

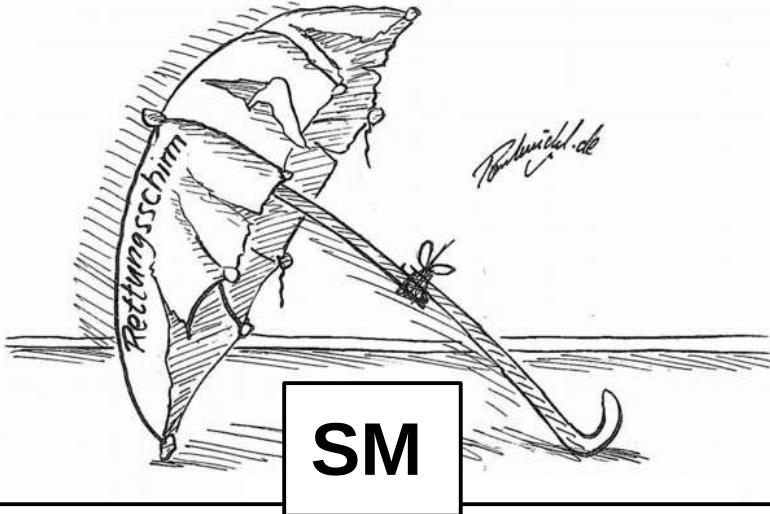
Moderne Experimentalphysik III: Hadronen und Teilchen (Physik VI)

Thomas Müller, Roger Wolf

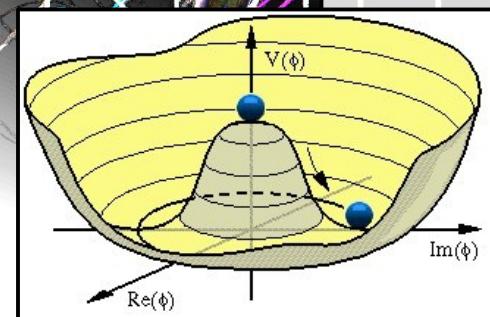
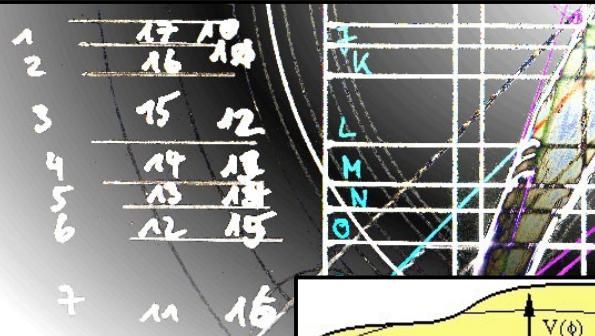
28. Juni 2018

INSTITUTE OF EXPERIMENTAL PARTICLE PHYSICS (IETP) – PHYSICS FACULTY





Problem: lokale Eichsymmetrien in Lagrangedichte sind durch massive Teilchen explizit gebrochen



Wie kann eine Symmetrie zur gleichen Zeit erhalten und gebrochen sein?

Spontane Symmetriebrechung:

$$f(x, y) = x^2 + y^2$$

$$x = r \cos \varphi$$

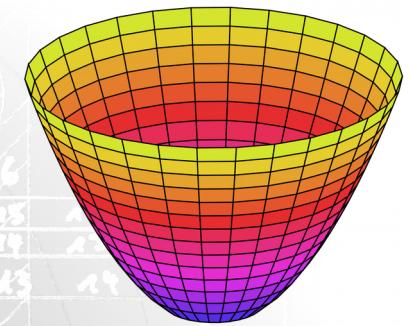
$$y = r \sin \varphi$$

$$f(x, y)|_{r, \varphi} = r^2 (\cos \varphi^2 + \sin \varphi^2) = r^2$$

$$\tilde{f}(x, y) = (x - 1)^2 + (y - 1)^2$$

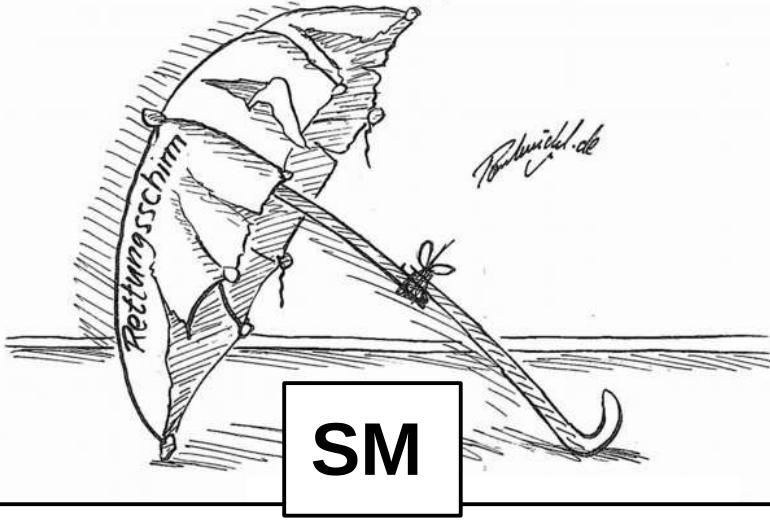
$$\tilde{f}(x, y)|_{r, \varphi} = r^2 + 2(1 - r(\sin \varphi + \cos \varphi))$$

("hidden symmetry")



Führe Potential ein das den Grundzustand des Universums aus der Symmetriearchse der Bewegungsgleichungen zwingt.

→ Teilchenmasse als Kopplung an nicht verschwindenden Vakuumerwartungswert.



Problem: lokale Eichsymmetrien in Lagrangedichte sind durch massive Teilchen explizit gebrochen

Erinnerung:

$$f_{H \rightarrow ff} = i \frac{m_f}{v} \quad (\text{Fermions})$$

$$f_{H \rightarrow VV} = i \frac{2m_V^2}{v} \quad (\text{Heavy Bosons trilinear})$$

$$f_{HH \rightarrow VV} = i \frac{2m_V^2}{v^2} \quad (\text{Heavy Bosons quartic})$$

$$f_{H \rightarrow HH} = i \frac{3m_H^2}{v} \quad (H \text{ Boson trilinear})$$

$$f_{HH \rightarrow HH} = i \frac{3m_H^2}{v^2} \quad (H \text{ Boson quartic})$$

Wie kann eine Symmetrie zur gleichen Zeit erhalten und gebrochen sein?

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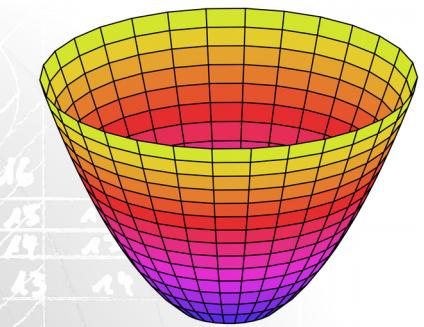
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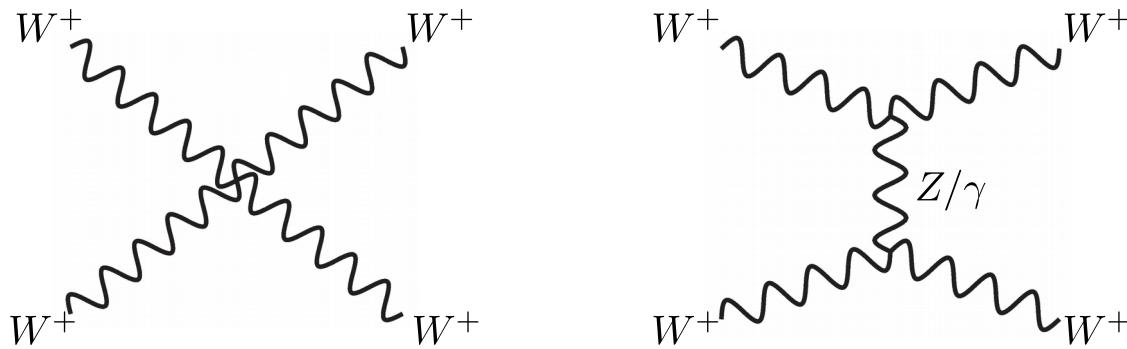


Führe Potential ein das den Grundzustand des Universums aus der Symmetriearchse der Bewegungsgleichungen zwingt.

→ Teilchenmasse als Kopplung an nicht verschwindenden Vakuumerwartungswert.

Higgs sector in the light of (tree-level) unitarity

- Unitarity problem demonstrated for $W^+W^+ \rightarrow W^+W^+$ scattering:

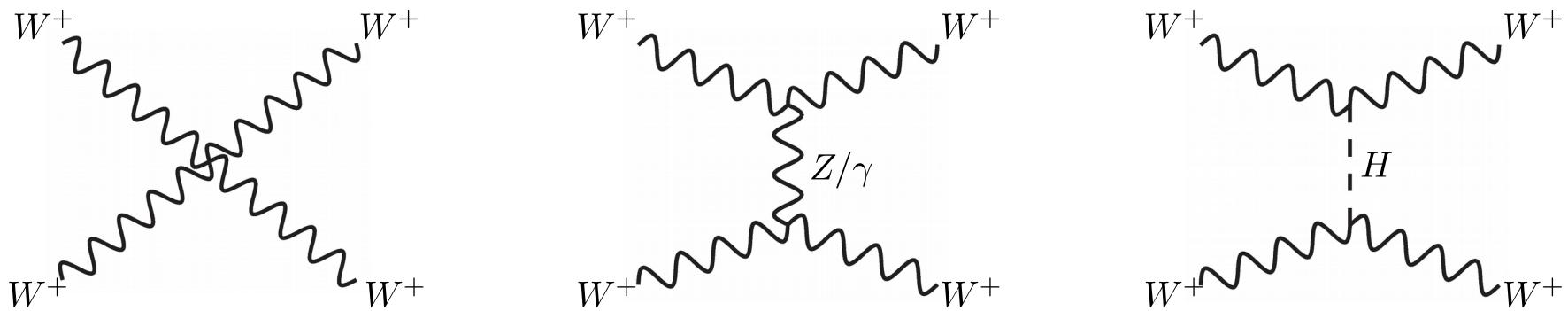


$$\mathcal{M}_{gauge} = -g^2 \frac{s}{4m_W^2} + \mathcal{O}(s^0)$$

constraint

Higgs sector in the light of (tree-level) unitarity

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$$\mathcal{M}_{gauge} = -g^2 \frac{s}{4m_W^2} + \mathcal{O}(s^0)$$

constraint

$$\mathcal{M}_H = g_{HWW}^2 \frac{s}{m_W^4} + \mathcal{O}(s^0)$$

Exact cancellation of **divergent behavior**
only if scalar exchange particle has coupling
of type $\propto m_W^2$.

$$g_{HWW} = \frac{2m_W^2}{v} = g \cdot m_W$$

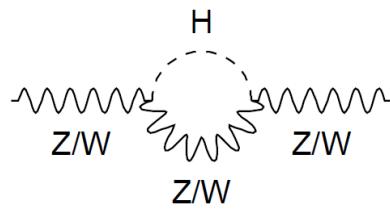
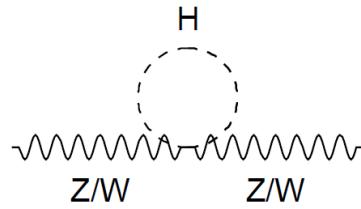
$$\text{with: } v = \frac{2m_W}{g}$$

- Any additional contribution to this process should preserve this cancellation.

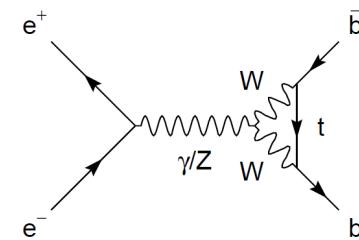
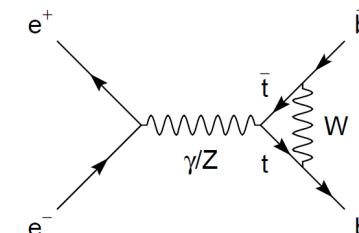
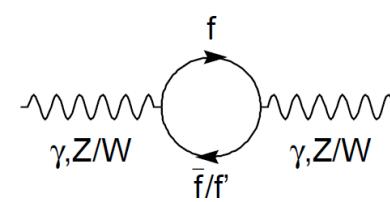
Higher orders and precision observables

- Particles, which cannot be directly observed at lower energy scales, still have influence on observables, due to higher order corrections in loops.

The Higgs/*top* in propagator loops:



The *top* in vertex loops:



- Introduce direct dependencies of (measurable) effective vector boson masses and couplings on m_H & m_t .

Higher order corrections to m_W

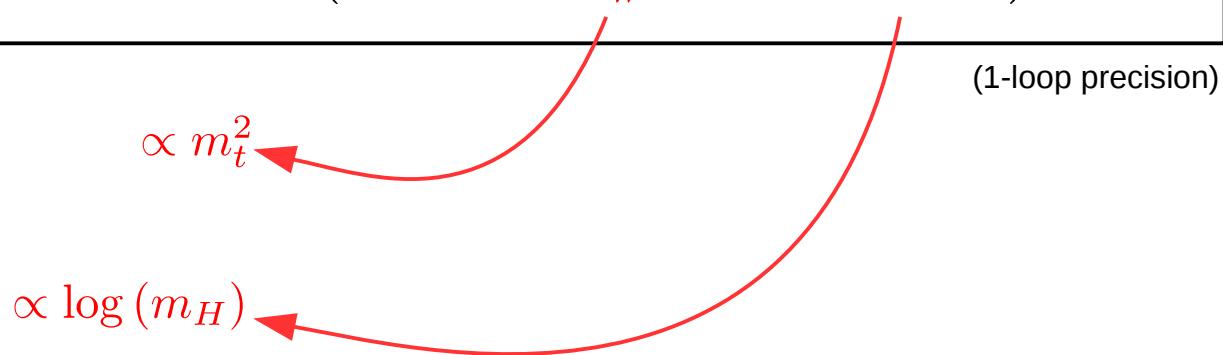
- Higher order corrections to m_W :

$$m_W^2 = \frac{m_Z^2}{2} \left(1 + \sqrt{1 - 4 \frac{\alpha\pi}{\sqrt{2}G_F m_Z^2} \cdot \frac{1}{1-\Delta r}} \right) \quad \Delta r = \Delta\alpha + \Delta r_W$$

$$\Delta\alpha = \Delta\alpha_{\text{lep}} + \Delta\alpha_{\text{top}} + \Delta\alpha_{\text{had}}^{(5)}$$

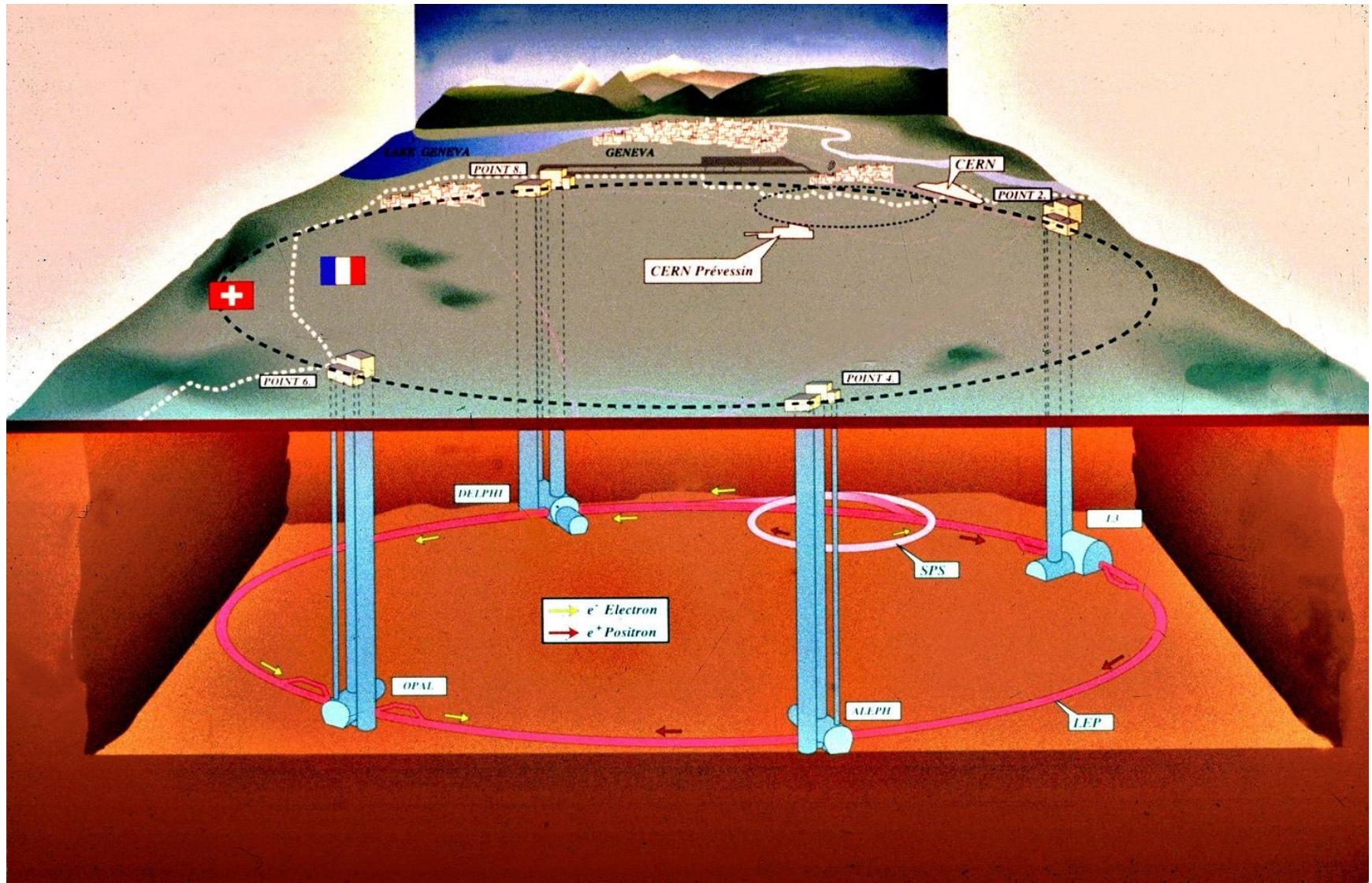
$$\Delta r_W(m_t, m_H) \simeq \frac{\alpha}{\pi \sin^2 \theta_W} \left(-\frac{3}{16} \frac{\cos^2 \theta_W}{\sin^2 \theta_W} \frac{m_t^2}{m_W^2} + \frac{11}{24} \log \left(\frac{m_H}{m_Z} \right) \right)$$

(1-loop precision)



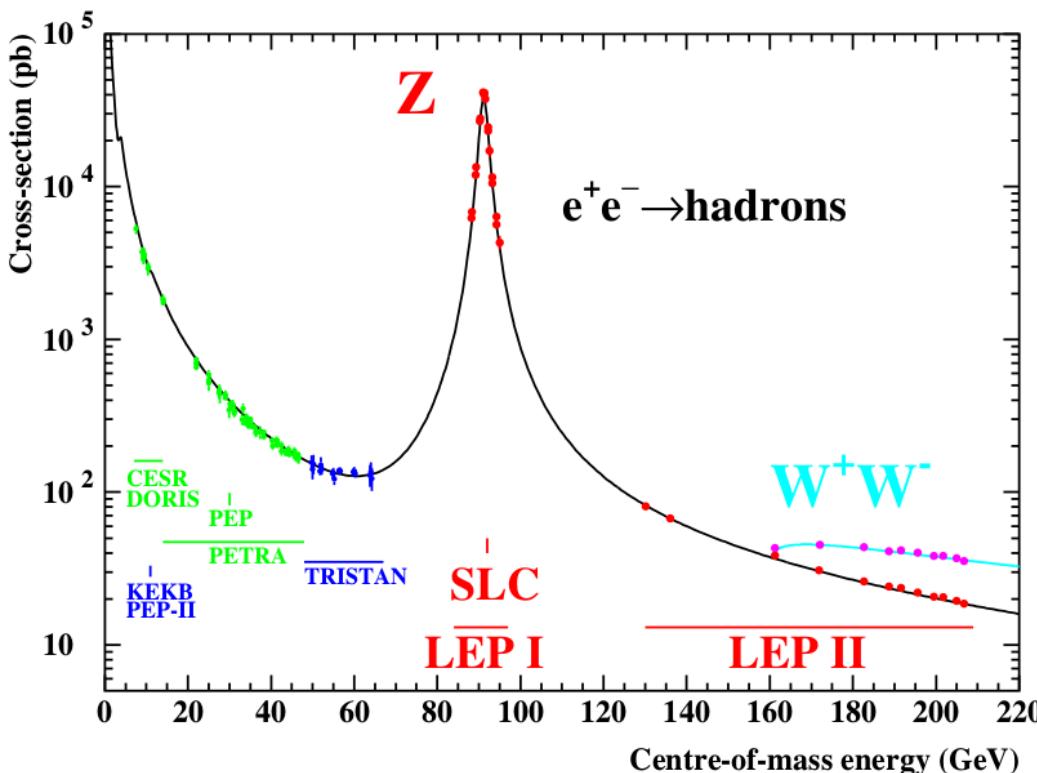
- Effects set in at $\mathcal{O}(\alpha^2) \approx \mathcal{O}(10^{-4}) \rightarrow$ high precision needed on observables and theoretical prediction!

High precision measurements @ LEP & SLAC



High precision observables @ LEP

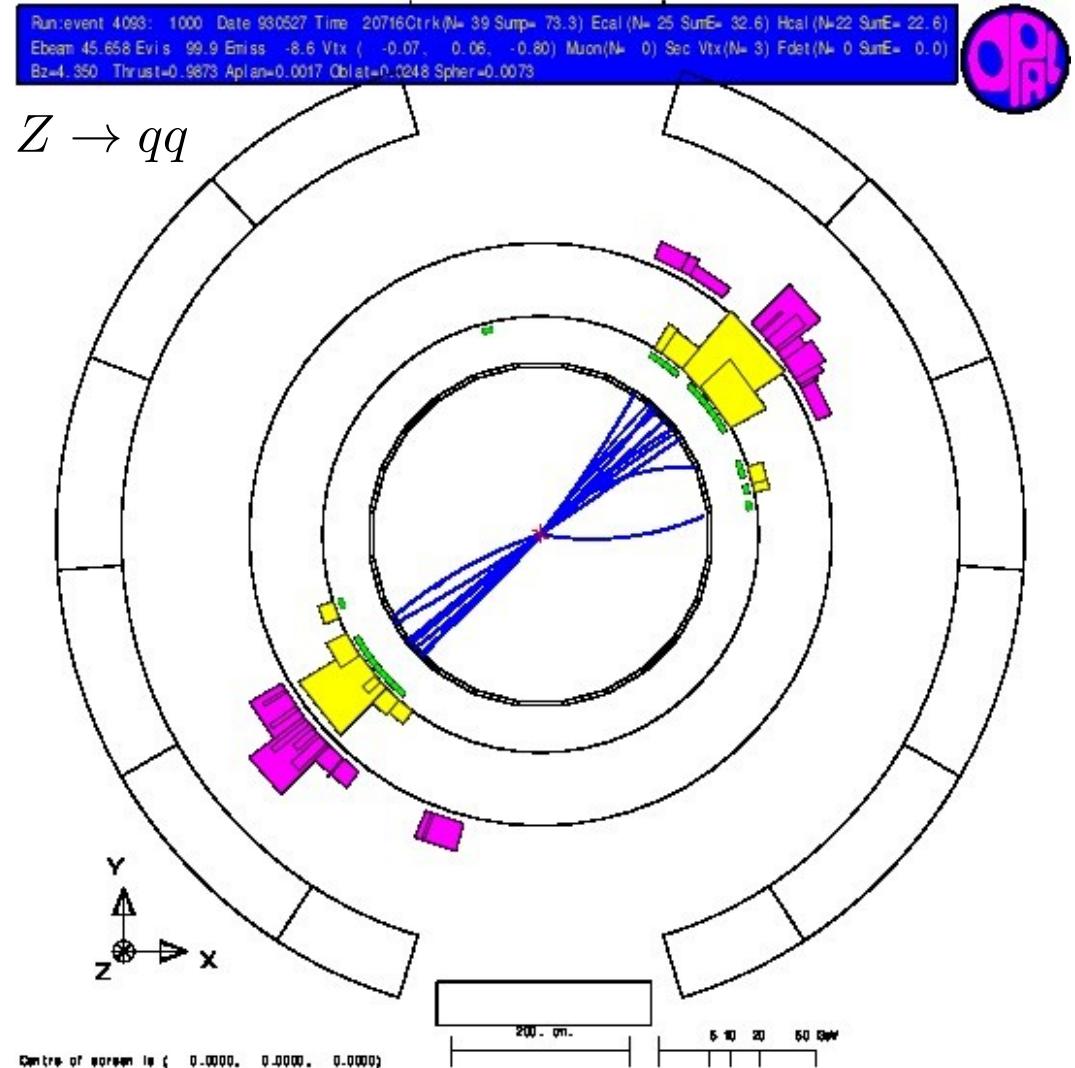
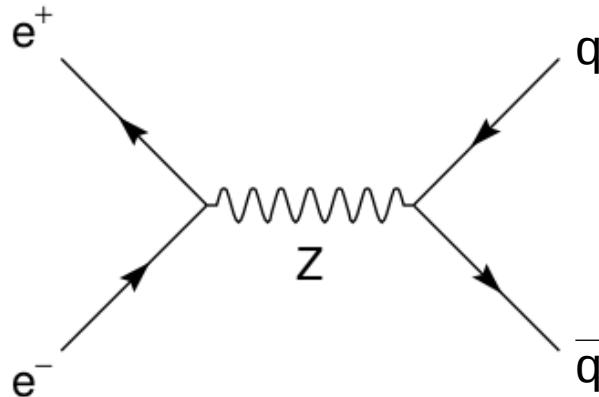
- High precision measurements made at $\sqrt{s} = m_Z$ during LEP-I run period:



Year	Centre-of-mass energy range [GeV]	Integrated luminosity [pb ⁻¹]
1989	88.2 – 94.2	1.7
1990	88.2 – 94.2	8.6
1991	88.5 – 93.7	18.9
1992	91.3	28.6
1993	89.4, 91.2, 93.0	40.0
1994	91.2	64.5
1995	89.4, 91.3, 93.0	39.8
	202.1	

- $15 \cdot 10^6 Z \rightarrow qq$ events
- $1.7 \cdot 10^6 Z \rightarrow \ell\ell$ events

Typical $Z \rightarrow qq$ event @ LEP



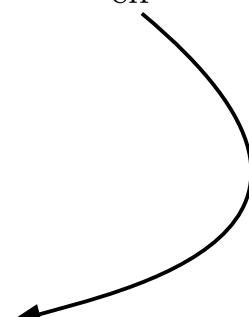
Z-pole electroweak precision observables

Pseudo-Observable	Measured Value		
$\Delta\alpha_{\text{had}}^{(5)}(m_Z)$	0.02758	\pm	0.00034
m_Z [GeV]	91.1875	\pm	0.0021
Γ_Z [GeV]	2.4952	\pm	0.0023
σ_{had}^0 [nb]	41.540	\pm	0.037
R_l^0	20.767	\pm	0.025
R_b^0	0.21629	\pm	0.00066
R_c^0	0.1721	\pm	0.0030
$A_{FB}^{0,l}$	0.0171	\pm	0.0010
$A_{FB}^{0,b}$	0.0992	\pm	0.0016
$A_{FB}^{0,c}$	0.0707	\pm	0.0035
$\sin^2 \theta_{\text{eff}}^{\text{lep}}$	0.2324	\pm	0.0012
$\mathcal{A}_l(\mathcal{P}_\tau)$	0.1465	\pm	0.0033
\mathcal{A}_b	0.923	\pm	0.020
\mathcal{A}_c	0.670	\pm	0.027
$\mathcal{A}_l(\text{SLD})$	0.1513	\pm	0.0021

(as of [hep-ex/0509008](#))

- 14(+1) observables.
- Precision between $\mathcal{O}(10^{-5})$ for m_Z & $\mathcal{O}(10^{-2})$ for $\mathcal{A}_l(\text{SLD})$ (including theoretical uncertainties).
- Exploit dependencies $\propto m_t^2$ and $\propto \log(m_H)$ of higher orders via relations in m_W and $\sin \theta_{\text{eff}}$.

NB: Using similar relations with the same dependencies as shown for m_W .



Parameter estimate

- Five parameter χ^2 fit:

Parameter	Best Fit Value	
$\Delta\alpha_{\text{had}}^{(5)}(m_Z)$	0.02759 ± 0.00035	$\Delta\alpha_{\text{had}}^{(5)}(m_Z)$
$\alpha_s(m_Z)$	0.1190 ± 0.0027	$\alpha_s(m_Z)$
m_Z	91.1874 ± 0.0021	m_Z
m_t	173 ± 11.5	m_t
$\log(m_H/\text{GeV})$	2.05 ± 0.29	$\log(m_H/\text{GeV})$

Fit of Z-pole observables only: ⁽¹⁾
 $\chi^2/ndof = 16/10$
 $\mathcal{P}(\chi^2) = 9.9\%$
(2005)

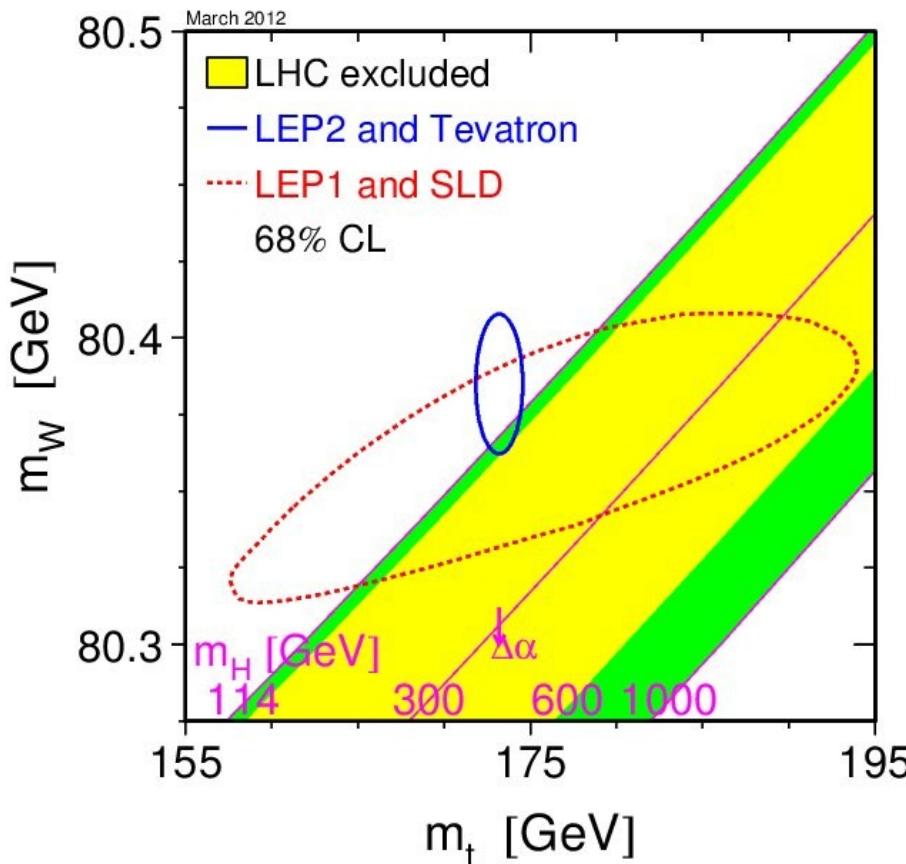
Fit of Z-pole observables +
 m_W, Γ_W, m_t : ⁽²⁾
 $\chi^2/ndof = 16.9/13$
 $\mathcal{P}(\chi^2) = 20.2\%$
(2012)



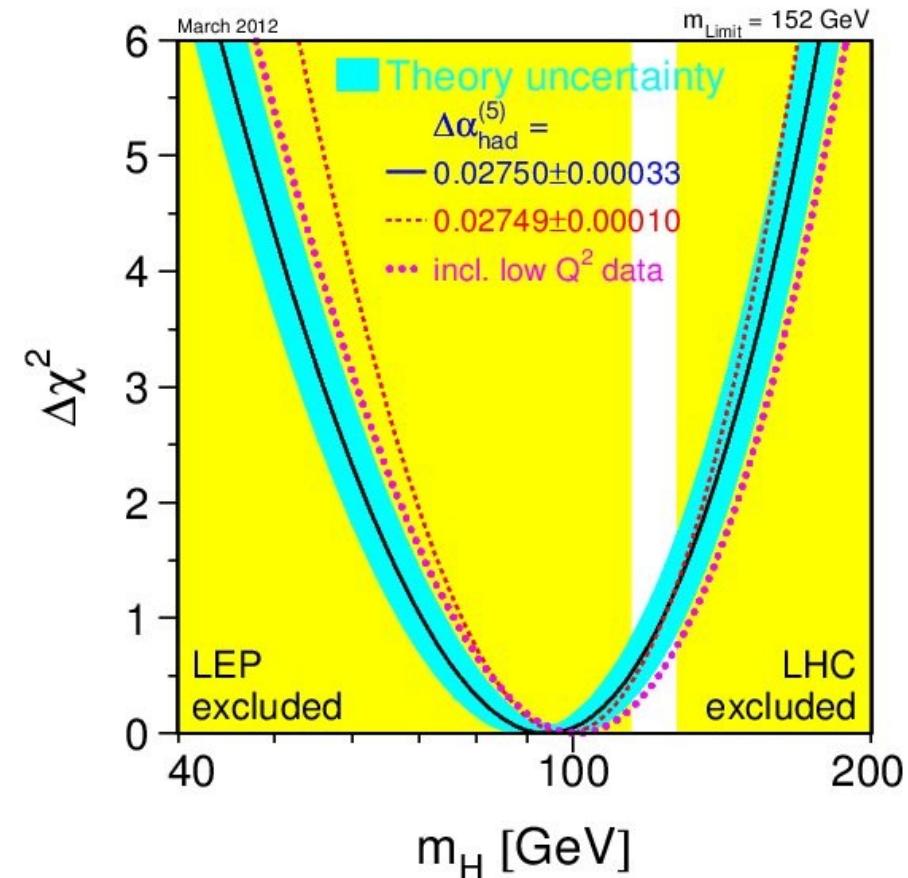
⁽¹⁾ (as of hep-ex/0509008)

⁽²⁾ http://lepewwwg.web.cern.ch/LEPEWWWG/winter12_results

Main result



Z-pole + $m_W + \Gamma_W$:
 $m_t = 178.1 \pm^{10.9}_{7.8}$ GeV



Z-pole + $m_W + \Gamma_W + m_t$:
 $m_H = 98 \pm^{25}_{21}$ GeV

Direkte Higgs Suchen und Entdeckung



Higgs Boson...

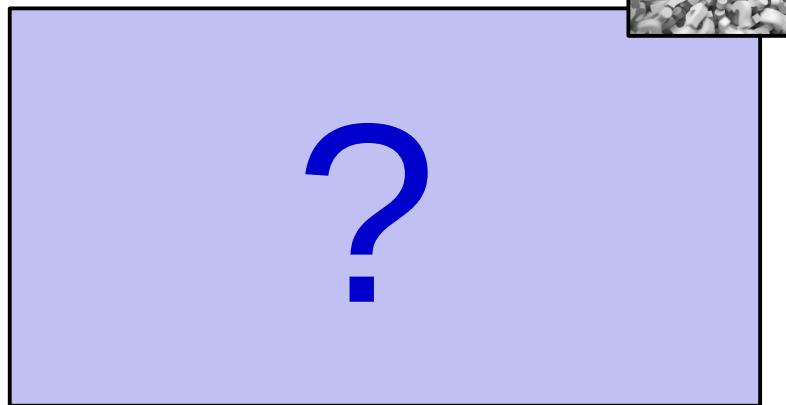
Google-Suche

Auf gut Glück!

[Google.de angeboten auf: English](#)

Direkte Suche bei LEP-II

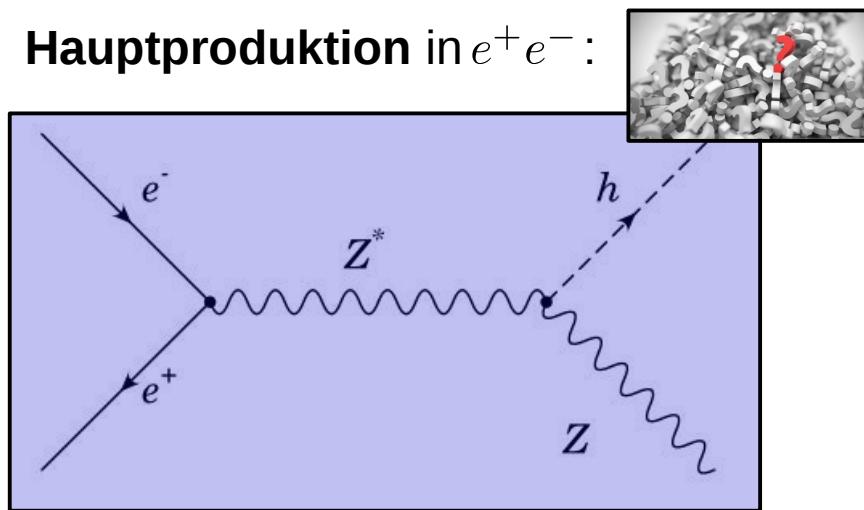
- **Hauptproduktion in e^+e^- :**



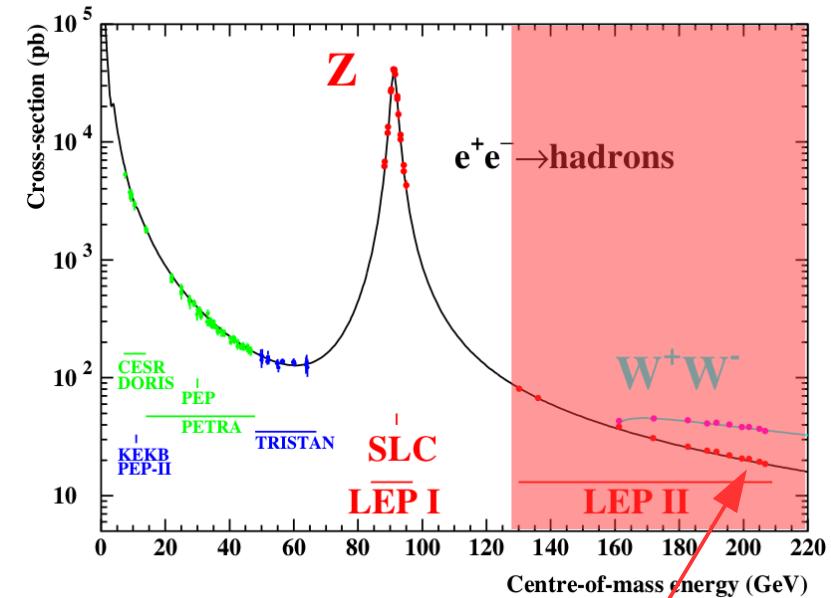
- Higgs Boson koppelt an Masse.
- Kopplung am stärksten für schwerste Objekte.

Direkte Suche bei LEP-II

- Hauptproduktion in e^+e^- :



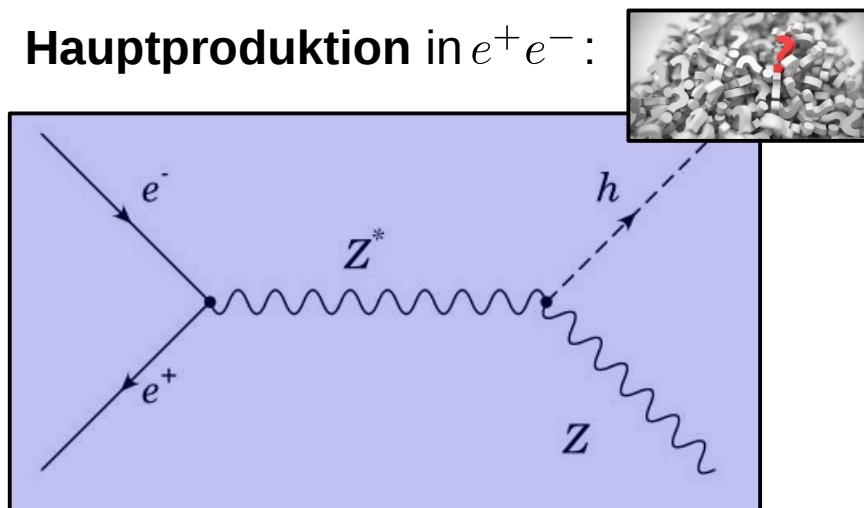
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Year	1996	1997	1998	1999	2000					
E_{CM} nominal [GeV]	161	172	183	189	192	196	200	202	205	207

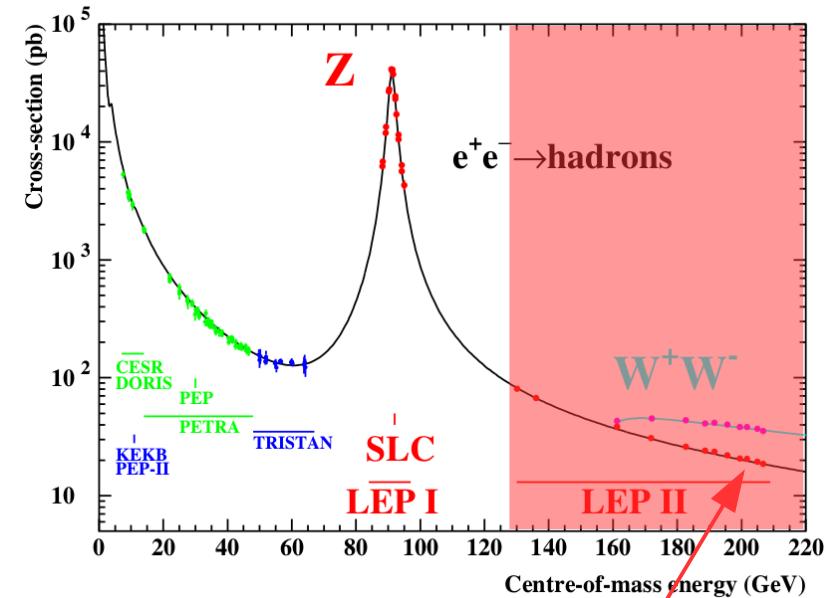
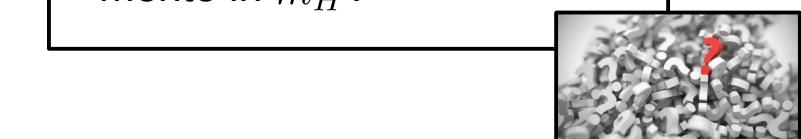
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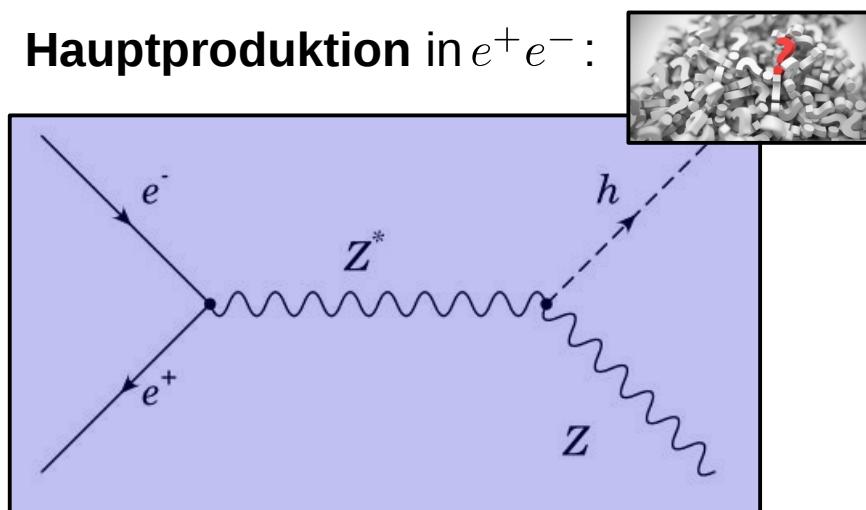
Wie groß war die maximale Reichweite der LEP-Experimente in m_H ?



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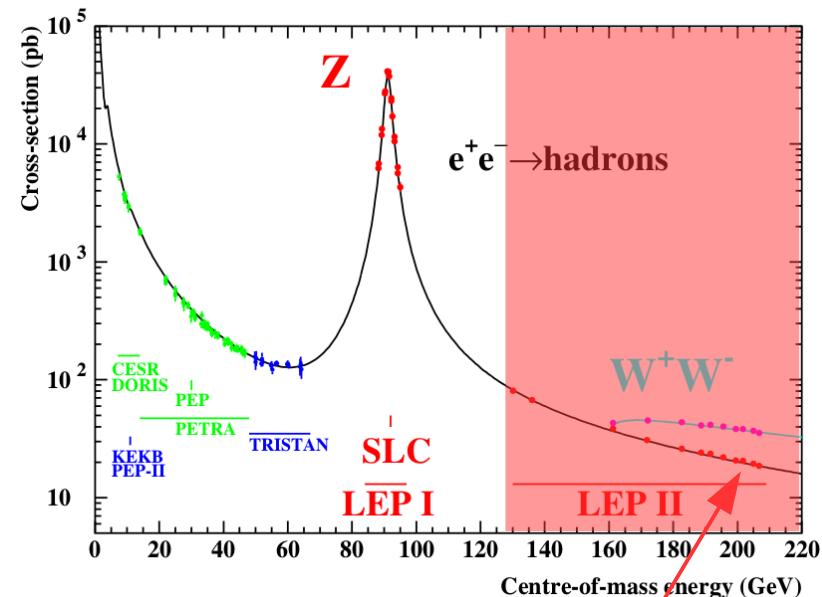


Zu vergleichen mit Folien 9 und 10

Wie groß war die maximale Reichweite der LEP-Experimente in m_H ? $\approx 117 \text{ GeV}$

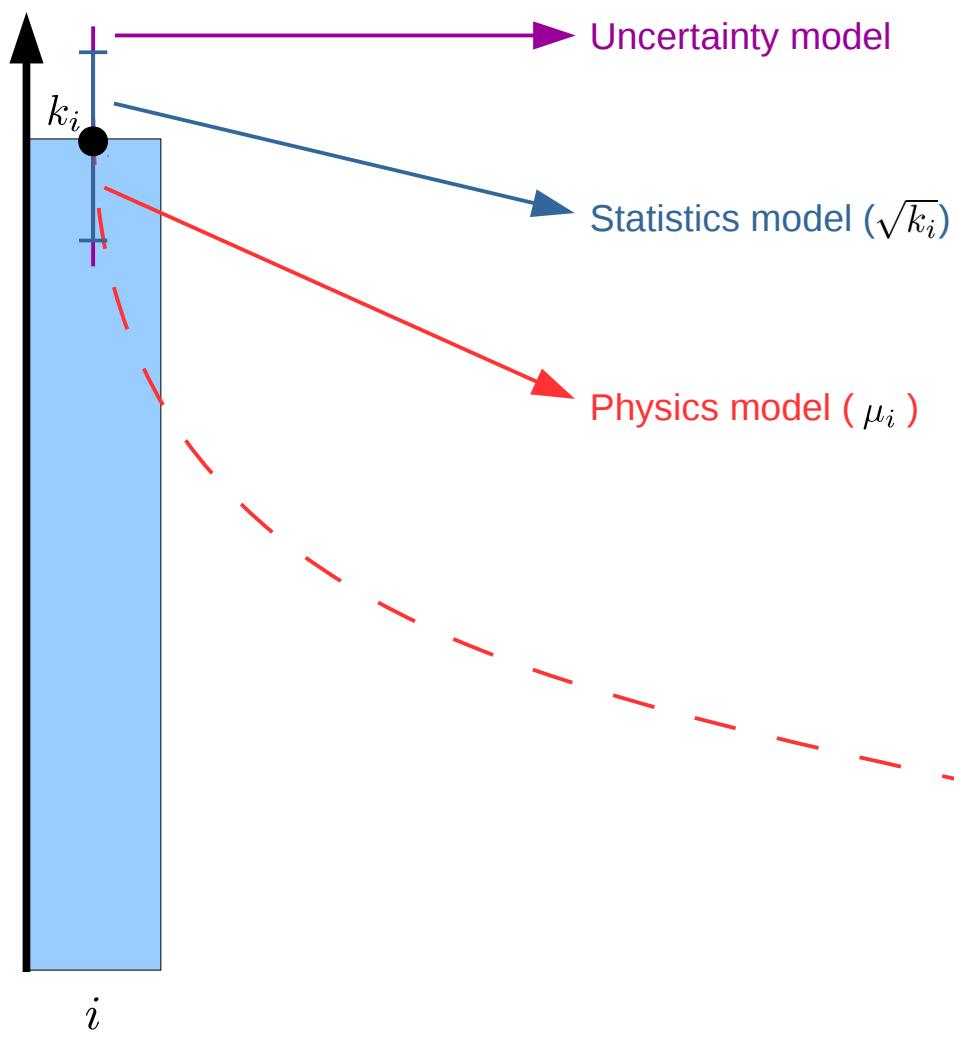


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Models in counting experiments

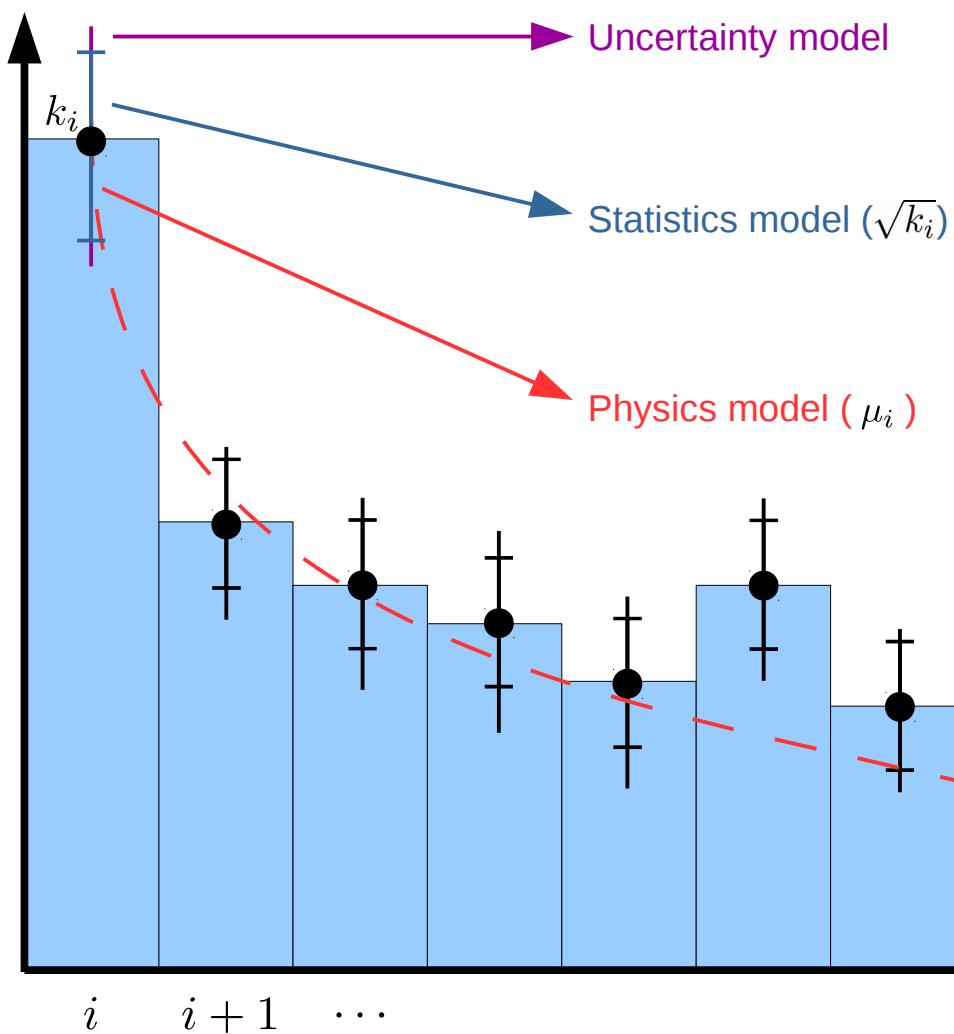


$$\mathcal{P}(k_i, \mu_i) = \frac{\mu_i^{k_i}}{k_i!} e^{-\mu_i}$$

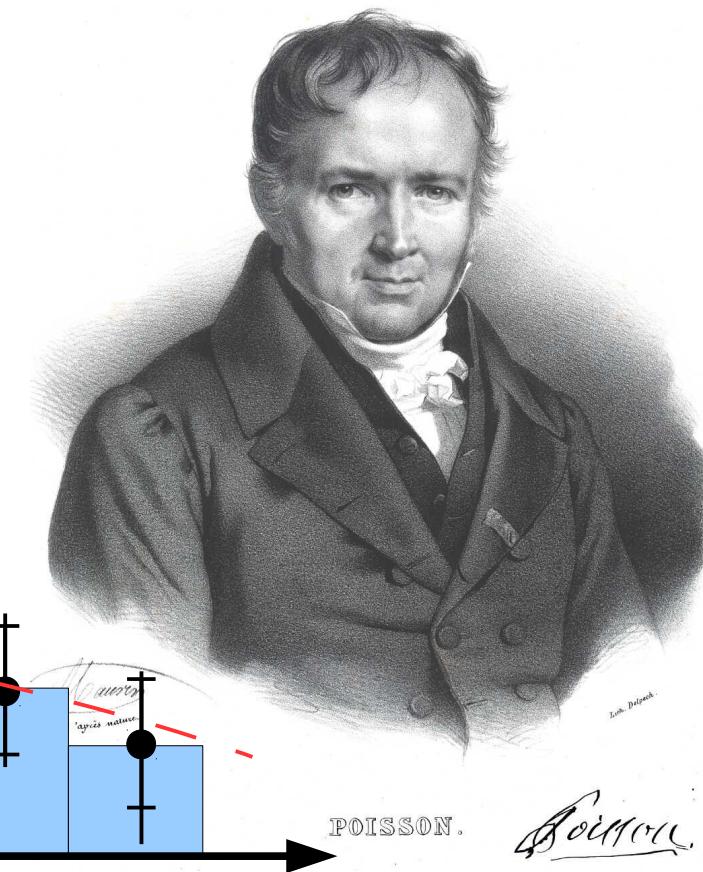


Siméon Denis Poisson
(21.07.1781 – 25.04.1840)

From one to many...



$$\prod_i \mathcal{P}(k_i, \mu_i) = \prod_i \frac{\mu_i^{k_i}}{k_i!} e^{-\mu_i}$$



Siméon Denis Poisson
(21.07.1781 – 25.04.1840)

Model building (likelihood functions)

- Likeliness of a model to be true quantified by *likelihood function* $\mathcal{L}(\{k_i\}, \{\kappa_j\})$

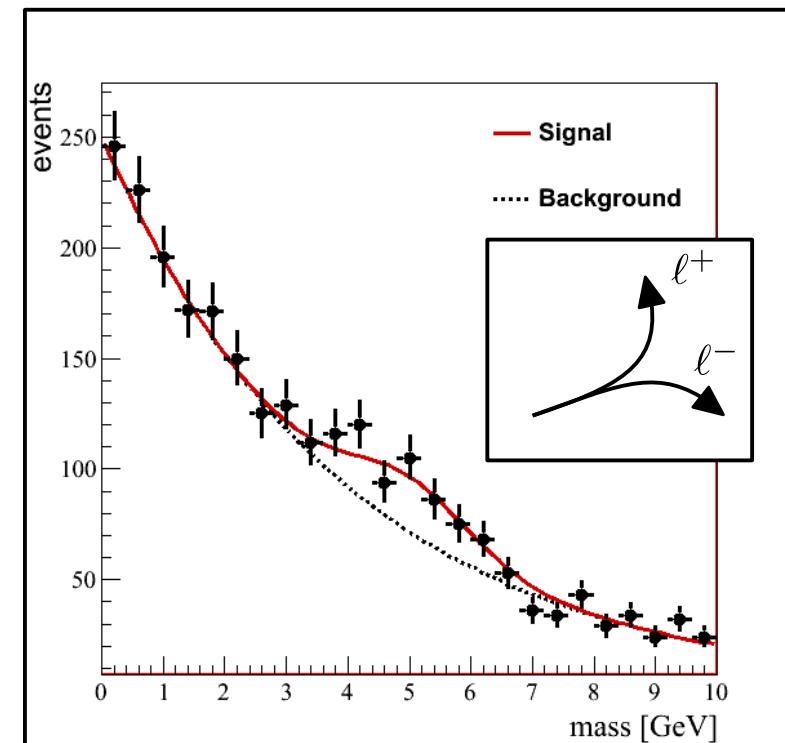
$$\prod_i \mathcal{P}(k_i, \mu_i) = \prod_i \frac{\mu_i^{k_i}}{k_i!} e^{-\mu_i}$$

model parameters.
measured number of events (e.g. in bins i).

- Simple example:
signal on top of known background in a binned histogram:

$$\mathcal{L}(\{k_i\}, \{\kappa_j\}) = \prod_i \underbrace{\mathcal{P}(k_i, \mu_i(\kappa_j))}_{\text{Product of pdfs for each bin (Poisson).}}$$

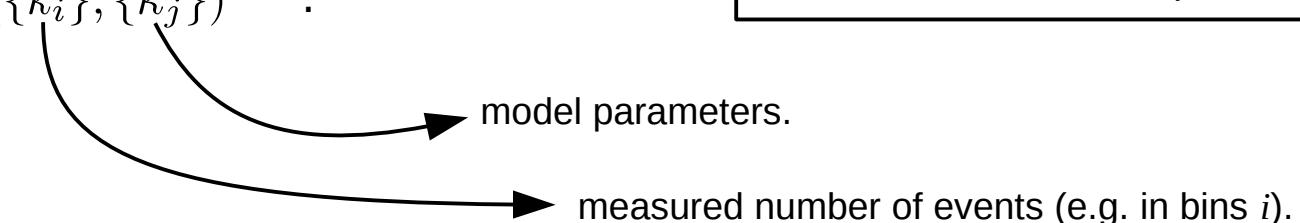
$$\mu_i(\kappa_j) = \underbrace{\kappa_0 \cdot e^{-\kappa_1 x_i}}_{\text{background}} + \underbrace{\kappa_2 \cdot e^{-(\kappa_3 - x_i)^2}}_{\text{signal}}$$



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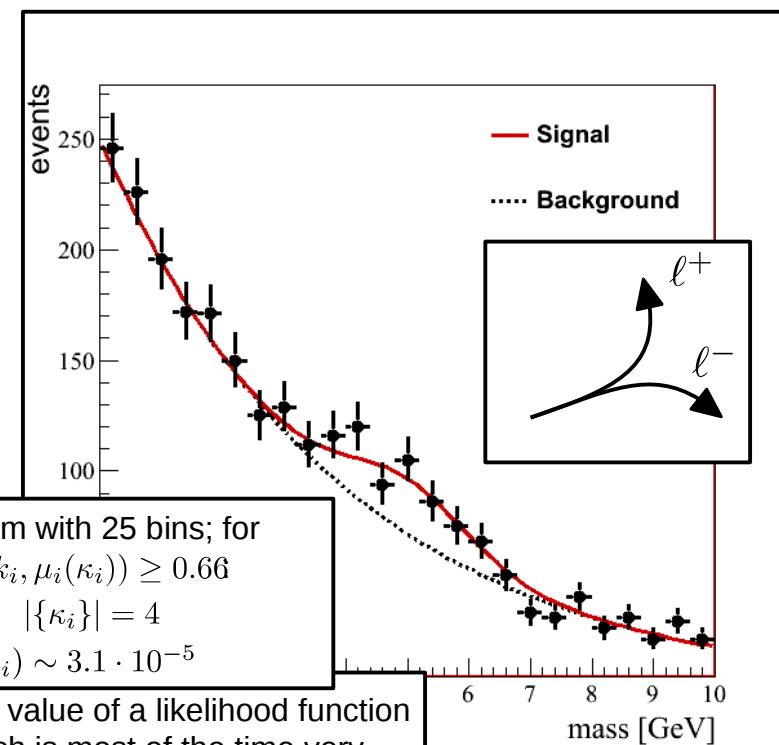
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Product of pdfs for each bin (Poisson).

$$\mu_i(\kappa_j) = \underbrace{\kappa_0 \cdot e^{-\kappa_1 x_i}}_{\text{background}} + \underbrace{\kappa_2 \cdot e^{-(\kappa_3 - x_i)^2}}_{\text{signal}}$$

EX: histogram with 25 bins; for each bin $\mathcal{P}(k_i, \mu_i(\kappa_i)) \geq 0.66$
 $|\{k_i\}| = 25 \quad |\{\kappa_i\}| = 4$
 $\prod \mathcal{P}(k_i, \mu_i(\kappa_i)) \sim 3.1 \cdot 10^{-5}$

NB: a value of a likelihood function as such is most of the time very close to zero, and w/o a reference in general w/o further meaning.



Example: test statistics (LEP ~2000)

- Test signal (H_1 , for fixed mass, m , and fixed signal strength, μ) vs. background-only (H_0).

$\mathcal{L}(n|b(\kappa_j)) = \prod_i \mathcal{P}(n_i|b_i(\kappa_j)) \times \prod_j \mathcal{C}(\kappa_j|\tilde{\kappa}_j)$

$\mathcal{L}(n|\mu s(\kappa_j) + b(\kappa_j)) = \prod_i \mathcal{P}(n_i|\mu s_i(\kappa_j) + b_i(\kappa_j)) \times \prod_j \mathcal{C}(\kappa_j|\tilde{\kappa}_j)$

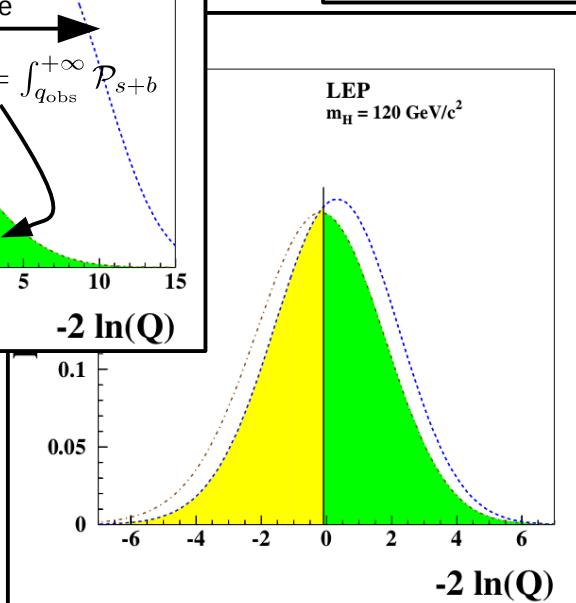
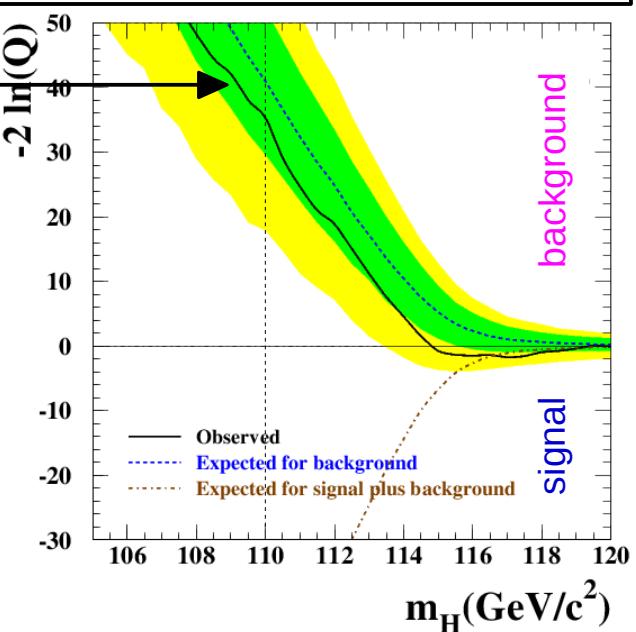
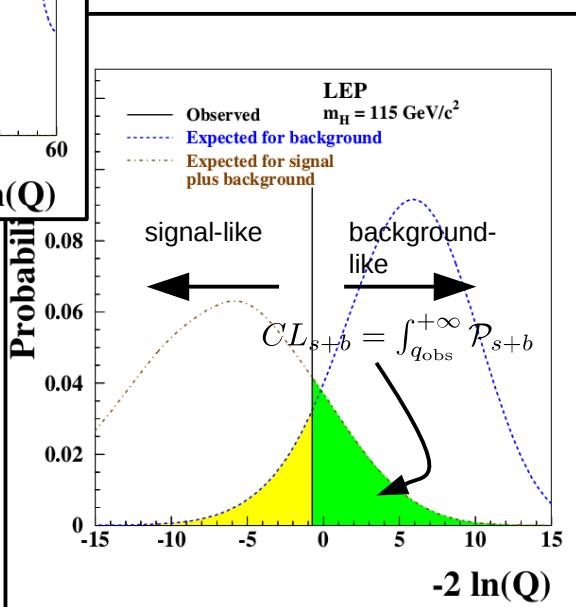
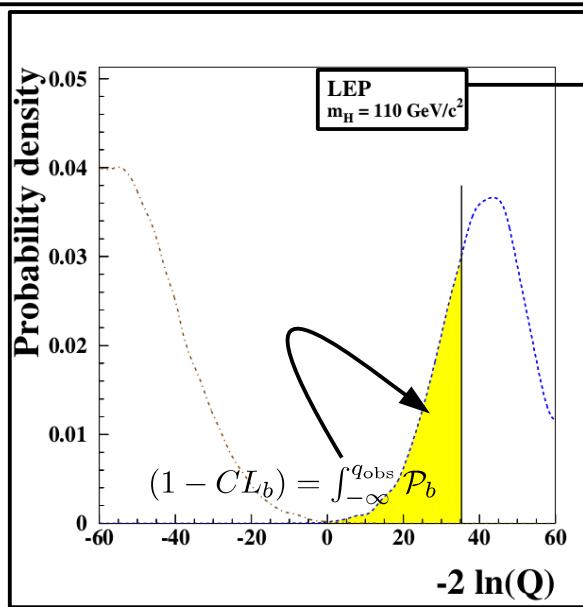
$q_\mu = -2 \ln \left(\frac{\mathcal{L}(n|\mu s+b)}{\mathcal{L}(n|b)} \right), \quad 0 \leq \mu$

nuisance parameters $\tilde{\kappa}_j$ integrated out before evaluation of q_μ (\rightarrow marginalization).

pdf's for nuisance parameters modified according to Bayes theorem.

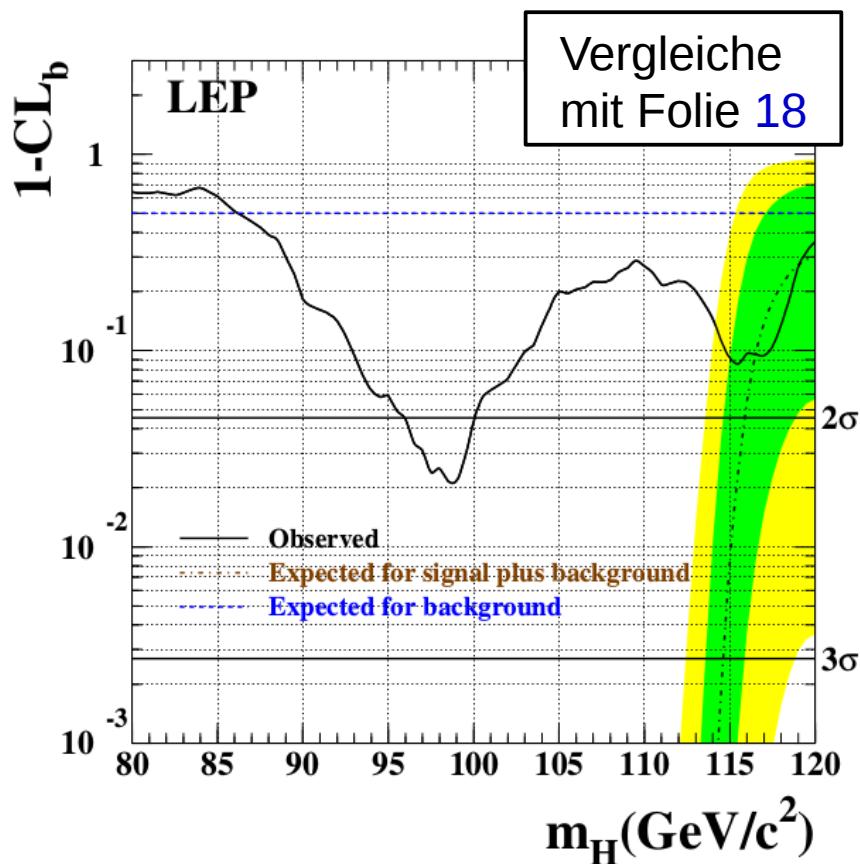
Example: test statistics (LEP)

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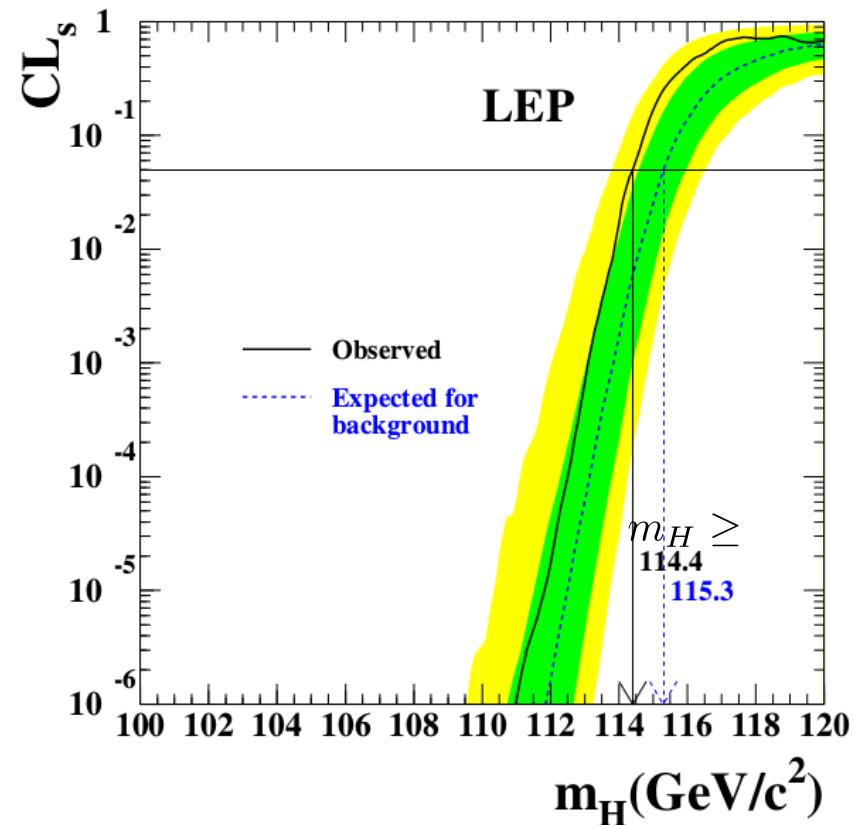


Ergebnis LEP-II

- p-Wert (Untergrundkompatibilität) :

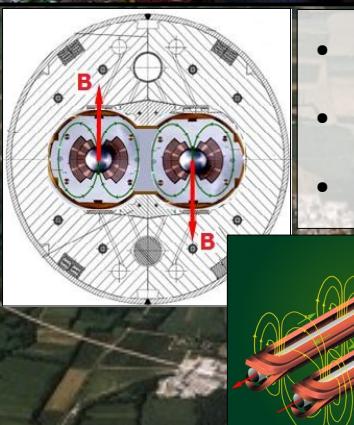
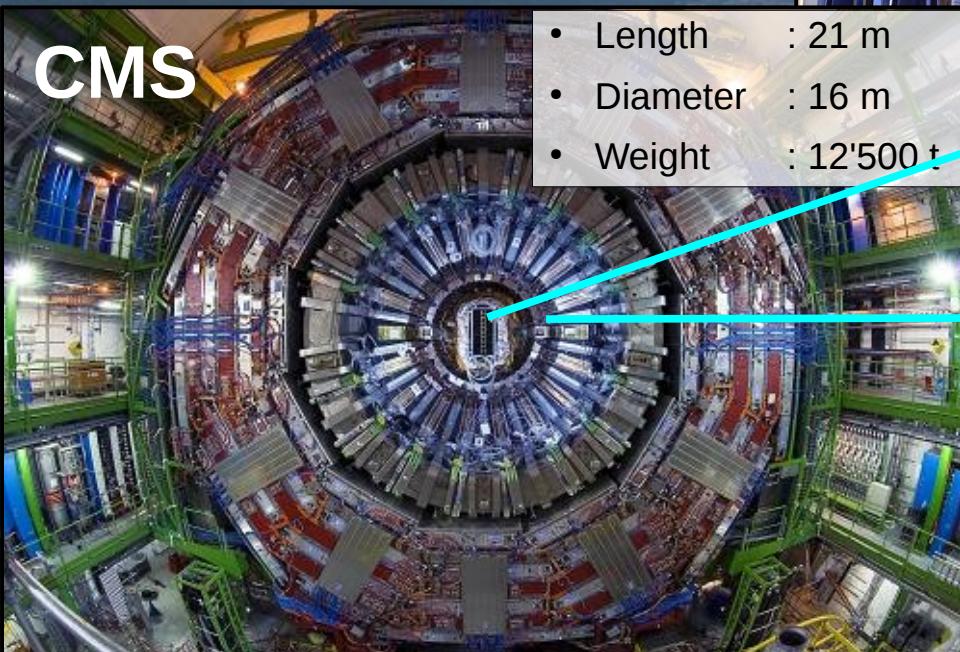


- Limit auf Masse (@ 95% CL):



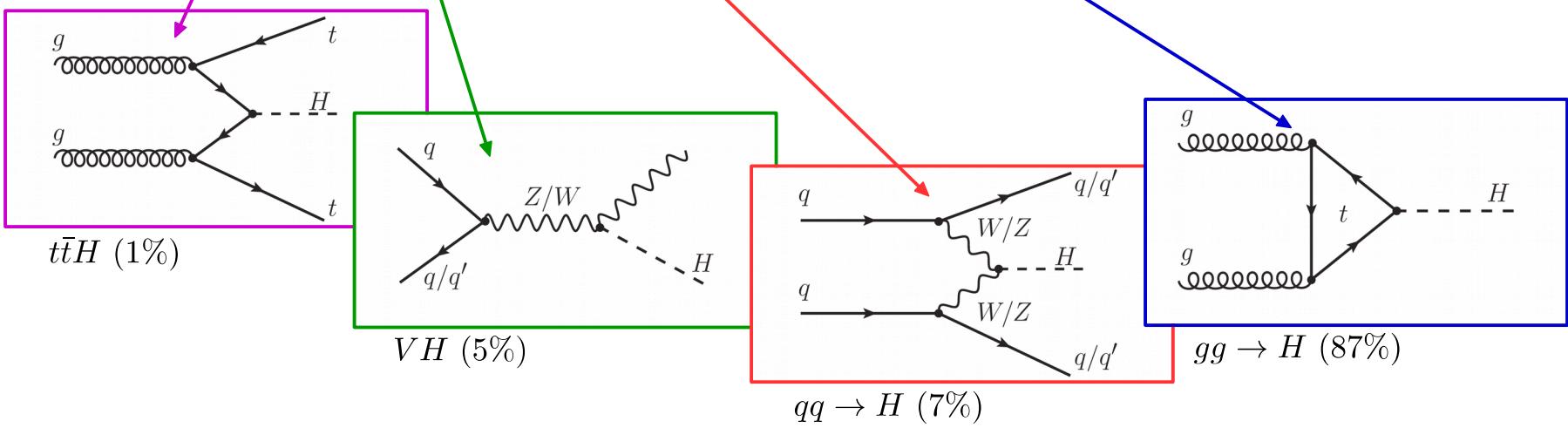
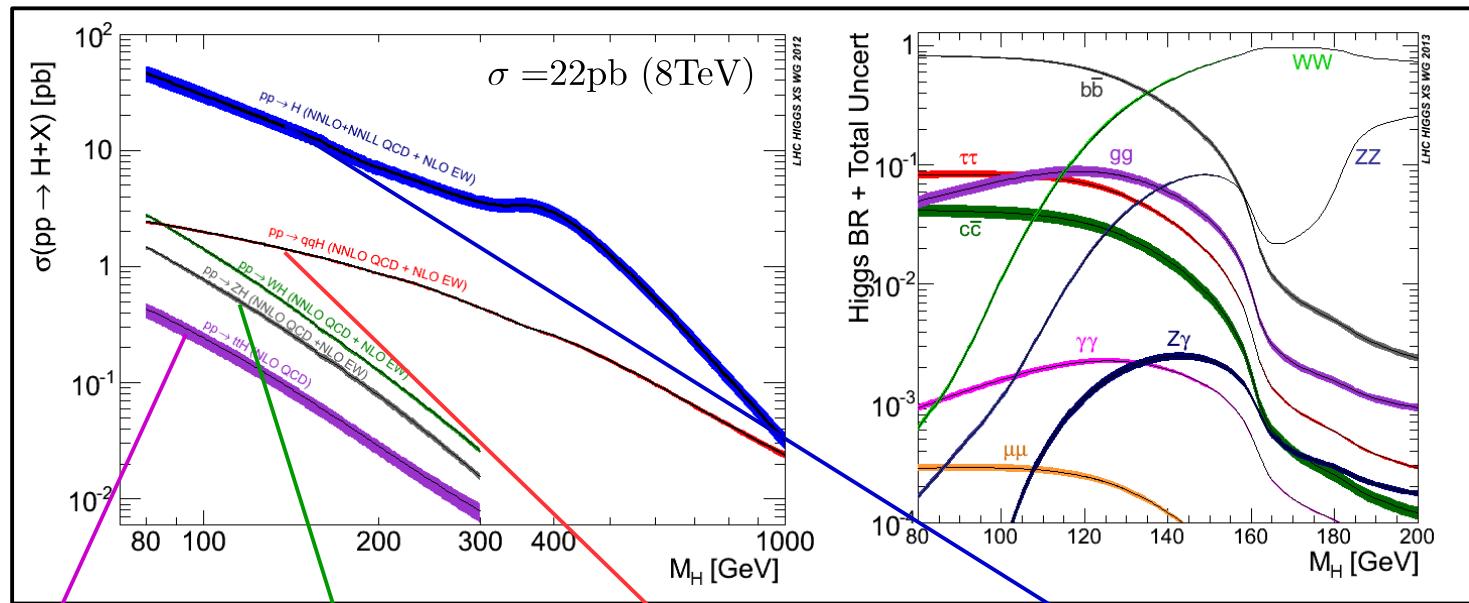
- Kein Signal beobachtet
- Untere Schranke auf Masse: $m_H > 114.4 \text{ GeV}$ (@ 95% CL)

Direkte Suche am LHC

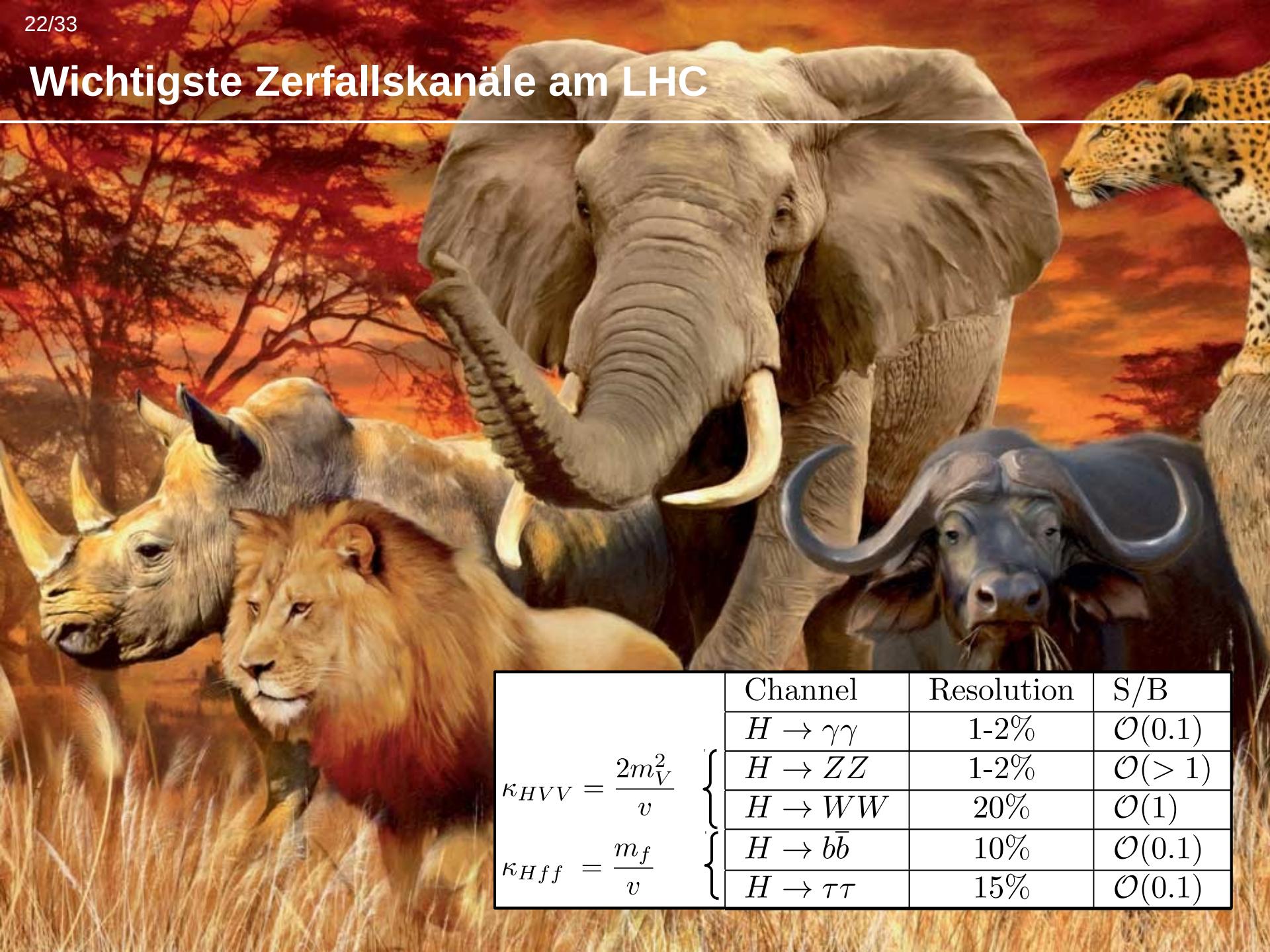


- Energy density 500 kJ/m.
- Tension 200'000 t/m.

Produktion und Zerfall



Wichtigste Zerfallskanäle am LHC

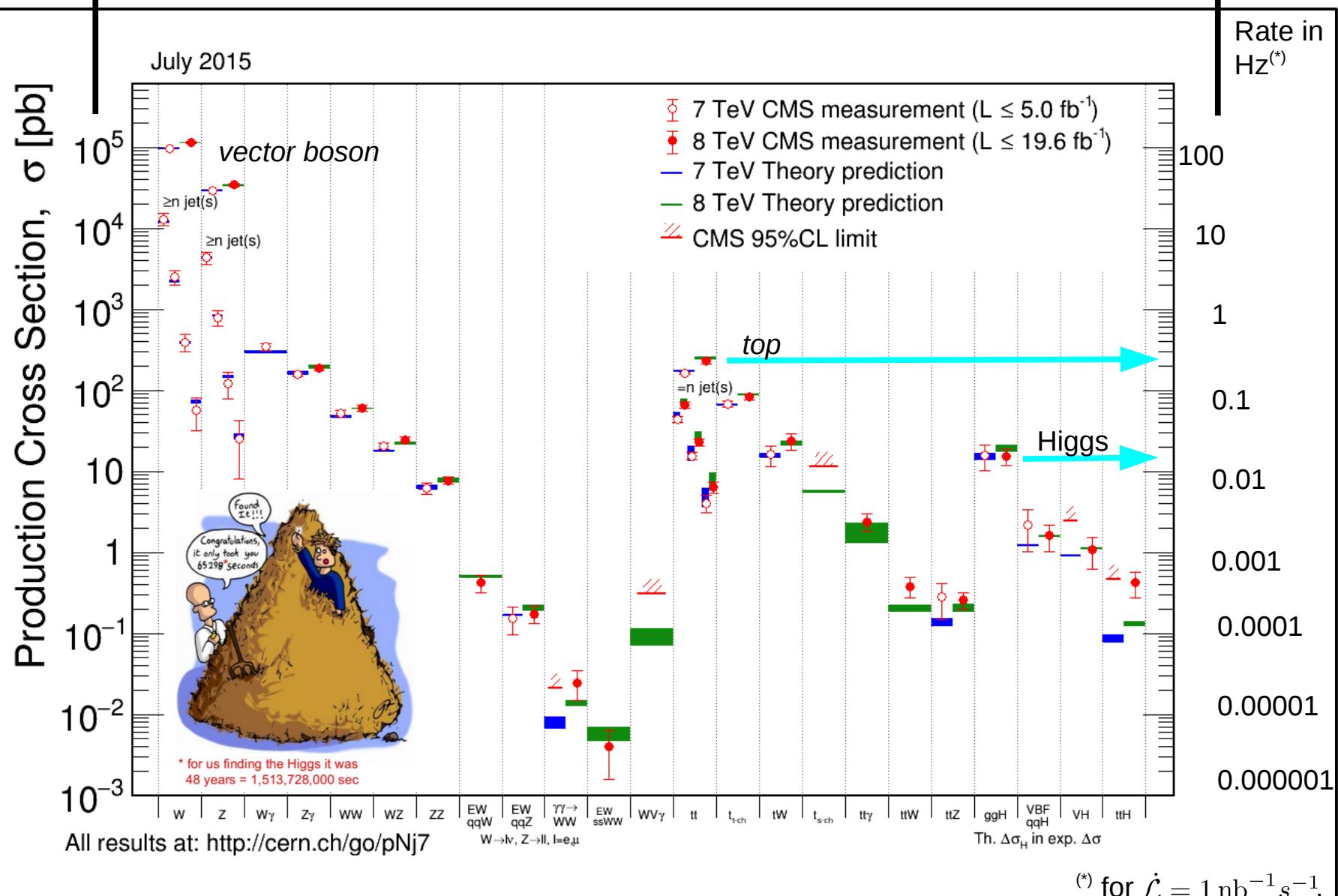


$$\kappa_{HVV} = \frac{2m_V^2}{v}$$

$$\kappa_{Hff} = \frac{m_f}{v}$$

	Channel	Resolution	S/B
$H \rightarrow \gamma\gamma$	1-2%	$\mathcal{O}(0.1)$	
$H \rightarrow ZZ$	1-2%	$\mathcal{O}(> 1)$	
$H \rightarrow WW$	20%	$\mathcal{O}(1)$	
$H \rightarrow b\bar{b}$	10%	$\mathcal{O}(0.1)$	
$H \rightarrow \tau\tau$	15%	$\mathcal{O}(0.1)$	

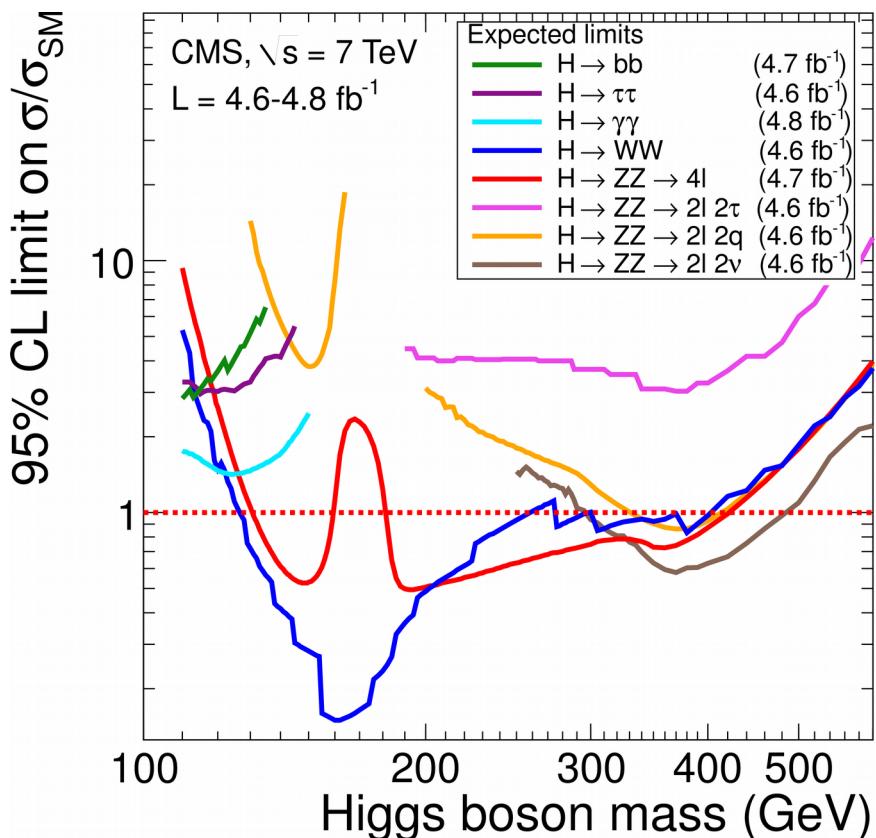
Die Herausforderung



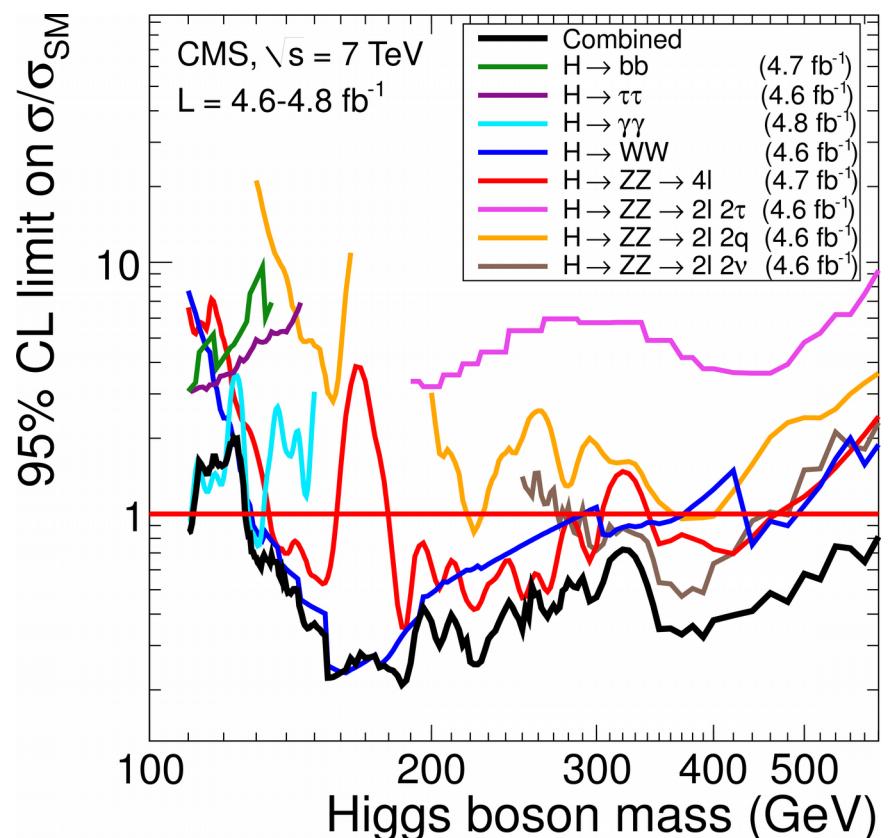
Verlauf der Suche bis Mitte 2012

Anm: ausgeschlossen sind alle Bereiche oberhalb der Kurven

- Ausschluß des SM bis auf **engen Massenbereich** des Higgs Bosons
- Erwartete (Ausschluß-)Sensitivität:



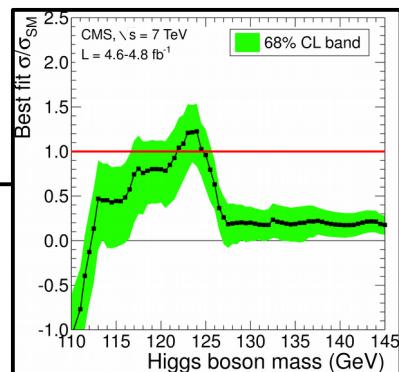
- Beobachteter Massenausschluß:



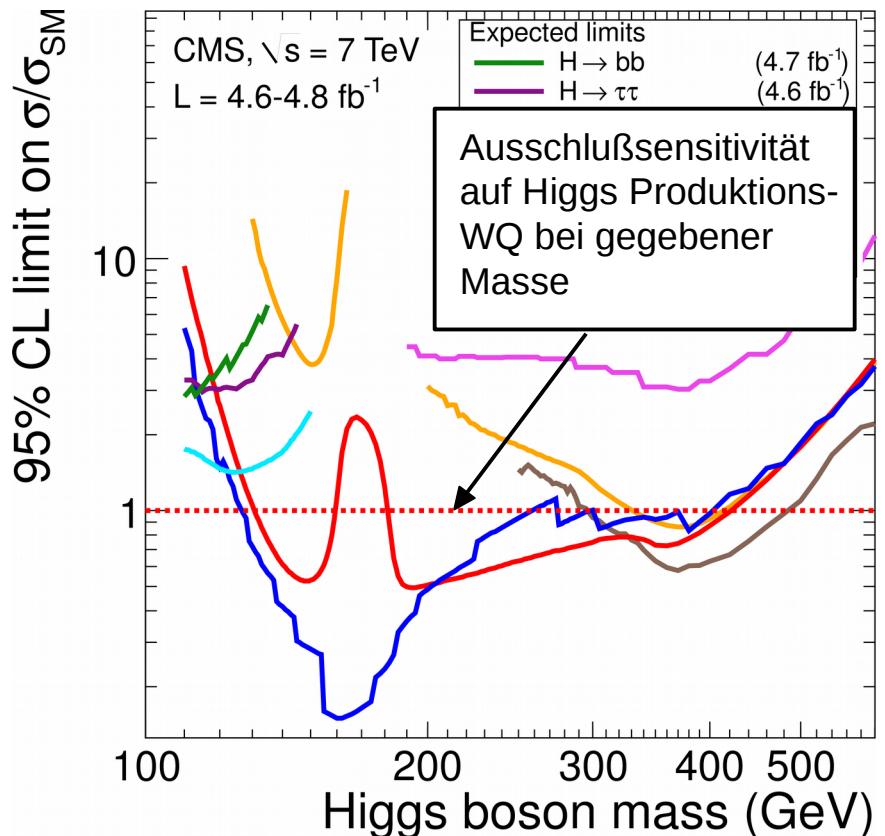
Verlauf bis Mitte 2012

- Ausschluß des SM bis auf **engen Massenbereich** des Higgs Bosons

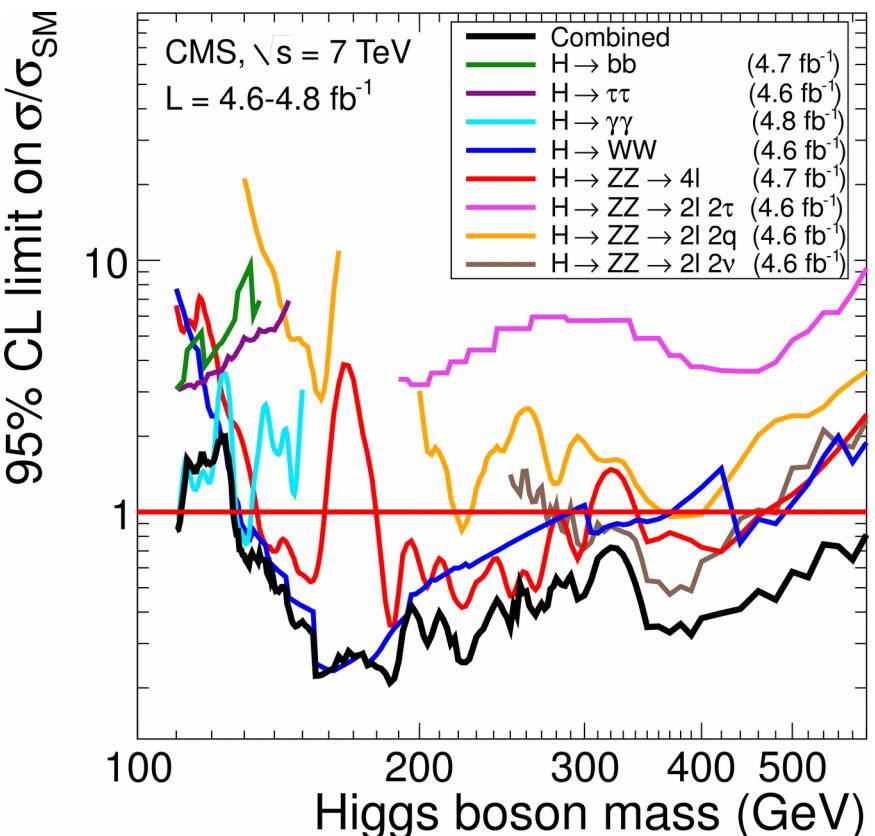
Signalstärke aus Anpassung an Daten



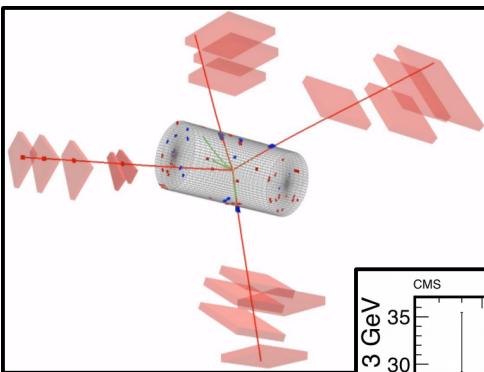
- Erwartete (Ausschluß-)Sensitivität:



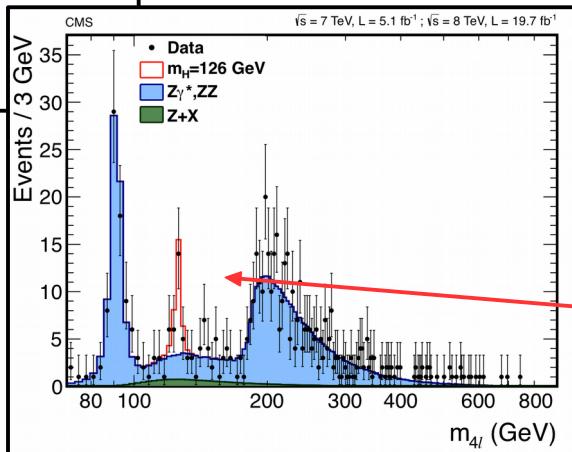
- Beobachteter Massenausschluß:



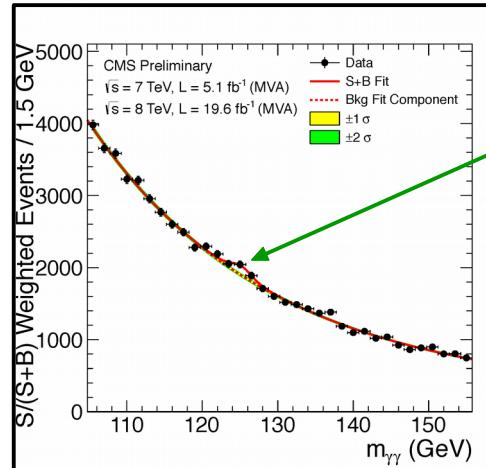
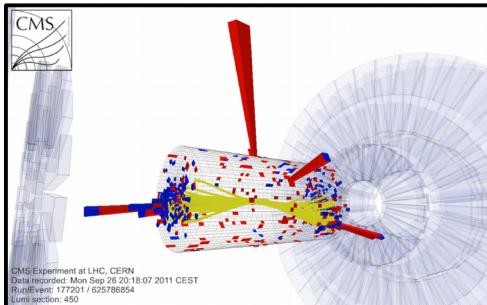
Die Entdeckung



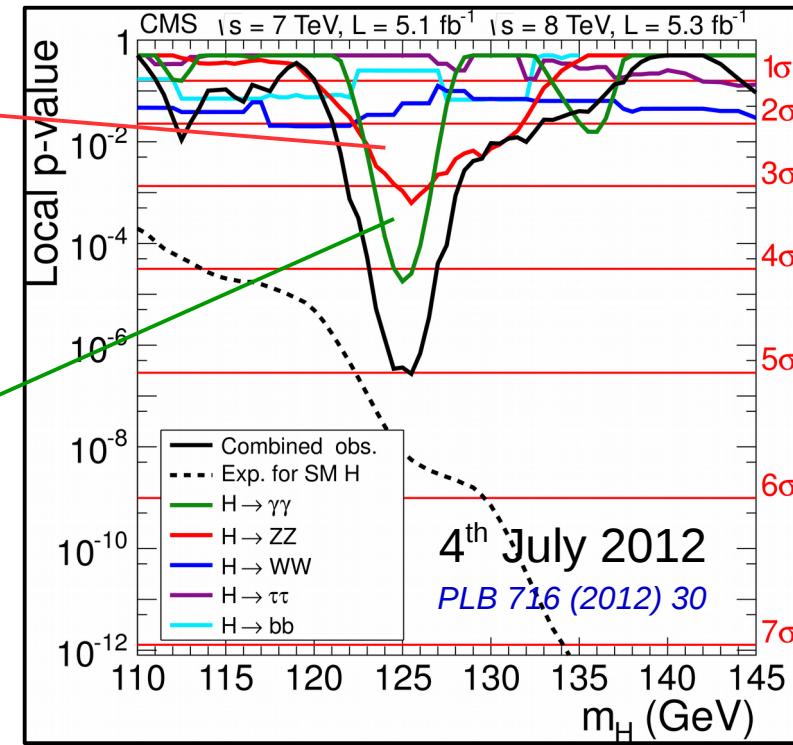
$H \rightarrow ZZ \rightarrow 4\ell$



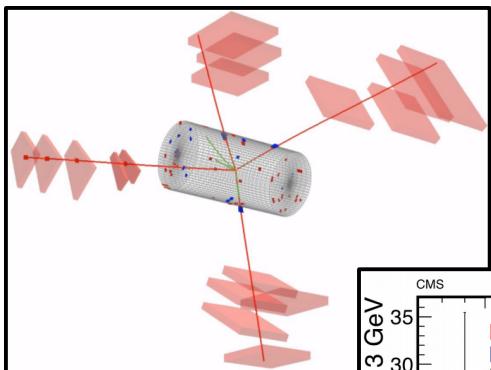
$H \rightarrow \gamma\gamma$



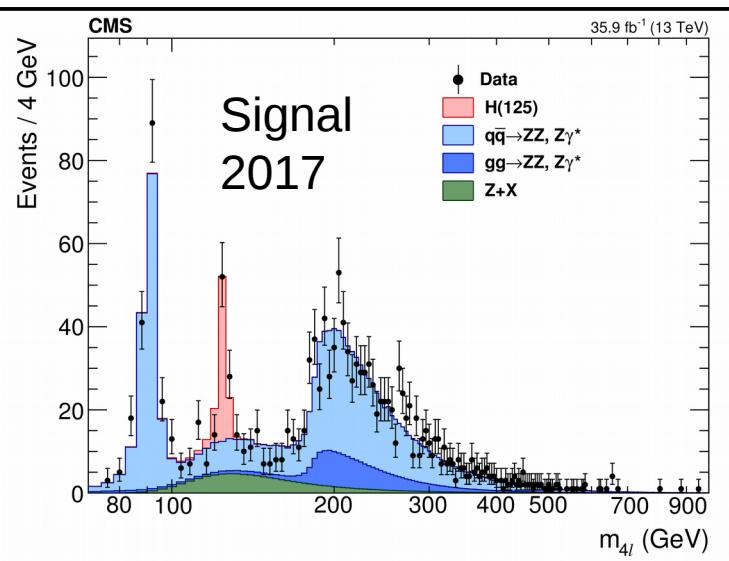
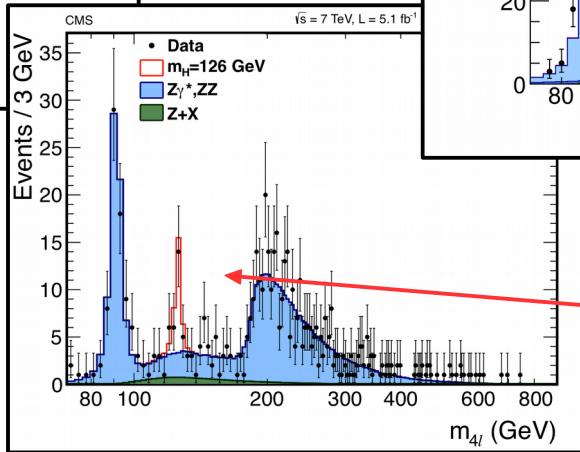
- **Klare Beobachtung** in bosonischen Zerfallskanälen
- Fermionische Kanäle noch nicht klar, weil noch nicht sensitiv genug (\rightarrow erfolgte etwas später)



Die Entdeckung

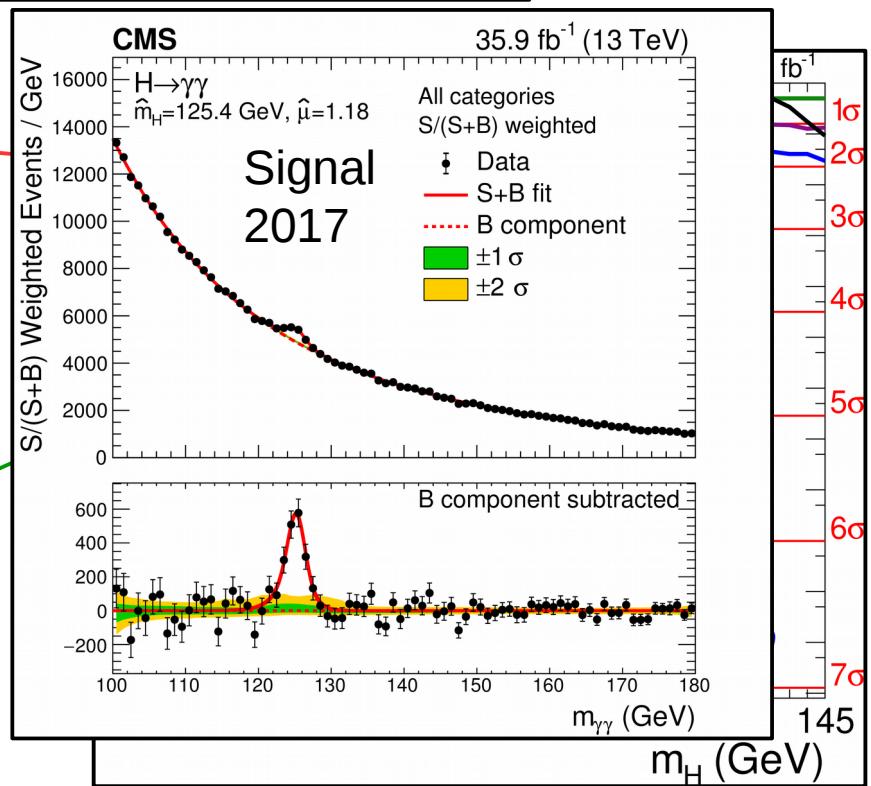
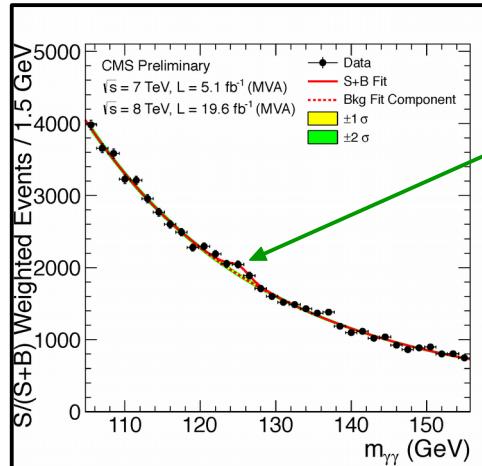
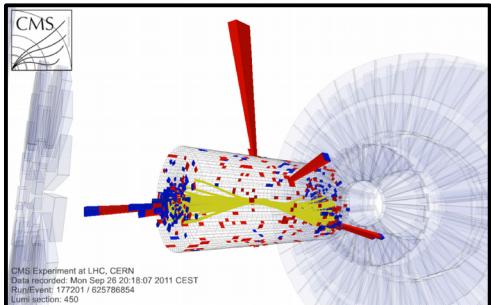


$H \rightarrow ZZ \rightarrow 4\ell$



bosonischen
noch nicht klar,
genug (\rightarrow)

$H \rightarrow \gamma\gamma$



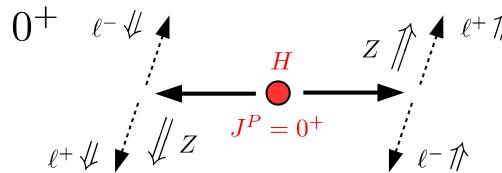
Spin & CP

- Golden decay channel:

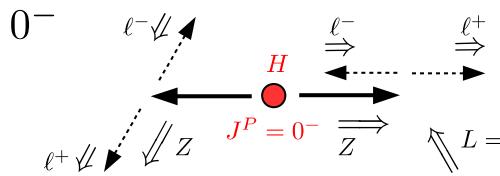
$$H \rightarrow ZZ \rightarrow 4\ell$$

$$P(Y_L^m(\theta, \varphi)) = (-1)^L \cdot Y_L^m(\theta, \varphi)$$

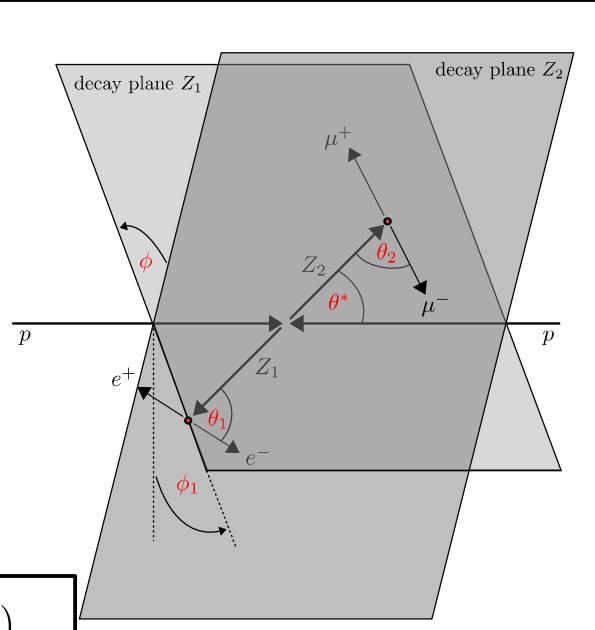
$$P(4\ell) = (-1)^L (-1)^2 (+1)^2 = (-1)^L$$



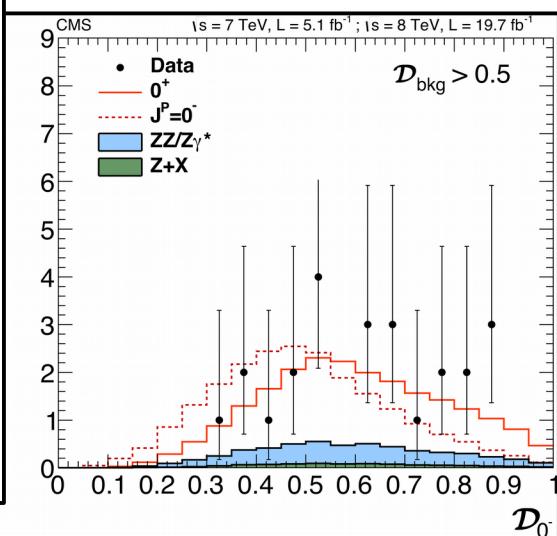
$$|0, 0\rangle = \sqrt{\frac{1}{3}}|1, -1\rangle \otimes |1, -1\rangle - \sqrt{\frac{1}{3}}|1, 0\rangle \otimes |1, 0\rangle + \sqrt{\frac{1}{3}}|1, -1\rangle \otimes |1, -1\rangle$$



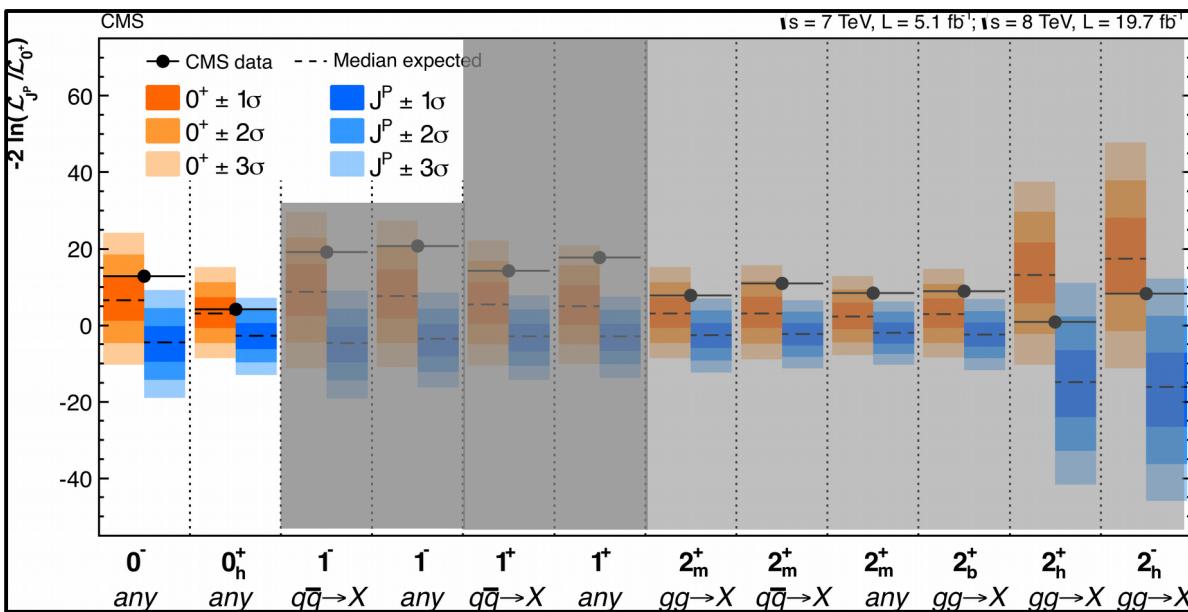
$$|1, \pm 1\rangle = \sqrt{\frac{1}{2}}|1, \pm 1\rangle \otimes |1, 0\rangle - \sqrt{\frac{1}{2}}|1, 0\rangle \otimes |1, \pm 1\rangle$$



PRD 89 (2014) 092007

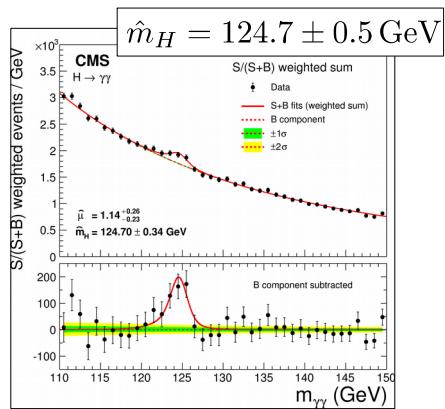


Test of pure spin hypotheses (based on $\mathcal{O}(50)$ evts):

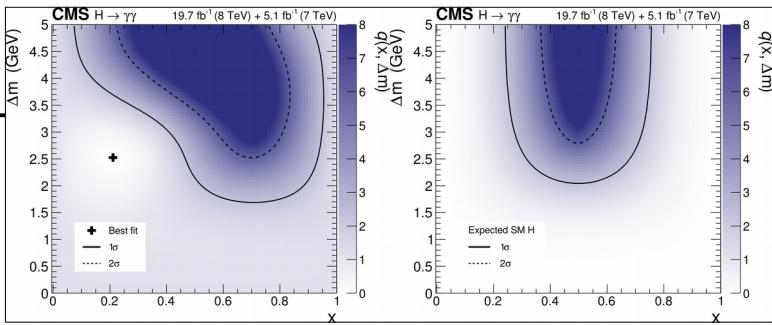


Compatibility

EPJC 74 (2014) 3076

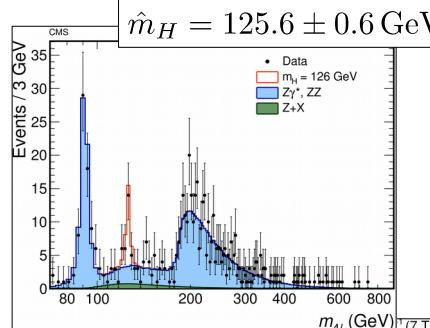


EPJC 74 (2014) 3076

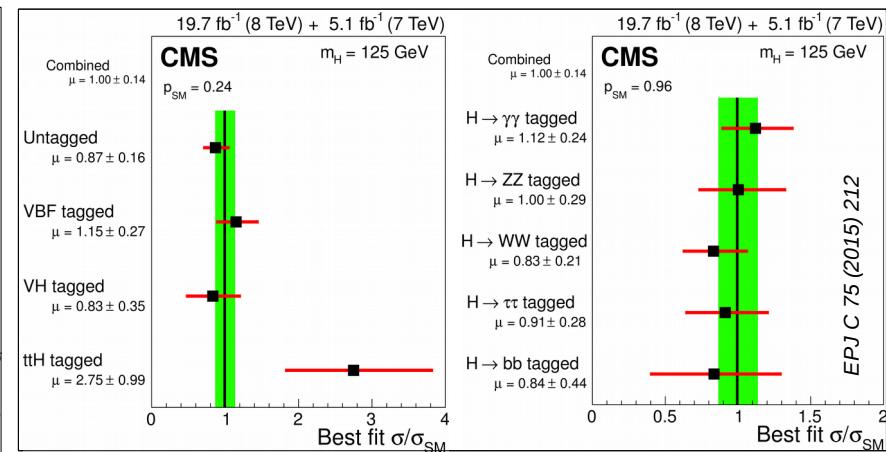
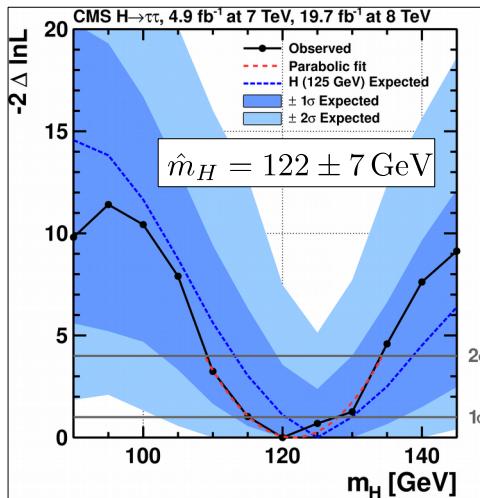


Coupling across production modes or decay channels:

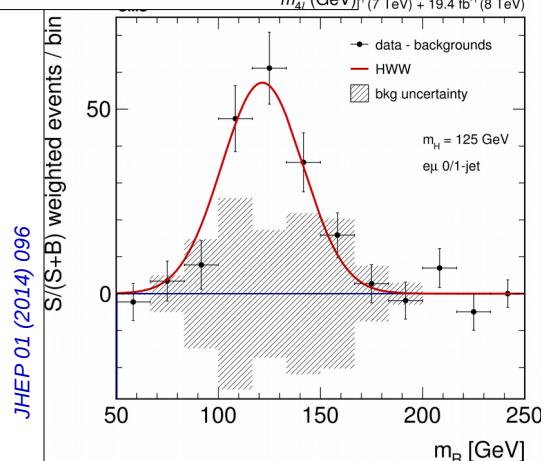
PRD 89 (2014) 092007



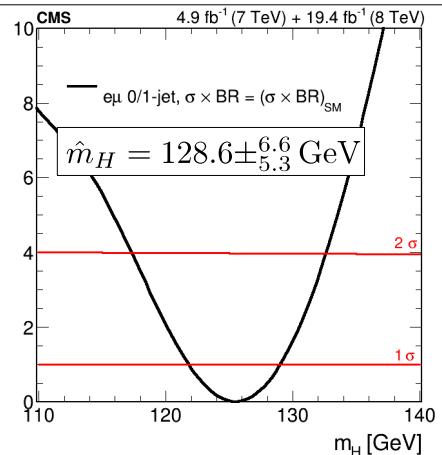
JHEP 05 (2014) 104



EPJC 75 (2015) 212



JHEP 01 (2014) 096



EPJC 75 (2015) 212

Overall coupling consistency:

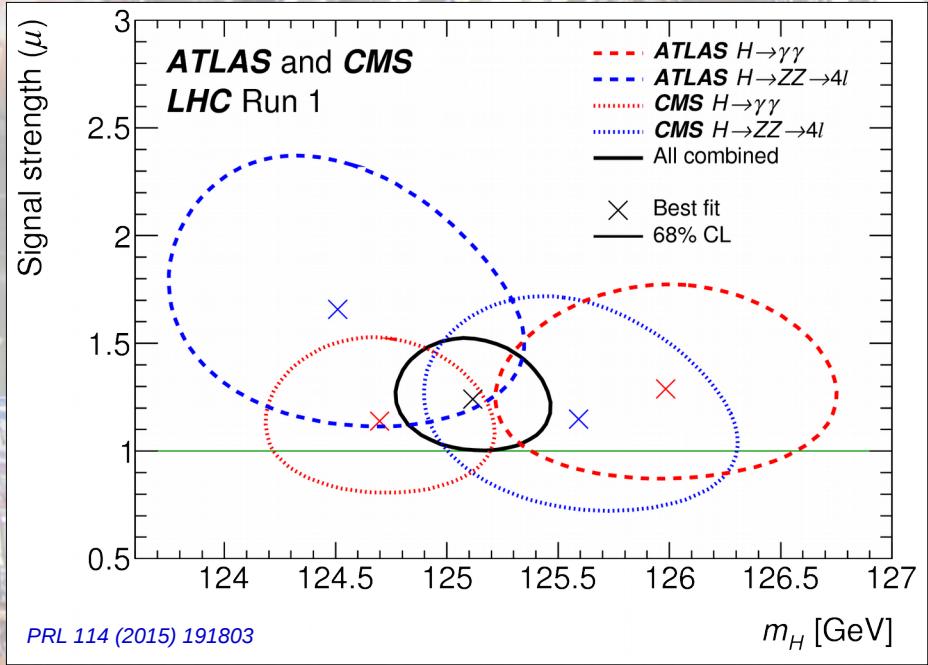
- Event categories : 227
- Nuisance parameters: $\mathcal{O}(2500)$
- 16 MB binary file of stat. model (~145 MB in human readable form).

$$\mu = \sigma/\sigma_{SM} = 1.00 \pm 0.14$$

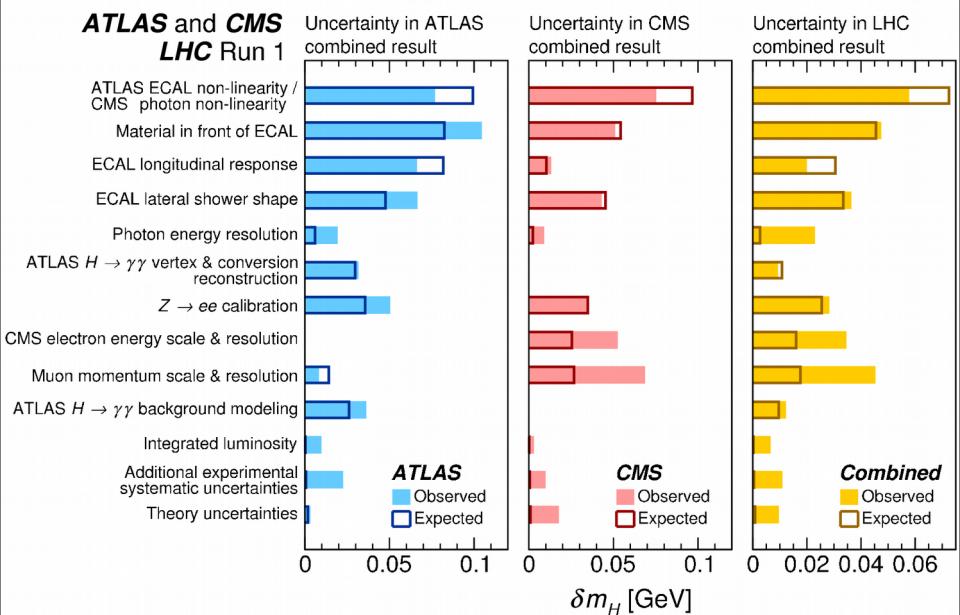
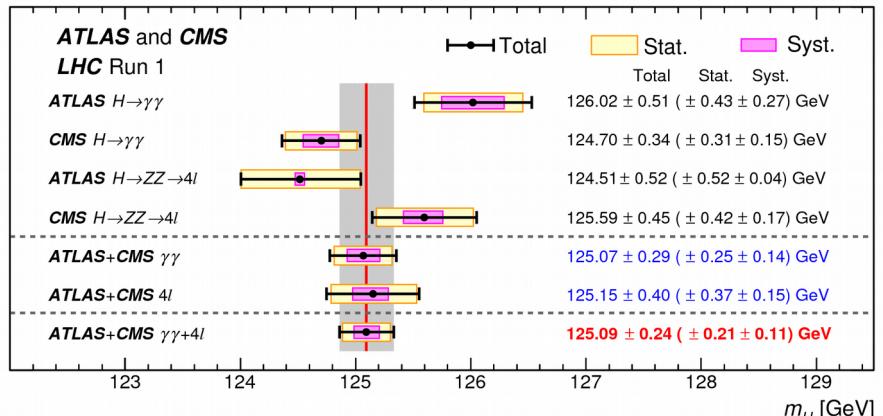
$$p\text{-value} = 84\%$$

Mass

- ATLAS+CMS LHC run-1 combination:



$125.06 \pm 0.21 \text{ (stat.)} \pm 0.19 \text{ (syst.)} \text{ GeV}$



- Event categories : 574
- Nuisance parameters: 4268
- $\mu = \sigma/\sigma_{SM} = 1.09 \pm 0.11$

Coupling structure

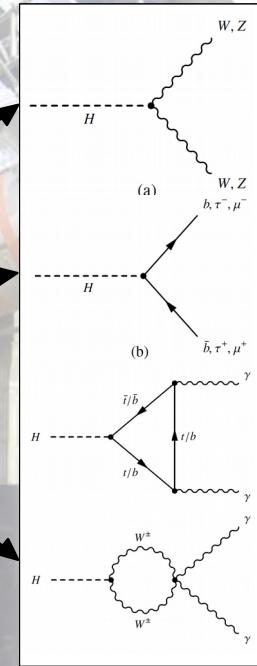
- ATLAS+CMS LHC run-1 combination:

Considered production modes:

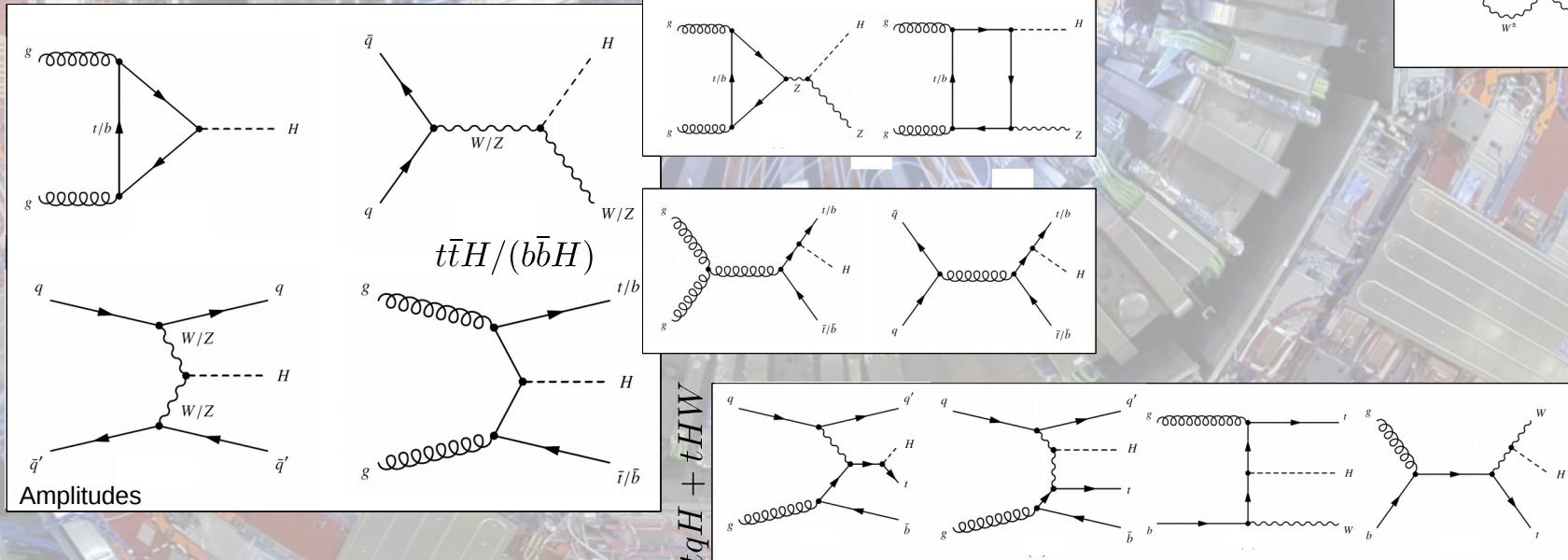
Production process	Cross section [pb]		Order of calculation
	$\sqrt{s} = 7 \text{ TeV}$	$\sqrt{s} = 8 \text{ TeV}$	
ggF	15.0 ± 1.6	19.2 ± 2.0	NNLO(QCD)+NLO(EW)
VBF	1.22 ± 0.03	1.58 ± 0.04	NLO(QCD+EW)+~NNLO(QCD)
WH	0.577 ± 0.016	0.703 ± 0.018	NNLO(QCD)+NLO(EW)
ZH	0.334 ± 0.013	0.414 ± 0.016	NNLO(QCD)+NLO(EW)
[ggZH]	0.023 ± 0.007	0.032 ± 0.010	NLO(QCD)
bbH	0.156 ± 0.021	0.203 ± 0.028	5FS NNLO(QCD) + 4FS NLO(QCD)
ttH	0.086 ± 0.009	0.129 ± 0.014	NLO(QCD)
tH	0.012 ± 0.001	0.018 ± 0.001	NLO(QCD)
Total	17.4 ± 1.6	22.3 ± 2.0	

Considered decay channels:

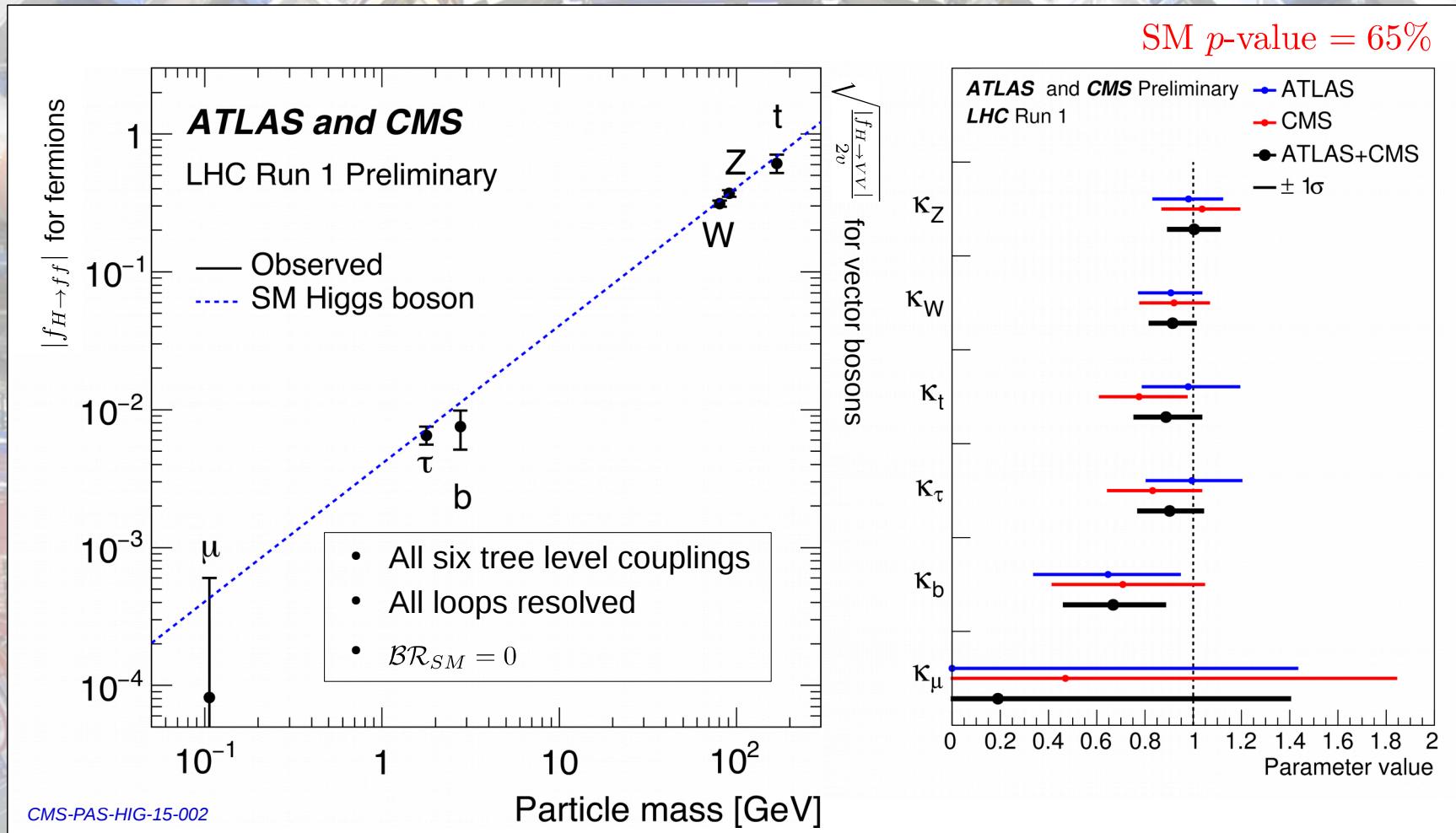
Decay channel	Branching ratio [%]
$H \rightarrow bb$	57.5 ± 1.9
$H \rightarrow WW$	21.6 ± 0.9
$H \rightarrow gg$	8.56 ± 0.86
$H \rightarrow \tau\tau$	6.30 ± 0.36
$H \rightarrow cc$	2.90 ± 0.35
$H \rightarrow ZZ$	2.67 ± 0.11
$H \rightarrow \gamma\gamma$	0.228 ± 0.011
$H \rightarrow Z\gamma$	0.155 ± 0.014
$H \rightarrow \mu\mu$	0.022 ± 0.001



Main production modes:



“Money plot”



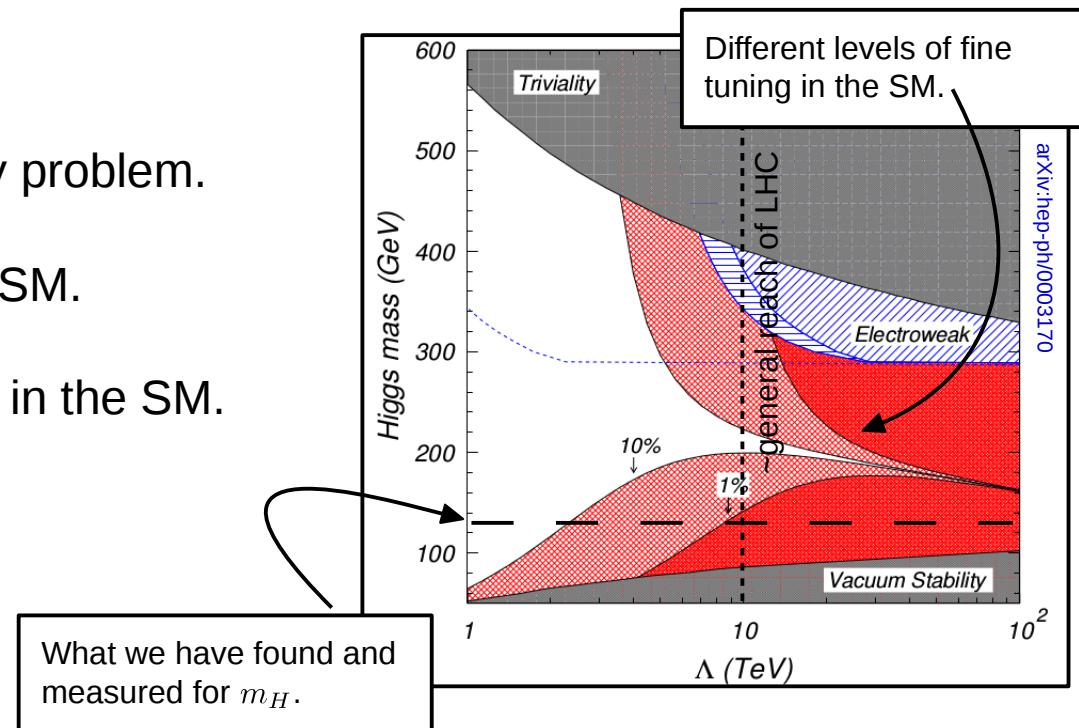
$$|f_{H \rightarrow ff}^{\text{obs}}| = \kappa_f \cdot |f_{H \rightarrow ff}^{\text{SM}}| = \kappa_f \cdot \frac{m_f}{v} \quad f = \mu, \tau, b, t$$

$$\sqrt{\frac{|f_{H \rightarrow VV}^{\text{obs}}|}{2v}} = \sqrt{\kappa_V} \cdot \sqrt{\frac{|f_{H \rightarrow VV}^{\text{SM}}|}{2v}} = \sqrt{\kappa_V} \cdot \frac{m_V}{v} \quad V = W, Z$$

Within measurement accuracy
unique scaling as expected
within the SM.

Why the Higgs boson still is not THE Higgs boson⁽¹⁾

- Gravity is not included in the SM.
- The SM suffers from the hierarchy problem.
- Dark matter is not included in the SM.
- Neutrino masses are not included in the SM.
- There are known deviations from the SM expectation in $a_\mu \equiv \frac{g_\mu - 2}{2}$ (3.6σ unresolved).



- There must be physics beyond the SM!
- At what scale does it set in?
- (How) Does it influence the Higgs sector?

⁽¹⁾ Arguments taken from S. Heinemeyer (HH Higgs workshop 2014)

Gliederung der Vorlesung

YOU
ARE
HERE

Vorlesung:	Vorlesungstag:	Übungsblatt:
VL-01 Einheiten, Relativistische Kinematik	Di 17.04.2018	-
VL-02 Teilchenstreuung	Do 19.04.2018	-
VL-03 Wirkungsquerschnitt	Di 24.04.2018	Blatt-01
VL-04 Teilchenbeschleunigung	Do 26.04.2018	-
Vorlesung fällt aus	Di 01.05.2018	Blatt-02
VL-05 Teilchennachweis durch Ionisation	Do 03.05.2018	-
VL-06 Elektromag. WW und Schauer	Di 08.05.2018	Blatt-03
Vorlesung fällt aus	Do 10.05.2018	-
VL-07 Detektoren der Teilchenphysik	Di 15.05.2018	Blatt-04
VL-08 Symmetrien und Erhaltungssätze	Do 17.05.2018	-
VL-09 Fundamentale Teilchen und Kräfte im SM	Di 22.05.2018	Blatt-05
VL-10 Diskrete Symmetrien des SM	Do 24.05.2018	-
VL-11 Teilchenzoo: vom Hadron zum Quark	Di 29.05.2018	Blatt-06
Vorlesung fällt aus	Do 31.05.2018	-
VL-12 Farbladung und QCD	Di 05.06.2018	Blatt-07
VL-13 Phänomenologie der schwachen WW	Do 07.06.2018	-
VL-14 Theorie der elektroschwachen WW	Di 12.06.2018	Blatt-08
VL-15 Higgs Mechanismus	Do 14.06.2018	-
VL-16 SM: Quarksektor	Di 19.06.2018	Blatt-09
VL-17 Top: Entdeckung und Eigenschaften	Do 21.06.2018	-
VL-18 Higgs: Entdeckung und Eigenschaften	Di 26.06.2018	-
VL-19 Neutrinosphysik	Do 28.06.2018	-

Backup