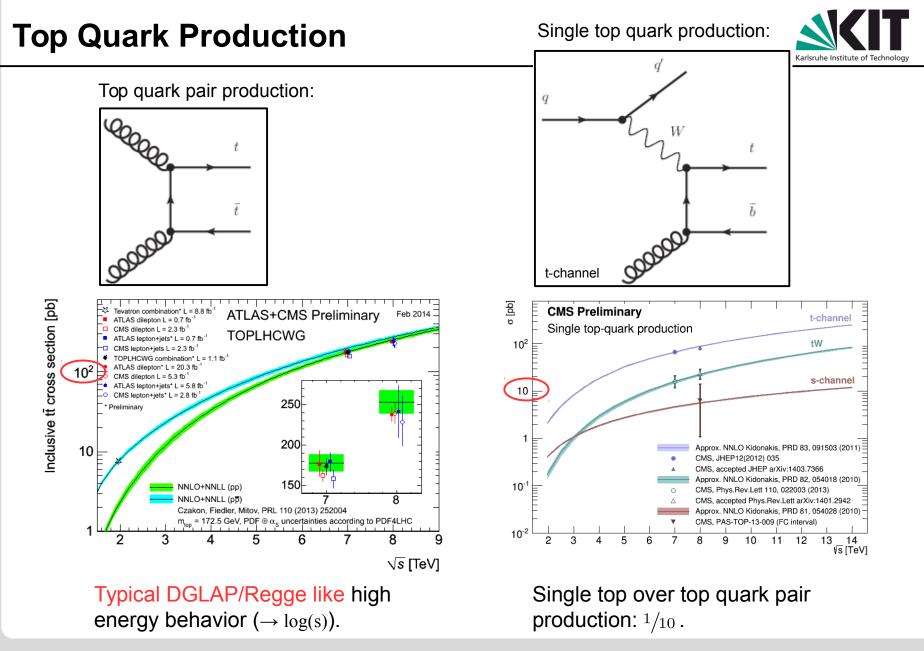
- Weak decay nearly exclusively via $t \rightarrow b$ ($BR \approx 0.998$).
- Lifetime $t = 5 \times 10^{-25}$ s (hadronization time scales $\delta t \approx \frac{\hbar}{\Lambda_{QCD}} \approx 3 \times 10^{-21}$ s).
- Fermi's Golden Rule: $\lambda_{if} = \frac{2\pi}{\hbar} \left| \mathcal{M}_{if} \right|^2 \cdot \rho(E_f)$



The Mass of the Top Quark

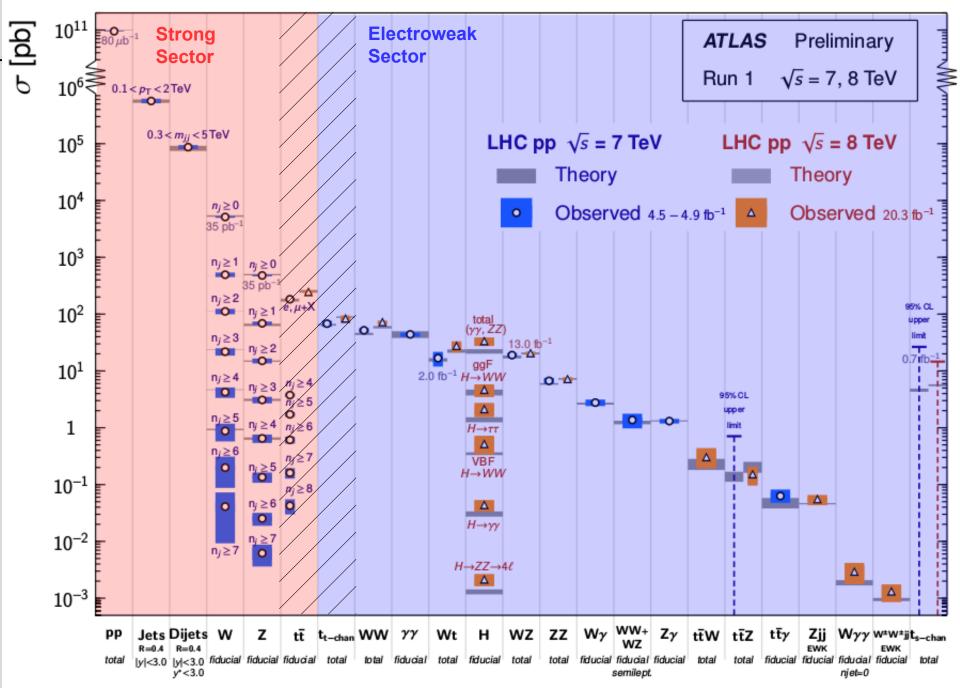


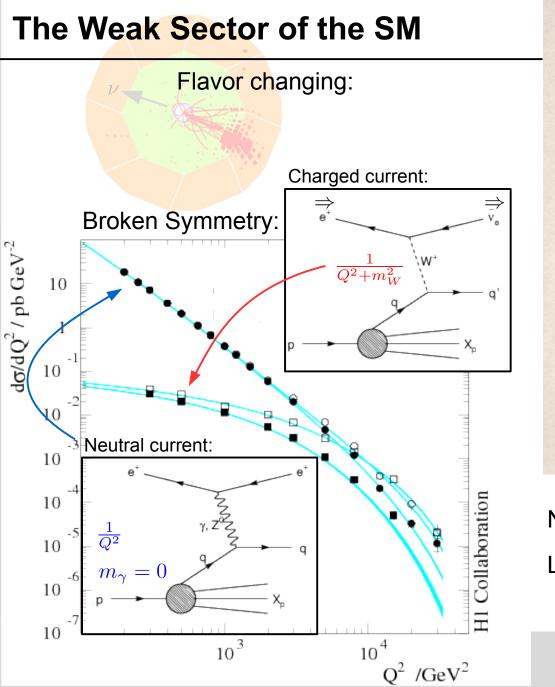


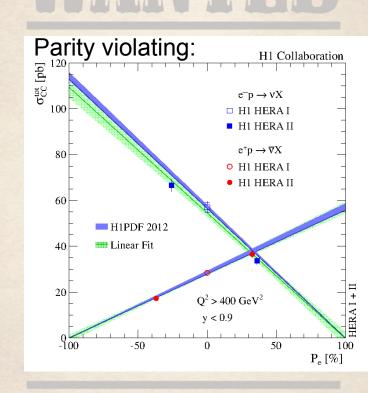


Standard Model Production Cross Section Measurements

Status: March 2015



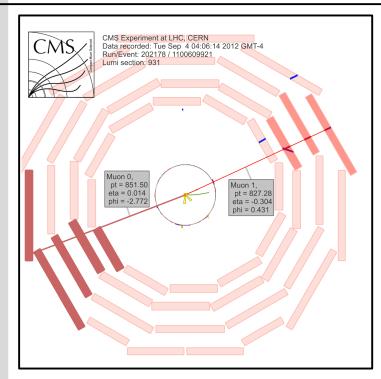




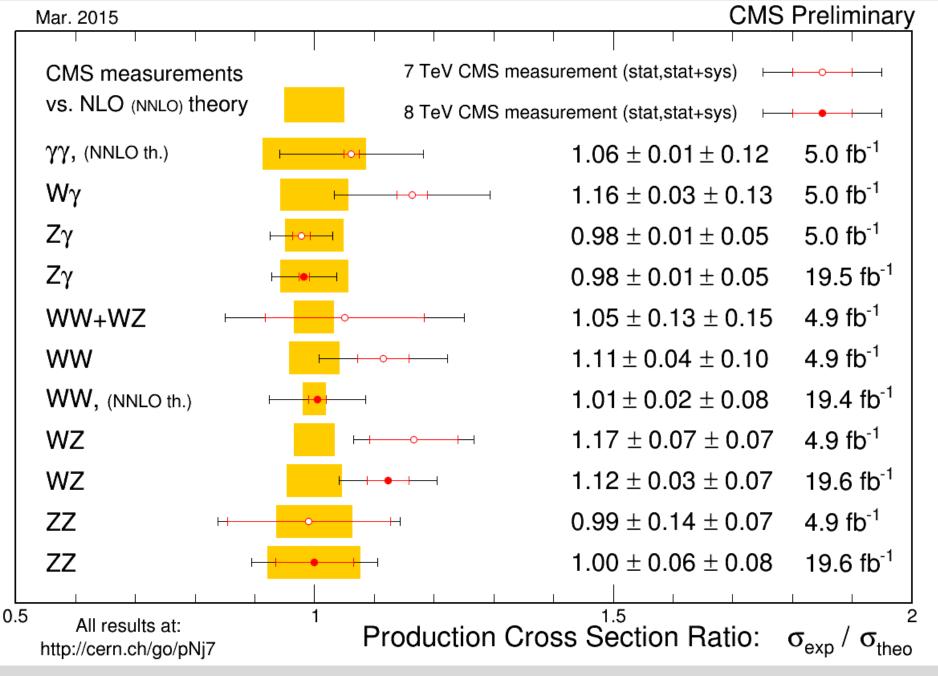
Non-Abelian gauge structure! Lepton universality!

Institute of Experimental Particle Physics (IEKP)

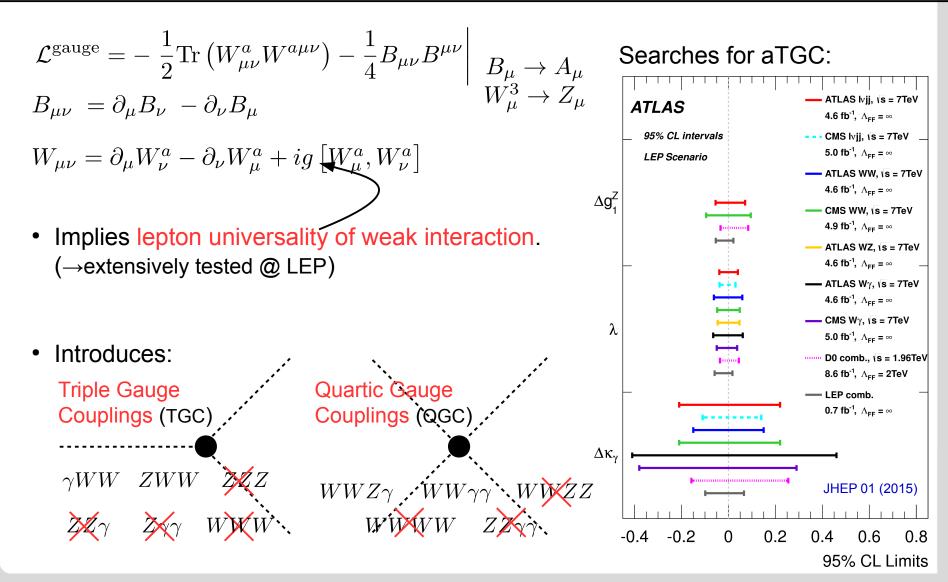
Z Bosons (abused) as Calibration Tools



- Tag & Probe to validate ℓ reconstr. efficiencies:
- Tag:
 - Strictly selected, isolated ℓ .
 - Well defined vertex with m_Z requirement.
- Probe:
 - Non-isolated ℓ .
 - Non-ID'ed ℓ .
 - A track pure track or cluster (validate linking efficiency).
- Tool to validate MET resolution (from recoil of $Z \rightarrow \mu\mu$): $\vec{U} + \vec{\mu}_1 + \vec{\mu}_2 = \vec{0}$
- Tool to validate efficiency to find hard interaction vertex.





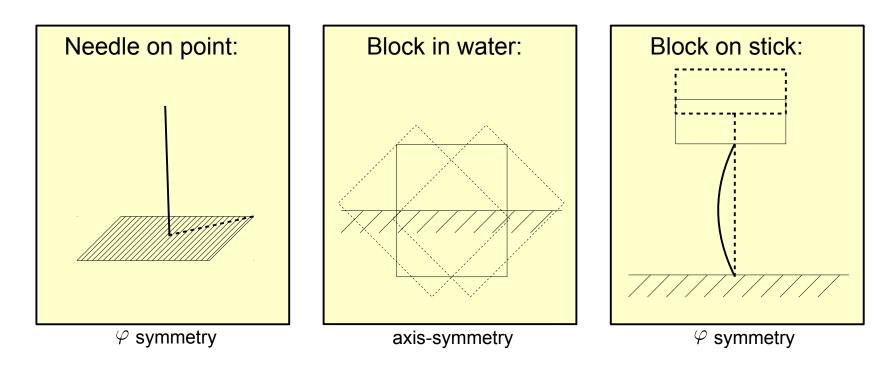


The Case of Electroweak Symmetry

Local gauge symmetries strictly require Weak interactions are described force mediating particle to have m = 0: by weak gauge symmetries! \rightarrow symmetry exists. Fermions Bosons Force mediating particles Quarks t explicitly break symmetry! \rightarrow U С symmetry not realized in nature. charm up top photon d b S Ζ down strange bottom Z boson Weak interaction makes a Leptons V_{τ} V_{P} difference between left- & right-W boson electron muon tau handed coordinate systems. neutrino neutrino neutrino auе и This property destroys local electron muon tau gluon gauge invariance for all weak interactions if fermions have $m_Z = 91.1876 \pm 0.0021 \text{ GeV}$ mass $m \neq 0$. $m_W = 85.385 \pm 0.015 \text{ GeV}$

Spontaneous Symmetry Breaking

- Symmetry present in the system (i.e. in Lagrangian density \mathcal{L}).
- BUT symmetry broken in energy ground state of the system (i.e. quantum vacuum).
- Three examples from classical mechanics:



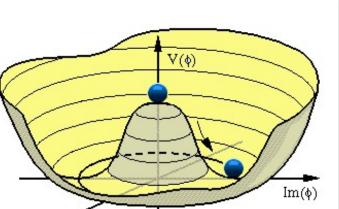
Institute of Experimental Particle Physics (IEKP)



The Higgs Mechanism

- Introduce new field ϕ with characteristic interaction potential.
- Leads to prediction of new particle: → Higgs boson!
- Allows to incorporate mass terms in the theory.
- Gauge symmetry compromising mass terms compensated by characteristic couplings to Higgs particle:

$$\kappa_V = \frac{2m_v^2}{v}$$
 (for force mediating *W* & *Z* boson).
 $\kappa_f = \frac{m_f}{v}$ (for weakly interacting fermions).

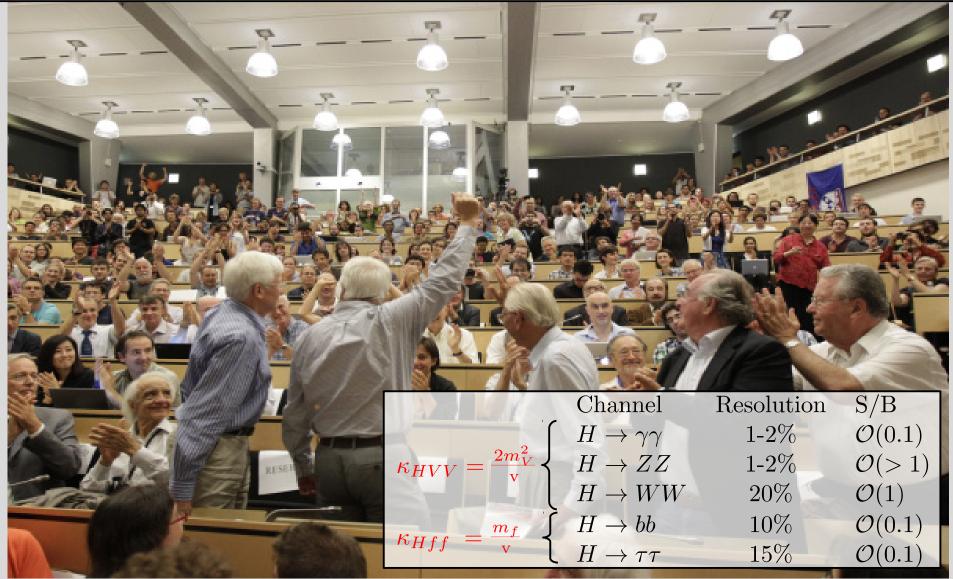


Re())



The Discovery of a New Particle 4th July 2012

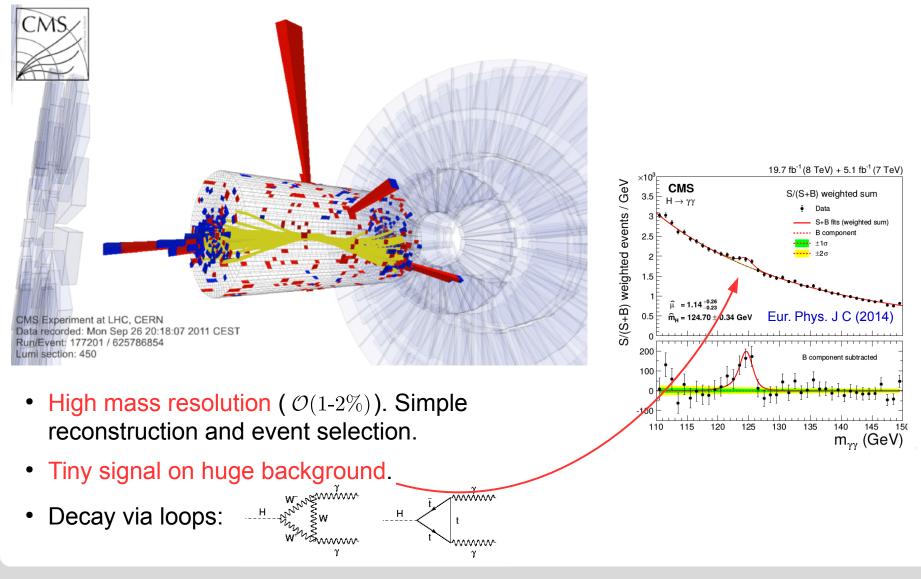




Institute of Experimental Particle Physics (IEKP)

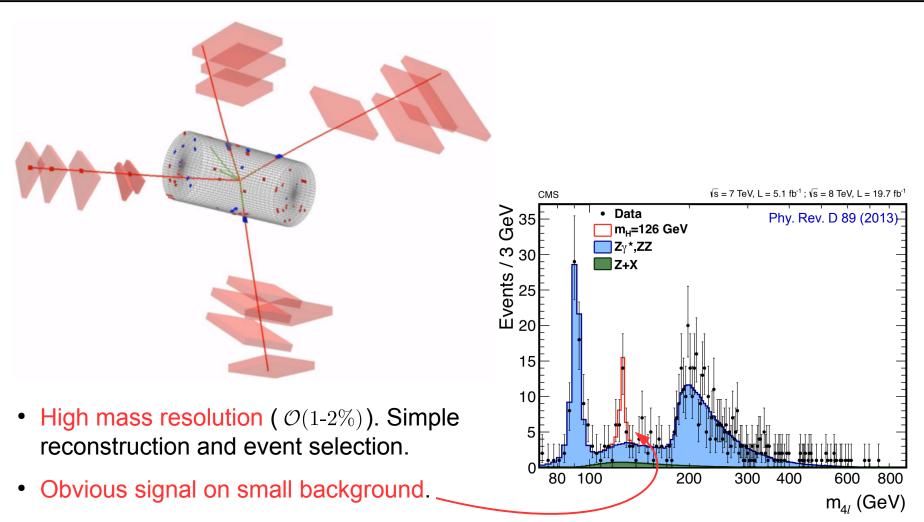
$H \rightarrow \gamma \gamma \,\,$ Decay Channel





$H \rightarrow ZZ$ Decay Channel





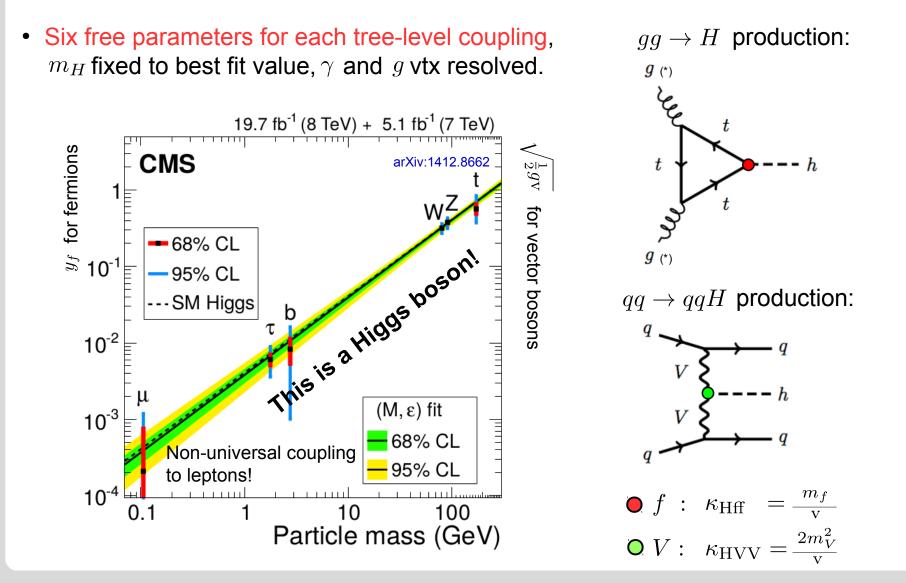
• Most important search channels: $4\mu \ 2\mu 2e \ 4e$

| | | 10 fb Significan | | Significar | 25 fb nce (σ) | $\hat{\mu}$ |
|---------------------|--|--|--------------|---|--------------------------|---|
| ATLAS | $ \begin{array}{c} H \to \gamma \gamma \\ H \to ZZ \\ H \to WW \end{array} $ | $ \begin{array}{r} 4.5 (2.5) \\ 3.6 (2.7) \\ 2.8 (2.3) \end{array} $ | [25] [25] | $5.2 (4.6) \\ 8.1 (6.2) \\ 6.1 (5.8)$ | [97] | 1.17 ± 0.27 1.44 ± 0.37 1.09 ± 0.21 |
| CMS | $ \begin{array}{c} H \rightarrow \gamma \gamma \\ H \rightarrow ZZ \\ H \rightarrow WW \end{array} $ | $\begin{array}{c} 4.1 \ (2.8) \\ 3.2 \ (3.8) \\ 1.6 \ (2.5) \end{array}$ | | 5.7 (5.2) 6.8 (6.7) 4.3 (5.8) | [89] [90] [91] | $\begin{array}{rrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrr$ |
| | | $10 { m ~fb^{-1}}$ | 1 | 2 | 25 fb^{-1} | |
| | | Significanc | | Significan | | $\hat{\mu}$ |
| AS | $H \to \tau \tau$ | _ | | 4.5(3.4) | [113] | 1.43 ± 0.40 |
| ATLAS | $H 	o b\bar{b}$ $t\bar{t}H(\gamma\gamma)$ | _ | | 1.4(2.6) – | $[114] \\ [115]$ | $ \begin{array}{r} 0.52 \pm 0.40 \\ 1.30 \pm 2.20 \end{array} $ |
| \mathbf{IS} | $H \to \tau \tau$ | · · · · | [27] | 3.2 (3.7) | [92] | $0.78~\pm~0.27$ |
| CMS | $H \to bb$ $t\bar{t}H$ | 0.7(1.9) | [27] | $\begin{array}{c} 2.1 \ (2.5) \\ 3.4 \ (1.2) \end{array}$ | [93] [111] | $ \begin{array}{r} 0.84 \pm 0.44 \\ 2.80 \pm 0.95 \end{array} $ |

Taken from "The Higgs Boson Discovery at the Large Hadron Collider" Springer Tract of Modern Physics (tbp soon).

Karlsruhe Institute of Technology

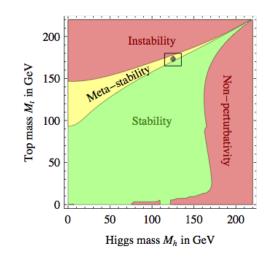


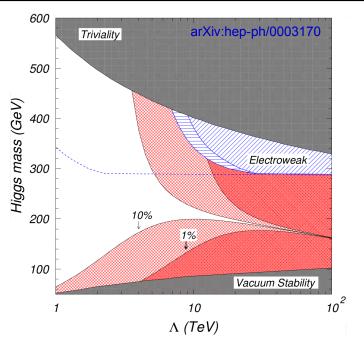


Why it is not THE Higgs Boson (of the SM)⁽¹⁾



- Gravity is not included in the SM.
- The SM suffers from the hierarchy problem.
- Dark matter is not included in the SM.
- Neutrino masses are not included in the SM.
- There are known deviations in $a_{\mu} \equiv \frac{g_{\mu}-2}{2}$ from the SM expectation (3.6 σ unresolved).



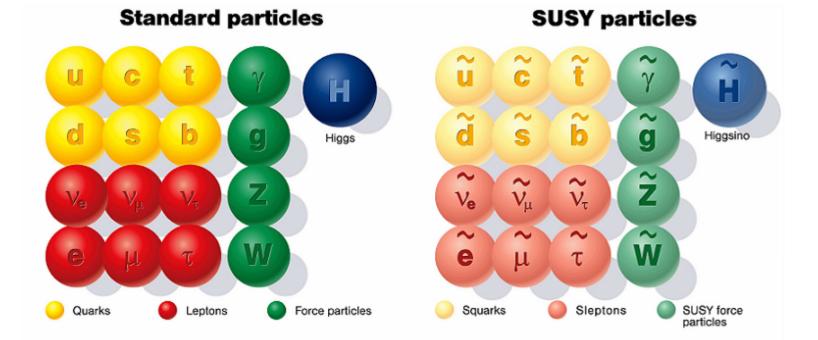


- There must be physics beyond the SM!
- At what scale does it set in?
- (How) Does it influence the Higgs sector?

SUSY Extension of the SM



 Extension of SM by a last remaining, non-trivial, symmetry operation (boson ↔ fermion), SUSY, can cure many shortcomings of SM:



- E.g. lightest SUSY particle (LSP) perfect candidate for DM .
- Problem: SUSY itself is broken!



• Five neutral Higgs bosons predicted:

$$H_1 = \begin{pmatrix} H_1^0 \\ H_1^- \end{pmatrix}$$
, $Y_{H_1} = -1$, v_1 : VEV₁

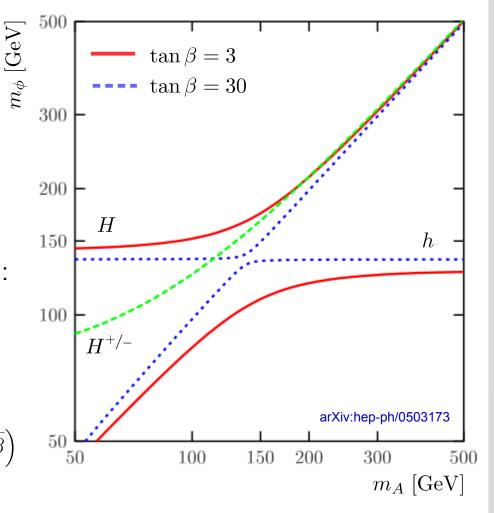
$$H_2 = \begin{pmatrix} H_2^+ \\ H_2^0 \end{pmatrix}, \quad Y_{H_2} = +1, \quad \mathbf{v}_2: \ \mathbf{VEV}_2$$

$$N_{\text{ndof}} = 8 \quad -3 = 5$$
$$W, Z \quad H^{+/-}, H, h, A$$

• MSSM mass requirements at tree level:

two free parameters: m_A , $\tan\beta = {
m v_1}/{
m v_2}$

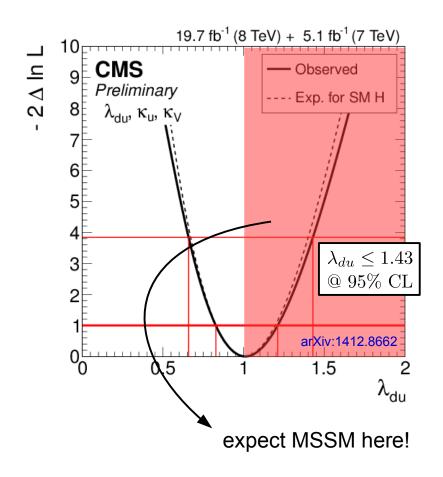
$$\begin{split} m_{H^{+/-}}^2 &= m_A^2 + m_W^2 \\ m_{H, h}^2 &= \frac{1}{2} \left(m_A^2 + m_Z^2 \pm \sqrt{\left(m_A^2 + m_Z^2\right)^2} \right) \\ & - \frac{1}{2} \left(m_A^2 + m_Z^2 \pm \sqrt{\left(m_A^2 + m_Z^2\right)^2} \right) \\ & - \frac{1}{2} \left(m_A^2 + m_Z^2 \pm \sqrt{\left(m_A^2 + m_Z^2\right)^2} \right) \\ & - \frac{1}{2} \left(m_A^2 + m_Z^2 \pm \sqrt{\left(m_A^2 + m_Z^2\right)^2} \right) \\ & - \frac{1}{2} \left(m_A^2 + m_Z^2 \pm \sqrt{\left(m_A^2 + m_Z^2\right)^2} \right) \\ & - \frac{1}{2} \left(m_A^2 + m_Z^2 \pm \sqrt{\left(m_A^2 + m_Z^2\right)^2} \right) \\ & - \frac{1}{2} \left(m_A^2 + m_Z^2 \pm \sqrt{\left(m_A^2 + m_Z^2\right)^2} \right) \\ & - \frac{1}{2} \left(m_A^2 + m_Z^2 \pm \sqrt{\left(m_A^2 + m_Z^2\right)^2} \right) \\ & - \frac{1}{2} \left(m_A^2 + m_Z^2 \pm \sqrt{\left(m_A^2 + m_Z^2\right)^2} \right) \\ & - \frac{1}{2} \left(m_A^2 + m_Z^2 \pm \sqrt{\left(m_A^2 + m_Z^2\right)^2} \right) \\ & - \frac{1}{2} \left(m_A^2 + m_Z^2 \pm \sqrt{\left(m_A^2 + m_Z^2\right)^2} \right) \\ & - \frac{1}{2} \left(m_A^2 + m_Z^2 \pm \sqrt{\left(m_A^2 + m_Z^2\right)^2} \right) \\ & - \frac{1}{2} \left(m_A^2 + m_Z^2 + m_Z^2 + m_Z^2 \right) \\ & - \frac{1}{2} \left(m_A^2 + m_Z^2 + m_Z^2 + m_Z^2 \right) \\ & - \frac{1}{2} \left(m_A^2 + m_Z^2 + m_Z^2 + m_Z^2 \right) \\ & - \frac{1}{2} \left(m_A^2 + m_Z^2 + m_Z^2 + m_Z^2 \right) \\ & - \frac{1}{2} \left(m_A^2 + m_Z^2 + m_Z^2 + m_Z^2 \right) \\ & - \frac{1}{2} \left(m_A^2 + m_Z^2 + m_Z^2 + m_Z^2 \right) \\ & - \frac{1}{2} \left(m_Z^2 + m_Z^2 + m_Z^2 + m_Z^2 + m_Z^2 \right) \\ & - \frac{1}{2} \left(m_Z^2 + m_Z^2 + m_Z^2 + m_Z^2 + m_Z^2 \right) \\ & - \frac{1}{2} \left(m_Z^2 + m_Z^2 + m_Z^2 + m_Z^2 + m_Z^2 \right) \\ & - \frac{1}{2} \left(m_Z^2 + m_Z^2 + m_Z^2 + m_Z^2 + m_Z^2 \right) \\ & - \frac{1}{2} \left(m_Z^2 + m_Z^2 + m_Z^2 + m_Z^2 + m_Z^2 \right) \\ & - \frac{1}{2} \left(m_Z^2 + m_Z^2 + m_Z^2 + m_Z^2 + m_Z^2 \right) \\ & - \frac{1}{2} \left(m_Z^2 + m_Z^2 + m_Z^2 + m_Z^2 + m_Z^2 + m_Z^2 + m_Z^2 \right) \\ & - \frac{1}{2} \left(m_Z^2 + m_$$



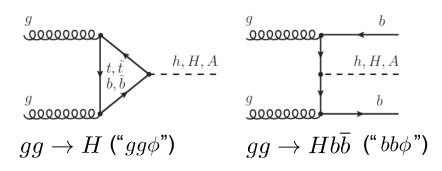
Enhancement of down-type Couplings



• In MSSM coupling to down-type fermions enhanced for $\tan \beta \gg 1$.



- Interesting decay channels:
 - H o au au ($\hat{\kappa}_{ au} = 0.84 \pm^{0.19}_{0.18}$)
 - $H \rightarrow bb$ ($\hat{\kappa}_b = 0.74 \pm 0.33_{0.29}$)
- Interesting production modes:



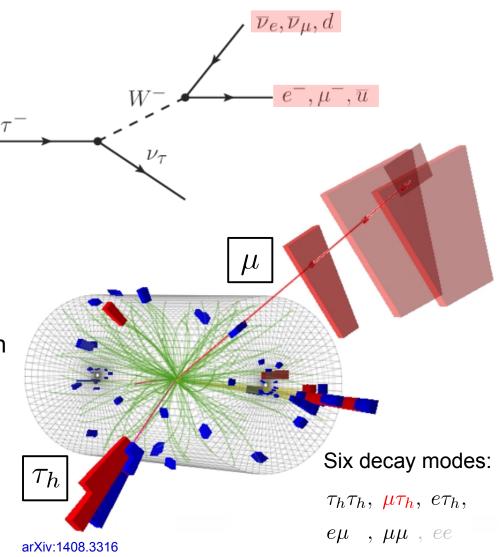


| Decay Mode | BR | |
|--|--------|-----------|
| $\tau \to e\nu_e\nu_\tau$ | 17.83% | all |
| $\tau \to \mu \nu_{\mu} \nu_{\tau}$ | 17.41% | , of a |
| $\tau \rightarrow 1$ -prong ν_{τ} | 37.10% | 20% |
| $\tau \rightarrow 3$ -prong ν_{τ} | 15.20% | 20 |

Search for 2 isolated high *p_T* leptons (*e*, μ, τ_h).

decay modes

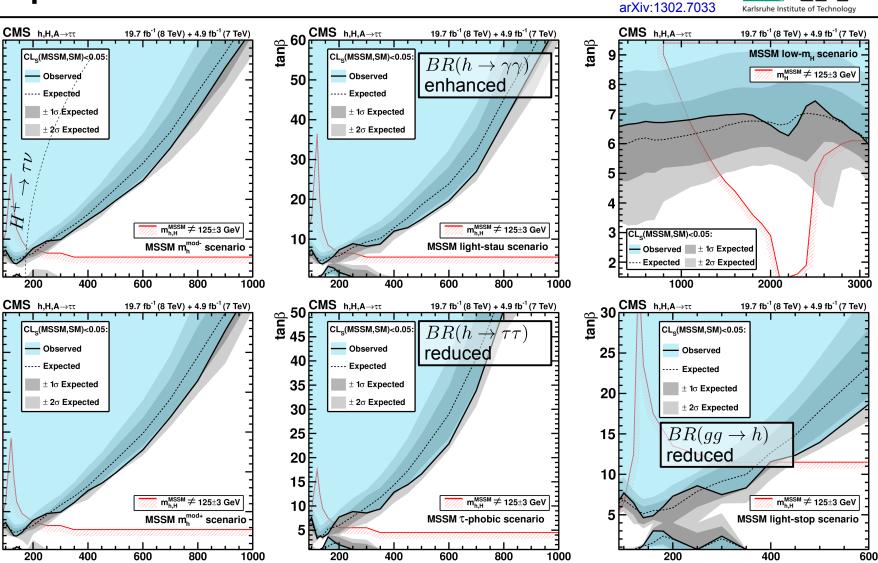
- Reduce obvious backgrounds (use on $\vec{E_T}$) & reconstruct $m_{ au au}$.
- Exploit characteristics of production mode to increase sensitivity.



Complete Set of Benchmark Scenarios

m₄ [GeV]





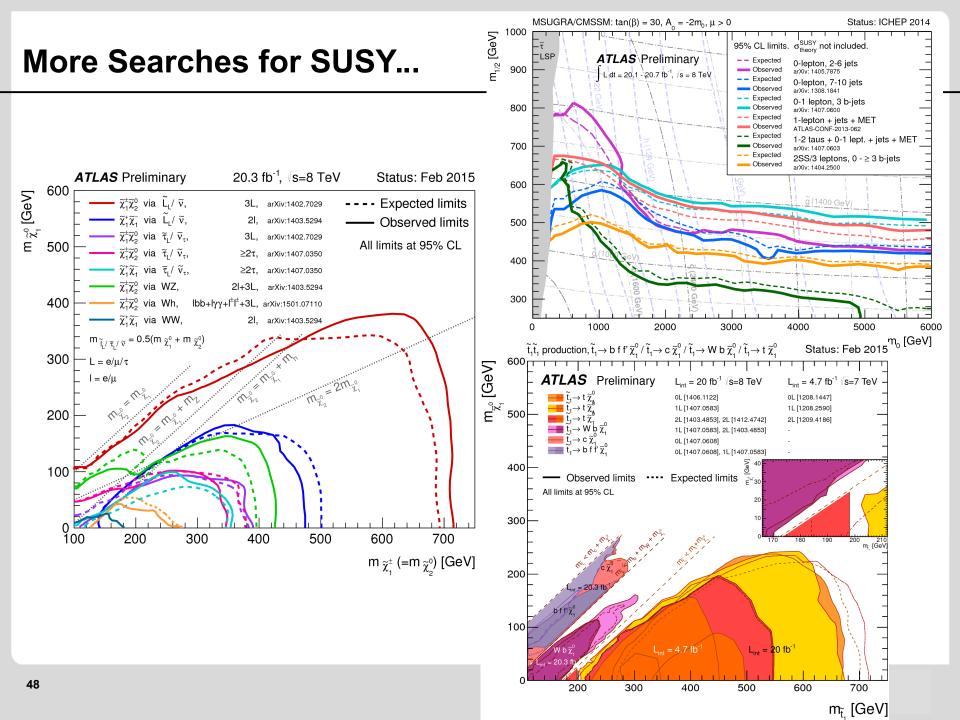
m_₄ [GeV]

m₄ [GeV]

,09 ,

tanβ 99

Institute of Experimental Particle Physics (IEKP)



ATLAS SUSY Searches* - 95% CL Lower Limits

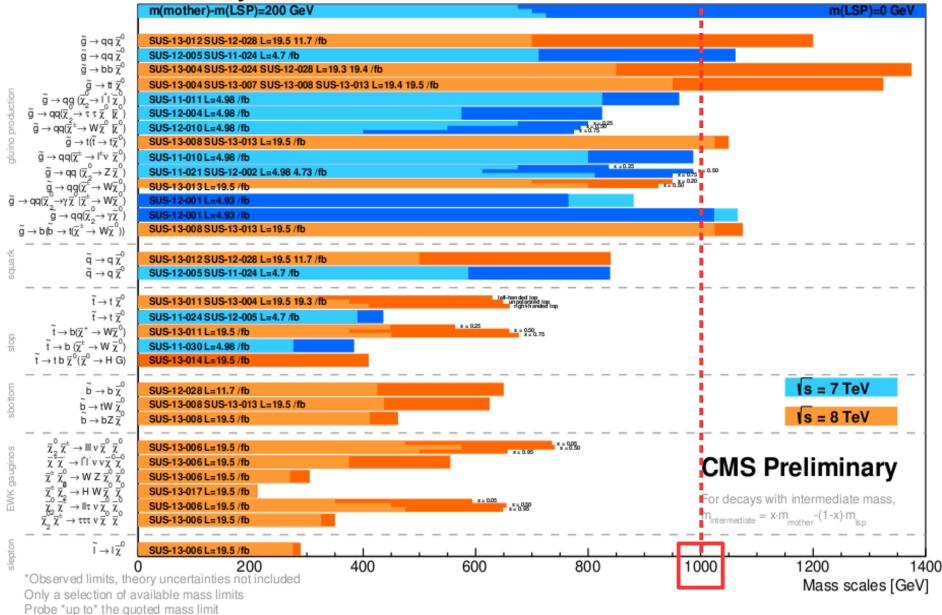
| | Model | e, μ, τ, γ | Jets | $E_{\rm T}^{\rm mass}$ | $\int \mathcal{L} dt [\text{fb}^-]$ | 1 Mass limit | Reference |
|---------------------------|--|--------------------------|--------------------------|------------------------|-------------------------------------|---|----------------------------|
| | MSUGRA/CMSSM | 0 | 2-6 jets | Yes | 20.3 | <i>q̃, g̃</i> 1.7 TeV m(<i>q̃</i>)=m(<i>g̃</i>) | 1405.7875 |
| Searches | $\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$ | 0 | 2-6 jets | Yes | 20.3 | \tilde{q} 850 GeV $m(\chi_1^0)=0$ GeV, $m(1^{st} \text{ gen}, \tilde{q})=m(2^{nd} \text{ gen}, \tilde{q})$ | 1405.7875 |
| | $\tilde{q}\tilde{q}\gamma, \tilde{q} \rightarrow q\tilde{\chi}_1^0$ (compressed) | 1γ | 0-1 jet | Yes | 20.3 | \tilde{q} 250 GeV $m(\tilde{q})-m(\tilde{\chi}_{1}^{0}) = m(c)$ | 1411.1559 |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\bar{q}\tilde{\chi}_{1}^{0}$ | 0 | 2-6 jets | Yes | 20.3 | \tilde{s} m(\tilde{x}_1^0)=0 GeV | 1405.7875 |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq\tilde{\chi}_1^{\pm} \rightarrow qqW^{\pm}\tilde{\chi}_1^0$ | 1 <i>e</i> , µ | 3-6 jets | Yes | 20 | \tilde{g} 1.2 TeV $m(\tilde{\chi}_1^0) < 300 \text{ GeV}, m(\tilde{\chi}_1^0) + m(\tilde{g}))$ | 1501.03555 |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq(\ell\ell/\ell\nu/\nu\nu)\tilde{\chi}_1^0$ | $2 e, \mu$ | 0-3 jets | - | 20 | \tilde{g} π_{32} TeV $m(\chi_1^0)=0$ GeV | 1501.03555 |
| | GMSB ($\tilde{\ell}$ NLSP) | 1-2 $	au$ + 0-1 ℓ | 0-2 jets | Yes | 20.3 | \tilde{g} 1.6 TeV $\tan\beta$ >20 | 1407.0603 |
| | GGM (bino NLSP) | 2γ | - | Yes | 20.3 | \tilde{s} 1.28 TeV $m(\tilde{\chi}_1^0)$ >50 GeV | ATLAS-CONF-2014 |
| | GGM (wino NLSP) | $1 e, \mu + \gamma$ | - | Yes | 4.8 | \tilde{s} 619 GeV m (χ_1^0) >50 GeV | ATLAS-CONF-2012 |
| | GGM (higgsino-bino NLSP) | γ | 1 <i>b</i> | Yes | 4.8 | ĝ 900 GeV m(λ ⁰ ₁)>220 GeV | 1211.1167 |
| | GGM (higgsino NLSP) | 2 e, µ (Z) | 0-3 jets | Yes | 5.8 | ĝ 690 GeV m(NLSP)>200 GeV | ATLAS-CONF-2012 |
| | Gravitino LSP | 0 | mono-jet | Yes | 20.3 | $F^{1/2}$ scale 865 GeV m(\tilde{G})>1.8 × 10 ⁻⁴ eV, m(\tilde{g})=m(\tilde{q})=1.5 TeV | 1502.01518 |
| ÷ | $\tilde{g} \rightarrow b \bar{b} \tilde{\chi}_{1}^{0}$ | 0 | 3 <i>b</i> | Yes | 20.1 | \tilde{g} 1.25 TeV $m(\tilde{\chi}_1^0) < 400 \text{ GeV}$ | 1407.0600 |
| 2 | $\tilde{g} \rightarrow t \bar{t} \tilde{\chi}_{1}^{0}$ | 0 | 7-10 jets | | 20.3 | \tilde{g} 1.1 To $m(\tilde{k}_1^0)$ <350 GeV | 1308.1841 |
| S IIICU. | $\tilde{g} \rightarrow t \bar{t} \tilde{\chi}_{1}^{0}$ | 0-1 <i>e</i> , μ | 3 <i>b</i> | Yes | 20.1 | \tilde{g} 1.34 TeV m(\tilde{k}_{1}^{0})<400 GeV | 1407.0600 |
| | $\tilde{g} \rightarrow b t \tilde{\chi}_1^+$ | 0-1 <i>e</i> , <i>µ</i> | 3 <i>b</i> | Yes | 20.1 | \tilde{g} 1.3 TeV m $(\tilde{\chi}_1^0)$ <300 GeV | 1407.0600 |
| | $\tilde{b}_1 \tilde{b}_1, \tilde{b}_1 \rightarrow b \tilde{\chi}_1^0$ | 0 | 2 <i>b</i> | Yes | 20.1 | \tilde{b}_1 100-620 GeV m($\tilde{\chi}_1^0$)<90 GeV | 1308.2631 |
| | $\tilde{b}_1 \tilde{b}_1, \tilde{b}_1 \rightarrow t \tilde{\chi}_1^{\pm}$ | 2 <i>e</i> , μ (SS) | 0-3 <i>b</i> | Yes | | \tilde{b}_1 275-440 GeV $m(\tilde{\chi}_1^{\pm})=2 m(\tilde{\chi}_1^0)$ | 1404.2500 |
| 5 | $\tilde{t}_1 \tilde{t}_1, \tilde{t}_1 \rightarrow b \tilde{\chi}_1^{\pm}$ | 1-2 <i>e</i> , μ | 1-2 <i>b</i> | Yes | | \tilde{t}_1 110-167 GeV $m(\tilde{\chi}_1^{\pm}) = 2m(\tilde{\chi}_1^0), m(\tilde{\chi}_1^0) = 55 \text{ GeV}$ | 1209.2102, 1407. |
| 5 | $\tilde{t}_1 \tilde{t}_1, \tilde{t}_1 \rightarrow Wb \tilde{\chi}_1^0 \text{ or } t \tilde{\chi}_1^0$ | 2 <i>e</i> , μ | 0-2 jets | Yes | | \tilde{t}_1 90-191 GeV 215-530 GeV m($\tilde{\chi}_1^0$)=1 GeV | 1403.4853, 1412. |
| $\tilde{t}_1 \tilde{t}_1$ | $\tilde{t}_1 \tilde{t}_1, \tilde{t}_1 \rightarrow t \tilde{\chi}_1^0$ | 0-1 <i>e</i> , μ | 1-2 <i>b</i> | Yes | | \tilde{t}_1 210-640 GeV m($\tilde{\chi}_1^0$)=1 GeV | 1407.0583,1406.1 |
| | $\tilde{t}_1 \tilde{t}_1, \tilde{t}_1 \rightarrow c \tilde{\chi}_1^0$ | | iono-jet/c-ta | - | | 7 <mark>1 90-240 Ge</mark> V m(ĩ ₁)-m(ξ ¹ ₁)<85 GeV | 1407.0608 |
| | $\tilde{t}_1 \tilde{t}_1$ (natural GMSB) | $2 e, \mu (Z)$ | 1 <i>b</i> | Yes | | \tilde{t}_1 150-580 GeV $m(\tilde{\chi}_1^0)$ >150 GeV | 1403.5222 |
| | $\tilde{t}_2 \tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z$ | 3 <i>e</i> , μ (Z) | 1 <i>b</i> | Yes | | <i>ĩ</i> ₂ 290-600 GeV m(<i>κ</i> ̃ ⁰)<200 GeV | 1403.5222 |
| | $\tilde{\ell}_{\mathrm{L,R}}\tilde{\ell}_{\mathrm{L,R}}, \tilde{\ell} \rightarrow \ell \tilde{\chi}_{1}^{0}$ | 2 <i>e</i> , µ | 0 | Yes | 20.3 | \tilde{t} 90-325 GeV m $(\tilde{\chi}_1^0)$ =0 GeV | 1403.5294 |
| | $\tilde{\chi}_1^+ \tilde{\chi}_1^-, \tilde{\chi}_1^+ \rightarrow \tilde{\ell} \nu(\ell \tilde{\nu})$ | 2 e, µ | 0 | Yes | | $\mathbf{\tilde{\chi}_{1}^{\pm}} \qquad \qquad \mathbf{m}(\tilde{\chi}_{1}^{0}) = 0 \text{ GeV}, \\ \mathbf{m}(\tilde{\ell}, \tilde{\nu}) = 0.5 (\mathbf{m}(\tilde{\chi}_{1}^{\pm}) + \mathbf{m}(\tilde{\chi}_{1}^{0})) $ | 1403.5294 |
| 3 | $\tilde{\chi}_1^+ \tilde{\chi}_1^-, \tilde{\chi}_1^+ \to \tilde{\tau} \nu(\tau \tilde{\nu})$ | 2 τ | - | Yes | 20.3 | $\tilde{\chi}_{1}^{\pm}$ 100-350 GeV $m(\tilde{\chi}_{1}^{0})=0$ GeV, $m(\tilde{\tau},\tilde{\nu})=0.5(m(\tilde{\chi}_{1}^{\pm})+m(\tilde{\chi}_{1}^{0}))$ | 1407.0350 |
| nirect | $\tilde{\chi}_{1}^{\pm}\tilde{\chi}_{2}^{0} \rightarrow \tilde{\ell}_{L} \nu \tilde{\ell}_{L} \ell(\tilde{\nu}\nu), \ell \tilde{\nu} \tilde{\ell}_{L} \ell(\tilde{\nu}\nu)$ | 3 <i>e</i> , µ | 0 | Yes | | $\texttt{m}(\tilde{\chi}_{1}^{\pm}) = \texttt{m}(\tilde{\chi}_{2}^{0}), \ \texttt{m}(\tilde{\chi}_{1}^{0}) = \texttt{0}, \ \texttt{m}(\tilde{\chi}_{1}^{0}) = \texttt{0}, \ \texttt{m}(\tilde{\chi}_{1}^{0}) + \texttt{m}(\tilde{\chi}_{1}^{0}))$ | 1402.7029 |
| | $\tilde{\chi}_1^{\pm} \tilde{\chi}_2^{\bar{0}} \rightarrow W \tilde{\chi}_1^0 Z \tilde{\chi}_1^0$ | 2-3 <i>e</i> , µ | 0-2 jets | Yes | 20.3 | $\tilde{\chi}_1^{\pm}, \tilde{\chi}_2^0$ 420 GeV $m(\tilde{\chi}_1^{\pm}) = m(\tilde{\chi}_2^0), m(\tilde{\chi}_1^0) = 0$, sleptons decoupled | 1403.5294, 1402. |
| | $\tilde{\chi}_{1}^{\pm}\tilde{\chi}_{2}^{\bar{0}} \rightarrow W \tilde{\chi}_{1}^{\dot{0}} h \tilde{\chi}_{1}^{\dot{0}}, h \rightarrow b \bar{b} / W W / \tau \tau / \gamma$ | | 0-2 <i>b</i> | Yes | 20.3 | $\tilde{\chi}_1^{\pm}, \tilde{\chi}_2^0$ 250 GeV $m(\tilde{\chi}_1^{\pm}) = m(\tilde{\chi}_2^0), m(\tilde{\chi}_1^0) = 0$, sleptons decoupled | 1501.07110 |
| _ | $\tilde{\chi}_{2}^{0}\tilde{\chi}_{3}^{0}, \tilde{\chi}_{2,3}^{0} \to \tilde{\ell}_{\mathrm{R}}\ell$ | 4 <i>e</i> ,μ | 0 | Yes | 20.3 | $\tilde{\chi}_{2,3}^{0} \qquad \qquad$ | 1405.5086 |
| | Direct $\tilde{\chi}_1^+ \tilde{\chi}_1^-$ prod., long-lived $\tilde{\chi}_1^\pm$ | Disapp. trk | 1 jet | Yes | 20.3 | $\tilde{\chi}_1^{\pm}$ 270 GeV $m(\tilde{\chi}_1^{\pm}) - m(\tilde{\chi}_1^{0}) = 160$ MeV, $\tau(\tilde{\chi}_1^{\pm}) = 0.2$ ns | 1310.3675 |
| 2 | Stable, stopped \tilde{g} R-hadron | 0 | 1-5 jets | Yes | 27.9 | \tilde{g} 832 GeV $m(\tilde{\chi}_1^0)=100$ GeV, 10 μ s< $\tau(\tilde{g})<1000$ s | 1310.6584 |
| 2 | Stable \tilde{g} R-hadron | trk | - | - | 19.1 | <u>8</u> 1. <mark>0</mark> 7 TeV | 1411.6795 |
| par incres | GMSB, stable $\tilde{\tau}, \tilde{\chi}_1^0 \rightarrow \tilde{\tau}(\tilde{e}, \tilde{\mu}) + \tau(e, \tilde{\mu})$ | | - | | 19.1 | $\tilde{\chi}_1^0$ 537 GeV 10 <tan<math>\beta<50</tan<math> | 1411.6795 |
| 2 | GMSB, $\tilde{\chi}_1^0 \rightarrow \gamma \tilde{G}$, long-lived $\tilde{\chi}_1^0$ | 2 γ 1 μ, displ. vtx | _ | Yes | 20.3 | \tilde{x}_{1}^{0} 435 GeV $2 < \tau(\tilde{x}_{1}^{0}) < 3$ ns, SPS8 model | 1409.5542 |
| _ | 44, 1 44/ () | | · . | | 20.3 | \tilde{q} 1.0 TeV 1.5 < $c\tau < 156$ mm, BR(μ)=1, m(\tilde{t}_1^0)=108 GeV | |
| | LFV $pp \rightarrow \tilde{v}_{\tau} + X, \tilde{v}_{\tau} \rightarrow e + \mu$ | 2 e, µ | - | - | 4.6 | $\tilde{\nu}_r$ 1.61 TeV $\lambda'_{311}=0.10, \lambda_{132}=0.05$ | 1212.1272 |
| | $LFV \ pp \rightarrow \tilde{v}_{\tau} + X, \tilde{v}_{\tau} \rightarrow e(\mu) + \tau$ | $1 e, \mu + \tau$ | - | - | 4.6 | $\tilde{\nu}_r$ 1.1 TeV $\lambda'_{311}=0.10, \lambda_{1(2)33}=0.05$ | 1212.1272 |
| X X X | Bilinear RPV CMSSM | 2 <i>e</i> , μ (SS) | 0-3 <i>b</i> | Yes | 20.3 | \tilde{q}, \tilde{g} .35 TeV $m(\tilde{q})=m(\tilde{g}), c\tau_{LSP} < 1 \text{ mm}$ | 1404.2500 |
| | $\tilde{\chi}_1^+ \tilde{\chi}_1^-, \tilde{\chi}_1^+ \to W \tilde{\chi}_1^0, \tilde{\chi}_1^0 \to e e \tilde{\nu}_{\mu}, e \mu \tilde{\nu}_e$ | $4 e, \mu$ | - | Yes | 20.3 | $\tilde{\chi}_1^{\pm}$ 750 GeV $m(\tilde{\chi}_1^0) > 0.2 \times m(\tilde{\chi}_1^{\pm}), \lambda_{121} \neq 0$ | 1405.5086 |
| | $\tilde{\chi}_1^+\tilde{\chi}_1^-, \tilde{\chi}_1^+ \to W \tilde{\chi}_1^0, \tilde{\chi}_1^0 \to \tau \tau \tilde{\nu}_e, e \tau \tilde{\nu}_\tau$ | $3 e, \mu + \tau$ | - | Yes | 20.3 | $\tilde{\chi}_{1}^{\pm}$ 450 GeV m($\tilde{\chi}_{1}^{0}$)>0.2×m($\tilde{\chi}_{1}^{\pm}$), $\lambda_{133} \neq 0$ | 1405.5086 |
| | $\tilde{g} \rightarrow qqq$ $\tilde{g} \rightarrow \tilde{t}_1 t, \tilde{t}_1 \rightarrow bs$ | 0 2 <i>e</i> , μ (SS) | 6-7 jets 0-3 <i>b</i> | - Yes | 20.3 20.3 | ĝ 916 GeV BR(t)=BR(c)=0% ĝ 850 GeV | ATLAS-CONF-201 1404.250 |
| | Scalar charm, $\tilde{c} \rightarrow c \tilde{\chi}_1^0$ | 0 | 2 c | Yes | 20.3 | č 490 GeV m(\tilde{k}_1^0)<200 GeV | 1501.01325 |
| er | Scalar charm, $c \rightarrow c \chi_1$ | 0 | ∠ C | ies | 20.3 | | 1301.01325 |

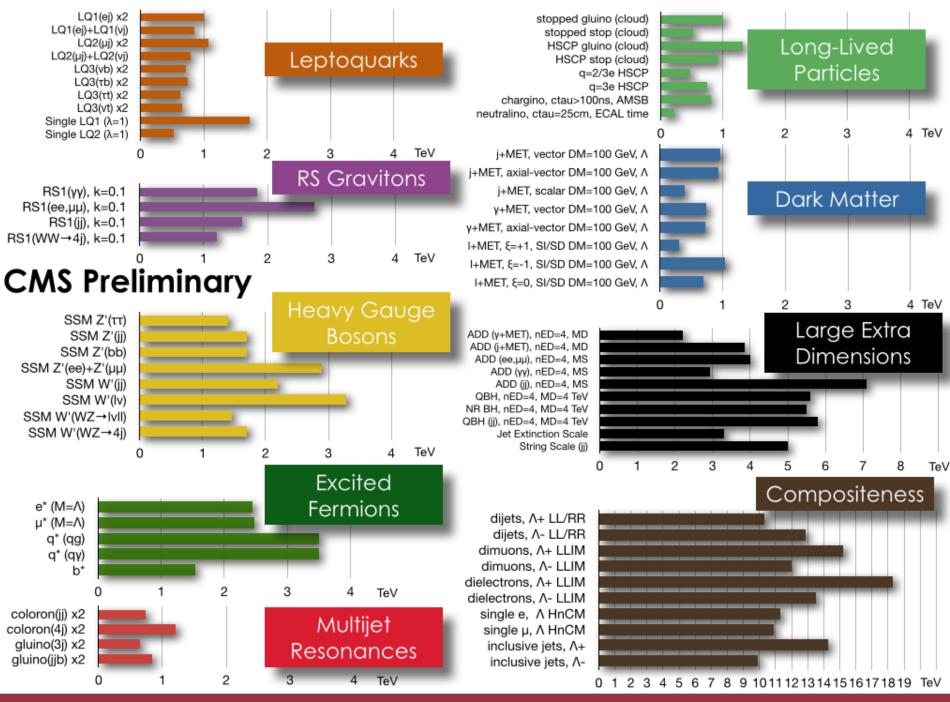
ATLAS Preliminary

*Only a selection of the available mass limits on new states or phenomena is shown. All limits quoted are observed minus 1 σ theoretical signal cross section uncertainty.

Summary of CMS SUSY Results* in SMS framework

SUSY 2013





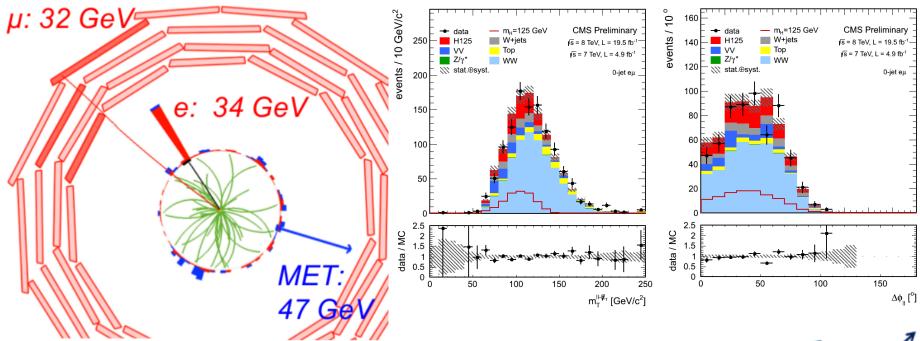
CMS Exotica Physics Group Summary – Moriond, 2015

Conclusions

- The SM is one of the best theories that mankind has ever come up with so far.
- The LHC has widely opened the door to the Terascale. This is where we expect the new physics to set in!
- It has brought the discovery of a Higgs boson thus the completion the SM!
 Spontaneous symmetry breaking is not a trick nor a back door solution, it is reality!
- It is clear that we have to account for physics beyond the SM. Where is it? Good arguments that LHC run-II has a good chance to bring it.
- If so the LHC will bring us the best times ever for particle physics!

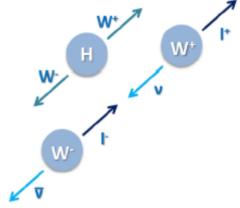






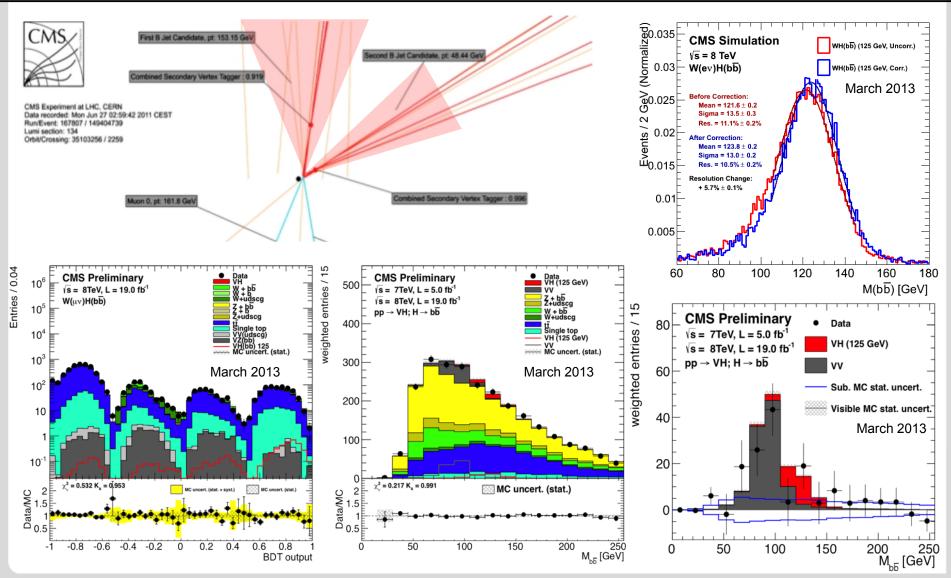
• High discovery potential, but bad mass resolution.

| ff | 0-jet | 1-jet | 2-jet(VBF) |
|------------|-------|-------|------------|
| $\int ff'$ | 0-jet | 1-jet | 2-jet(VBF) |



$H \rightarrow bb$ Decay Channel

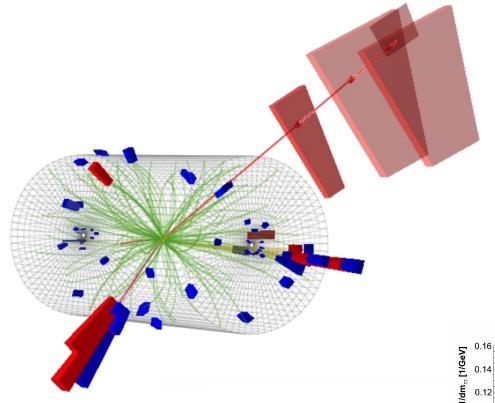




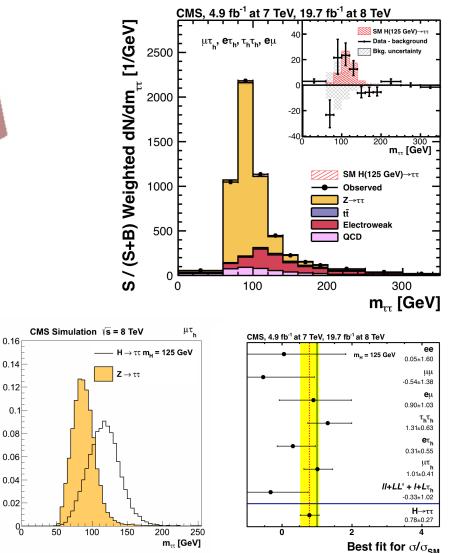
Institute of Experimental Particle Physics (IEKP)

$H \rightarrow \tau \tau$ Decay Channel





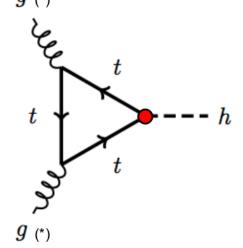
- $m_{\tau\tau}$ as main discriminating variable.
- Separation between irreducible $Z \rightarrow \tau \tau$ background and $H \rightarrow \tau \tau$ signal.

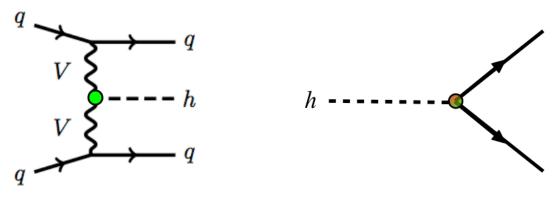




• Determine couplings from production mode and decay channel:

 $gg \rightarrow H$ production: $qq \rightarrow qqH$ production: Decay to f or V: g (*)



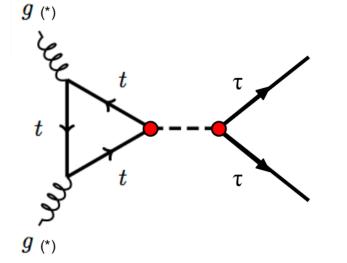


- f : $\kappa_{\text{Hff}} = \frac{m_f}{v}$ • V : $\kappa_{\text{HVV}} = \frac{2m_V^2}{v}$
- Coupling to gluon can be f or effective (*).
- Coupling to γ can be effective or a mixture of f + V.
- Direct measurement not possible since κ_i appear in nominator and denominator of $BR_i = \frac{\kappa_i}{\Gamma_h} = \frac{\kappa_i}{\sum \kappa}$

Narrow Width Approximation



• Propagator: $\frac{1}{(q^2-m^2+m^2\Gamma^2)} \rightarrow \frac{\pi}{m\Gamma}\delta(q^2-m^2)$ for $\Gamma \rightarrow 0$.



- i.e. put propagating particle on shell.
- Calculate cross section as $\sigma \times BR$.

• BR_X =
$$\frac{\Gamma_X}{\Gamma_H}$$
, $\Gamma_H = \sum_i \Gamma_i$.

- $\sigma \propto (\kappa_t \kappa_\tau)^2 \propto (\kappa_u \kappa_d)^2 \propto (\kappa_q \kappa_f)^2 \propto (\kappa_g \kappa_f)^2$.
- For each production mode and decay channel collect κ_i and express Γ_H as sum of individual κ_i .