

Teilchenphysik 2 — W/Z/Higgs an Collidern

Sommersemester 2019

Matthias Schröder und Roger Wolf | Vorlesung 8

INSTITUT FÜR EXPERIMENTELLE TEILCHENPHYSIK (ETP)



Date	Room	Type	Topic
Wed Apr 24.	KI. HS B	LE 01	1. Organisation and introduction: particle physics at colliders + W/Z/H history
Tue Apr 30.	30.23 11/12	—	<i>no class</i>
Wed May 01.	KI. HS B	—	<i>no class</i>
Tue May 07.	30.23 11/12	LE 02	2.1 Gauge theory & 2.2 The electroweak sector of the SM I
Wed May 08.	KI. HS B	LE 03, EX 01	2.3 Discovery of the W and Z bosons & EX gauge theories
Tue May 14.	30.23 11/12	LE 04	2.4 The Higgs mechanism
Wed May 15.	KI. HS B	EX 02	Exercise “SM Higgs mechanism”
Tue May 21.	30.23 11/12	—	<i>no class</i>
Wed May 22.	KI. HS B	LE 05	2.5 The electroweak sector of the SM II (Higgs mechanism + Yukawa couplings)
Tue May 28.	30.23 11/12	SP 01	Specialisation of 2.4 and 2.5
Wed May 29.	KI. HS B	LE 06	3.1 From theory to observables & 3.2 Reconstruction + analysis of exp. data
Tue Jun 04.	30.23 11/12	EX 03	Exercise “Trigger efficiency measurement”
Wed Jun 05.	KI. HS B	LE 07	3.3 Measurements in particle physics (part 1)
Tue Jun 11.	30.23 11/12	EX 04	Exercise on statistical methods
Wed Jun 12.	KI. HS B	LE 08	3.3 Measurements in particle physics (part 2)
Tue Jun 18.	30.23 11/12	SP 02	Specialisation “Limit setting”
Wed Jun 19.	KI. HS B	SP 03	Specialisation “Unfolding”
Tue Jun 25.	30.23 11/12	LE 09	4.1 Determination of SM parameters
Wed Jun 26.	KI. HS B	LE 10	4.2 Measurement and role of W/Z bosons at the LHC
Tue Jul 02.	30.23 11/12	EX 05	Paper seminar “Z pole measurements”
Wed Jul 03.	KI. HS B	LE 11	4.3 Processes with several W/Z bosons
Tue Jul 09.	30.23 11/12	EX 06	Paper seminar Higgs
Wed Jul 10.	KI. HS B	LE 12	5.1 Discovery and first measurements of the Higgs boson
Tue Jul 16.	30.23 11/12	EX 07	Exercise “Machine learning in physics analysis”
Wed Jul 17.	KI. HS B	LE 13	5.2 Measurement of couplings and kinematic properties
Tue Jul 23.	30.23 11/12	EX 08	Presentations: results of ML challenge
Wed Jul 24.	KI. HS B	LE 14	5.3 Search for Higgs physics beyond the SM & 5.4 Future Higgs physics

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Analysis Chain

Nature



Theory



Detector: data recording
calibrated digitised data
online selection (trigger)

MC simulation
physics process
detector signals



Physics object **reconstruction**
Event **selection**



Statistical analysis: **results**
Comparison with theory

3. From Theory to Experiment (and Back)

3.1 From theory to observables

- Cross-section calculation: basic picture
- Fermion propagator and perturbation theory
- Scattering matrix and Feynman rules

3.2 Reconstruction of experimental data

- Reminder: accelerators and particle detectors
- Trigger
- Reconstruction of physics objects

3.3 Measurements in particle physics

- Parameter estimation
- Hypothesis testing
- Search for new physics (exclusion limits)

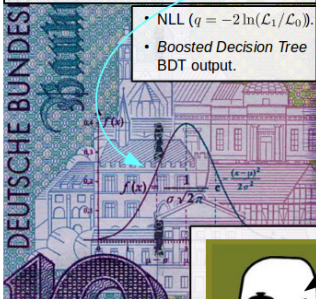
3.4 Monte Carlo simulation

3.3 Measurements in particle physics

Test statistic:

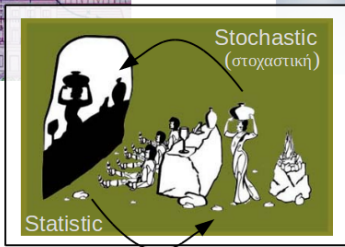
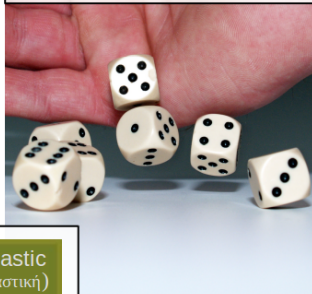
$$\Omega^n \rightarrow \mathbb{R} : x \rightarrow f(x)$$

- NLL ($q = -2 \ln(\mathcal{L}_1/\mathcal{L}_0)$).
- *Boosted Decision Tree* BDT output.



Probability (density) function:

$$\Omega^n \rightarrow [0, 1] \subset \mathbb{R} : x \rightarrow \mathcal{P}(x)$$



- Problem of statistics is usually *ill-defined*.
- Deduce **truth** from **shadows** in Platon's cave...

From R. Wolf

Experiment

- All measurements are derived from **counting experiments**
- **Millions of billions** of particle collision events each year
- Each event **independent** from and (to many extents) identical to all others

Theory

- Nature intrinsically stochastic
- QM wave functions interpreted as **probability density functions**
- Event-by-event simulation using **Monte Carlo** methods

Particle physics experiments perfect application of statistics

- R. J. Barlow, “Statistics: A Guide to the Use of Statistical Methods in the Physical Sciences”, The Manchester Physics Series **highly recommended**
- G. Cowan, “Statistical Data Analysis”, Oxford Science Publications
- Review “Statistics”, Chin. Phys. C, 40, 100001 (2016) **very good review**
- The ATLAS and CMS Collaborations, “Procedure for the LHC Higgs boson search combination in Summer 2011”, CMS-NOTE-2011-005, ATL-PHYS-PUB-2011-11 **expert document**

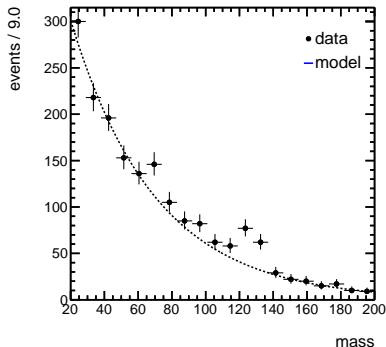
- **Likelihood** $\mathcal{L}(n_{\text{obs}}; k) = \prod_i \mathcal{P}(n_{\text{obs},i}; k)$ **quantifies compatibility of the observed data with a given model**

- i : independent counting experiments, e. g. bins of a histogram
- $n_{\text{obs}} = \{n_{\text{obs},i}\}$: number of observed events in i , distributed as \mathcal{P}
- $k = \{k_j\}$: model parameters in \mathcal{P}

→ **function of model parameters and observed data**

- **Example**
 - Data: observed events in 20 bins
 - Model: known SM process (background)

$$\mathcal{L}(n_{\text{obs}}; b) = \prod_i \text{Poisson}(n_{\text{obs},i}, b)$$



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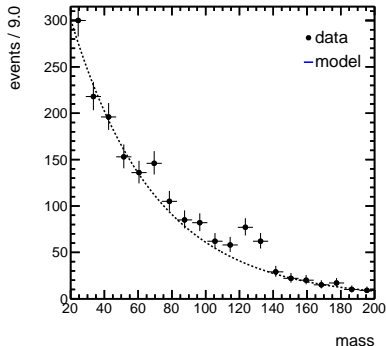
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- **Example**
 - Data: observed events in 20 bins
 - Model: known SM process (background)

$$\mathcal{L}(n_{\text{obs}}; b) = \prod_i \frac{(n_{\text{pred},i})^{n_{\text{obs},i}}}{n_{\text{obs},i}!} e^{-n_{\text{pred},i}}$$

$$n_{\text{pred},i}(b) = \underbrace{b \cdot e^{-\alpha X_i}}_{\text{background}}$$



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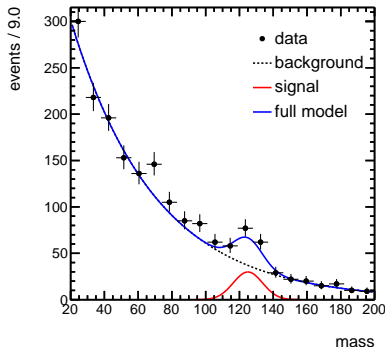
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→ **function of model parameters and observed data**

- **Example**
 - Data: observed events in 20 bins
 - Model: known SM process (background) + **signal**

$$\mathcal{L}(n_{\text{obs}}; b, s) = \prod_i \frac{(n_{\text{pred},i})^{n_{\text{obs},i}}}{n_{\text{obs},i}!} e^{-n_{\text{pred},i}}$$

$$n_{\text{pred},i}(b, s) = \underbrace{b \cdot e^{-\alpha x_i}}_{\text{background}} + \underbrace{s \cdot e^{-(m-x_i)^2}}_{\text{signal}}$$

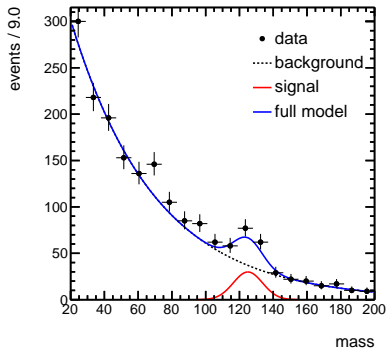


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→ **function of model parameters and observed data**

- \mathcal{L} is **not the probability** of a model!
 - Also not the probability of observing the data given the model
- \mathcal{L} is the product of pdfs of the model, used to quantify the compatibility of the data with the model



3.3.1. Parameter estimation

- Which parameters of a given model describe best the data?
- **Maximum likelihood (ML) method** (“maximum likelihood fit”):
Which parameter values lead to the maximum of \mathcal{L} ?

→ **Maximum likelihood estimators $\hat{\mathbf{k}}$** of the true values

$$\max [\mathcal{L}(n_{\text{obs}}, k)] = \mathcal{L}(n_{\text{obs}}, \hat{k})$$

- Technically: **minimisation of negative log-likelihood**

$$\text{NLL}(n_{\text{obs}}, k) = -2 \ln \mathcal{L}$$

- Logarithm is monotonic: retains minimum
- Turns product of pdfs into sum: easier to minimise
- Arbitrarily complex problem, vast amount of techniques and literature
 - e. g. ATLAS+CMS Higgs couplings combination: ≈ 4250 parameters
 - e. g. CMS tracker-alignment problem: $\mathcal{O}(10^6)$ parameters

Maximum-Likelihood Method

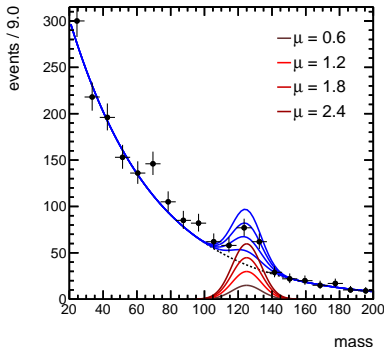
- Example: which *signal-strength modifier* μ describes best the data?
 - Common in Higgs physics: $\mu = (\sigma \cdot \mathcal{B}) / (\sigma_{\text{SM}} \cdot \mathcal{B}_{\text{SM}})$

$$n_{\text{pred},i}(\mu) = \underbrace{b \cdot e^{-\alpha x_i}}_{\text{background}} + \underbrace{\mu \cdot s_{\text{SM}} \cdot e^{-(m-x_i)^2}}_{\text{signal}}$$

$$\mathcal{L}(n_{\text{obs}}; \mu) = \prod_i \frac{(n_{\text{pred},i})^{n_{\text{obs},i}}}{n_{\text{obs},i}!} e^{-n_{\text{pred},i}}$$

→

$\hat{\mu}$: maximises $\mathcal{L}(n_{\text{obs}}; \mu)$



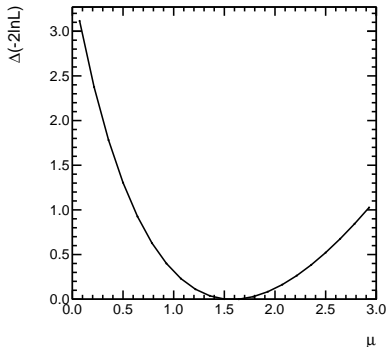
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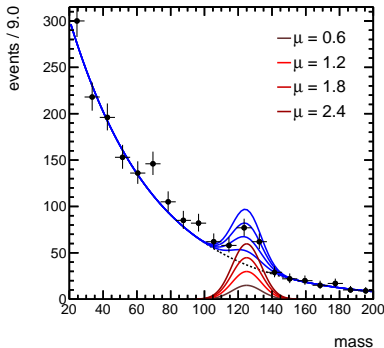
→ $\hat{\mu}$: minimises $-2 \ln \mathcal{L}(n_{\text{obs}}; \mu)$



- Example: which *signal-strength modifier* μ describes best the data?
 - Common in Higgs physics: $\mu = (\sigma \cdot \mathcal{B}) / (\sigma_{\text{SM}} \cdot \mathcal{B}_{\text{SM}})$
- Important case: **Gaussian distributed measurements**
 - e. g. approximation of Poisson for large number of events (in practice > 10)
- NLL becomes

$$-2 \ln \mathcal{L} = \sum_i \frac{(n_{\text{obs},i} - n_{\text{pred},i})^2}{n_{\text{pred},i}} + \text{const}$$

$$\chi^2 \equiv \sum_i \frac{(n_{\text{obs},i} - n_{\text{pred},i})^2}{n_{\text{pred},i}}$$

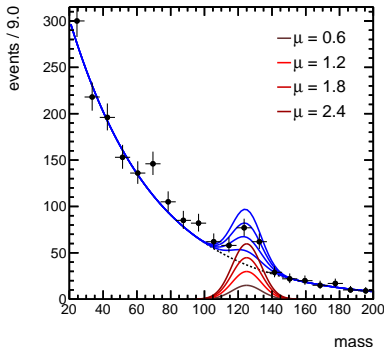


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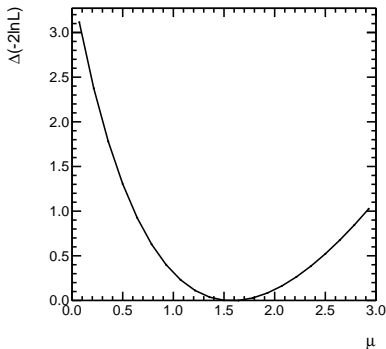
ML method obtains parameter values most compatible with the data *for a given model* — **a ML fit does not find the correct model!**

ML estimator is function of the observed data



- Uncertainty on $\hat{\mu}$ from scan of $\text{NLL} = -2 \ln \mathcal{L}$ around minimum
- **Uncertainty due to fluctuations of data:** “statistical uncertainty”
- For Gaussian pdfs: parabola, standard deviation σ follows from

$$\Delta(\text{NLL}) = \text{NLL}(\hat{\mu} \pm r \cdot \sigma) - \text{NLL}(\hat{\mu}) = r$$

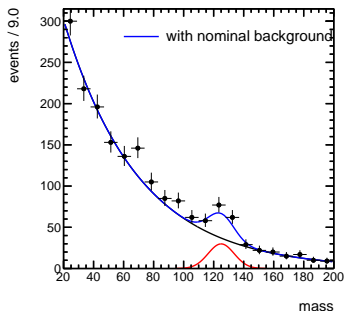


- Other cases can often be approximated by Gaussian case (and true for $n \rightarrow \infty$)
- For strongly asymmetric NLL
 - Often asymmetric intervals quoted
 - Better: variable transformation such that NLL symmetric

- **Uncertainty** reflects degree of precision with which one can deduce a parameter value from the data if one does everything correctly
 - Statistical uncertainties stem from the inherent stochastic nature of the data and finite number of observed events (can not be eliminated, only minimised)
 - Systematic uncertainties stem e. g. from the limited knowledge of the precision of the detector or approximations in theory calculations
 - Uncertainties can (in principle) be quantified
- In contrast, **errors are mistakes**
 - e. g. a lose cable or a wrong method
 - Errors can (and should) be eliminated
- In context of statistical data analysis, we mean uncertainties
 - NB: “Error bars” denote uncertainties!

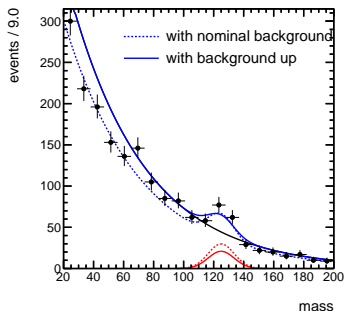
Incorporation of Systematic Uncertainties

- Often, background (and signal) models subject to systematic uncertainties
 - Some (theory or experimental) parameters not exactly known, e. g. cross section or trigger efficiency
- For example, background normalisation b not precisely known but with some uncertainty



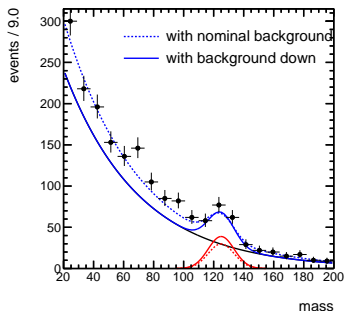
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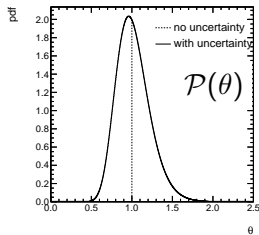
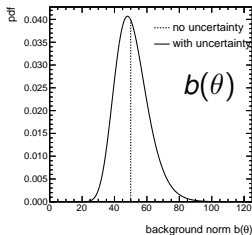
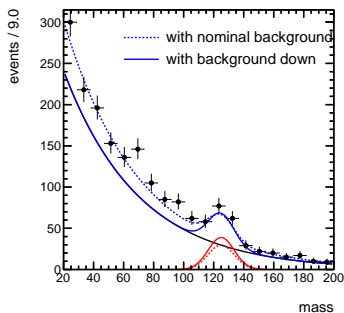


Incorporation of Systematic Uncertainties

- Incorporated into likelihood via **nuisance parameters** θ

$$\mathcal{L}(\mathbf{n}_{\text{obs}}; \mu, \theta) = \prod_i \mathcal{P}(n_{\text{obs}i}; \mu, \theta) \cdot \mathcal{P}_{\tilde{\theta}}(\tilde{\theta}|\theta)$$

- Background normalisation becomes function of θ : $b \rightarrow b(\theta)$
 - θ : assumed true value, parameter of the fit
 - $\tilde{\theta}$: best knowledge, e. g. estimate from independent measurement, distributed as $\mathcal{P}_{\tilde{\theta}}(\tilde{\theta}|\theta)$

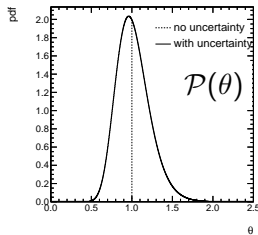
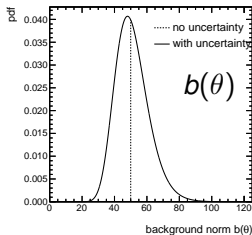
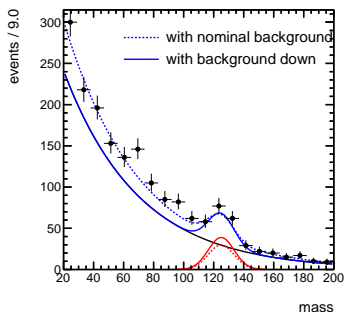


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- Background normalisation becomes function of θ : $b \rightarrow b(\theta)$
- ML fit can adjust θ to a value different than $\tilde{\theta}$ to achieve better description of data but at the cost of reducing value of $\mathcal{P}_{\tilde{\theta}}(\tilde{\theta}|\theta)$



Incorporation of Systematic Uncertainties

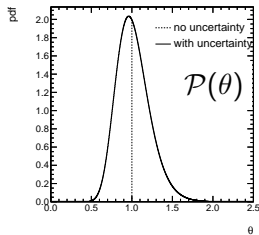
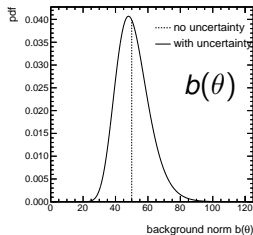
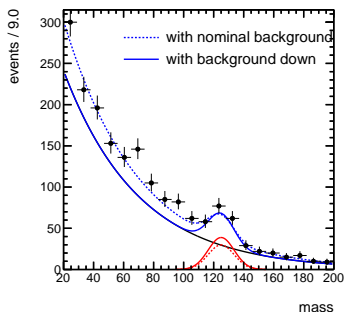
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- Nomenclature

- μ : *parameter of interest* (POI)
- θ : *nuisance parameter* (NP)

generally, (anti-)correlated with μ : increased uncertainty on $\hat{\mu}$



Incorporation of Systematic Uncertainties

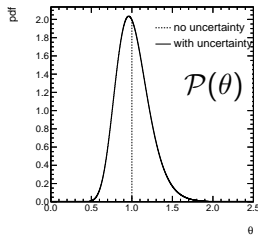
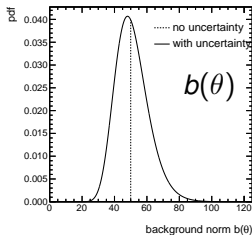
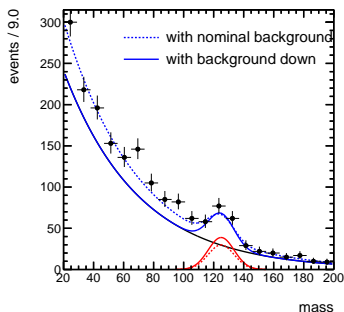
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- Nomenclature

- μ : *parameter of interest* (POI)
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generally, (anti-)correlated with μ : smaller sensitivity



Incorporation of Systematic Uncertainties

- Incorporated into likelihood via **nuisance parameters** θ

$$\mathcal{L}(\mathbf{n}_{\text{obs}}; \mu, \theta) = \prod_i \mathcal{P}(n_{\text{obs}i}; \mu, \theta) \cdot \mathcal{P}_{\tilde{\theta}}(\tilde{\theta}|\theta)$$

- Nomenclature
 - μ : *parameter of interest* (POI)
 - θ : *nuisance parameter* (NP)
- \mathcal{L} function of μ and θ , but only interested in μ
→ can rewrite θ as function of μ

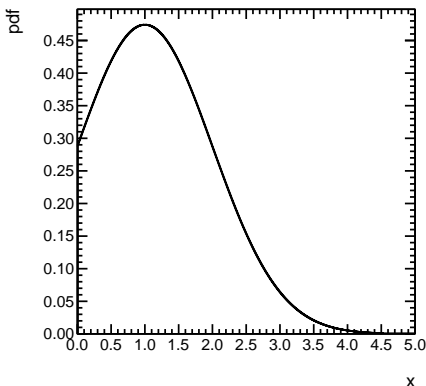
Profile likelihood $\mathcal{L}_p(\mu) \equiv \mathcal{L}(\mu, \hat{\theta}(\mu))$

$\hat{\theta}(\mu)$ maximises \mathcal{L} for given μ (“profiled values” of θ)

- \mathcal{L}_p in practice computationally advantageous
 - e. g. need only to consider μ instead of full parameter space (μ, θ) when computing uncertainties

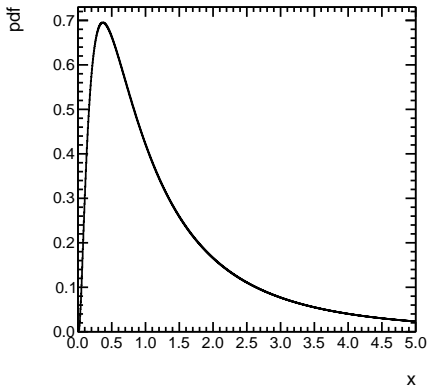
Typical Choices for $\mathcal{P}_\theta(\theta|\tilde{\theta})$

(truncated) Gaussian



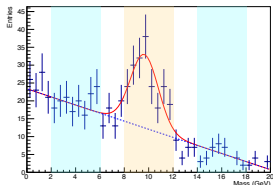
$$\frac{1}{\sqrt{2\pi}\sigma} e^{-\frac{1}{2}\left(\frac{x-\mu}{\sigma}\right)^2}$$

Log-normal

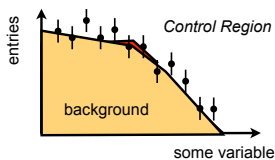
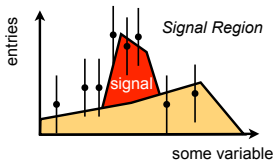


$$\frac{1}{\sqrt{2\pi}\sigma} \frac{1}{x} e^{-\frac{1}{2}\left(\frac{\ln(x)-\mu}{\sigma}\right)^2}$$

sidebands



signal/control region

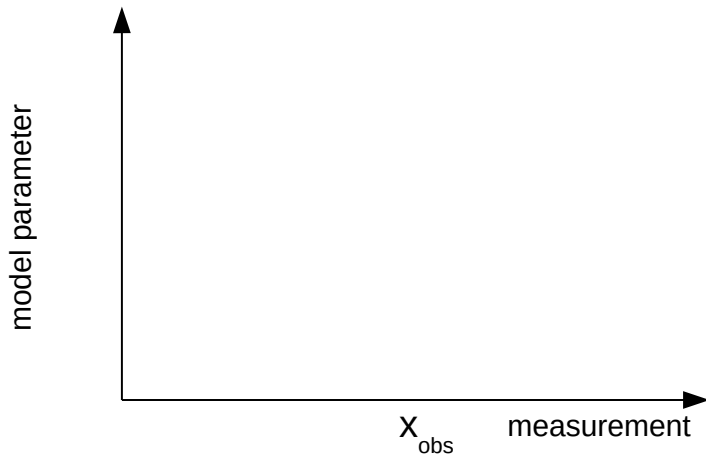


- Signal-depleted sideband/control region
 - Determine **rate** of background process
 - Determine **shape** of background processes
 - **Constrain uncertainties** on background prediction
- Can be incorporated into likelihood fit via **nuisance parameters**:
simultaneous determination of POI and background parameters

- Experiment to determine value of some parameter x measures x_{obs}
 - e. g. from maximum likelihood estimate, tag-and-probe measurement
- Want to quote “uncertainty”: interval that reflects statistical precision of measurement
 - Can be estimate of standard deviation from likelihood fit
 - More generally, in particular if non-Gaussian \mathcal{P} : **confidence interval**
- **Confidence interval** covers on average the true value with a given probability
 - e. g. *90 % confidence level*: if experiment repeated many time, the interval covers the true value in 90 % of the cases
 - Careful: this is not the probability of the true value to lie within the interval (there is only one true value either within or not)!

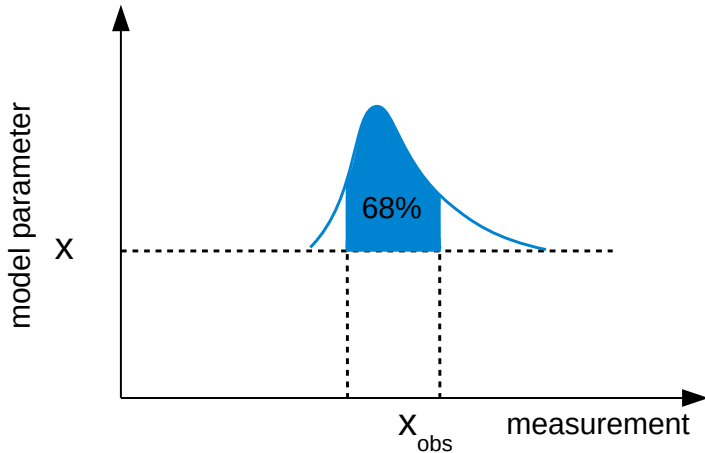
Confidence Intervals

Neyman construction



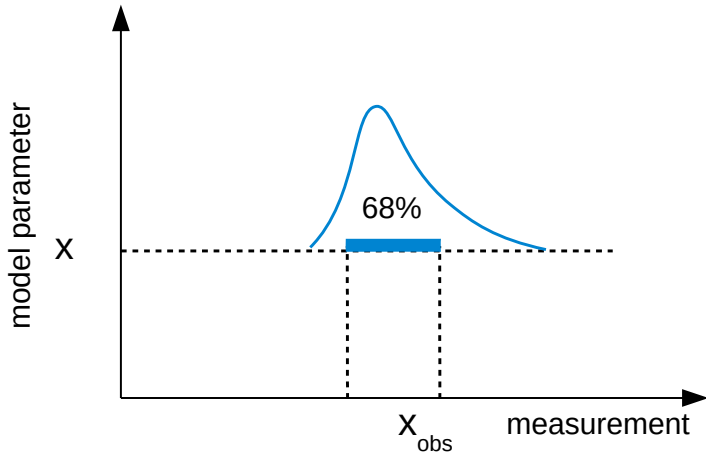
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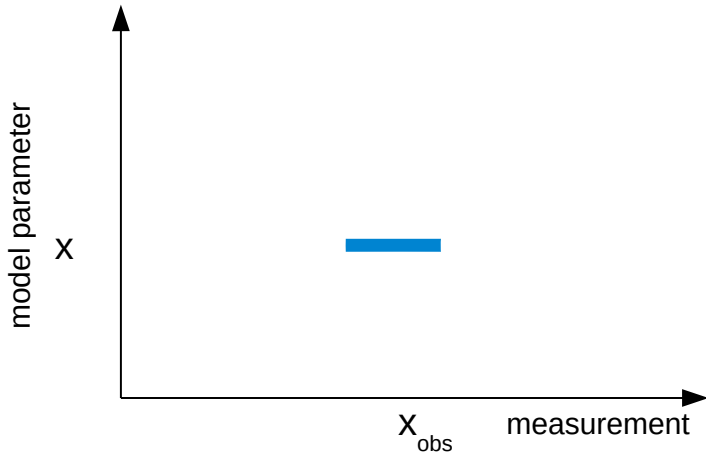
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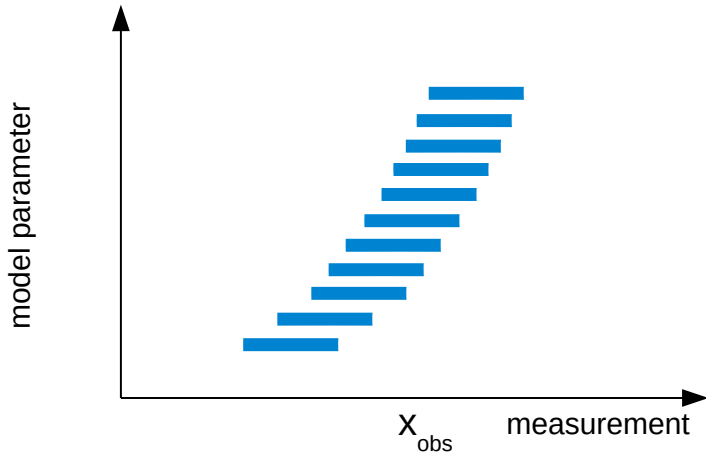
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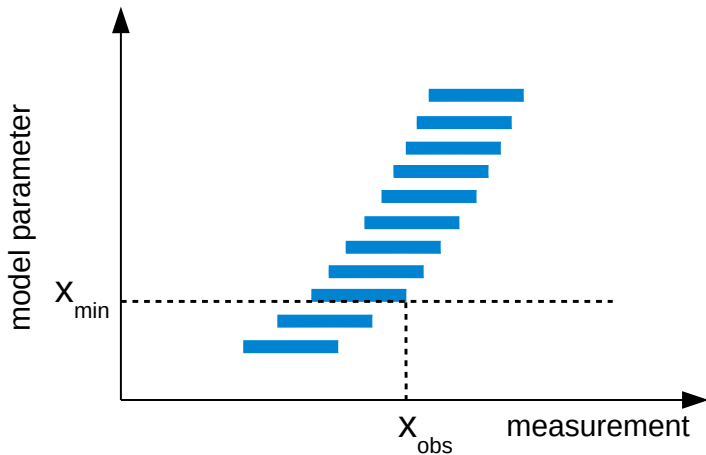
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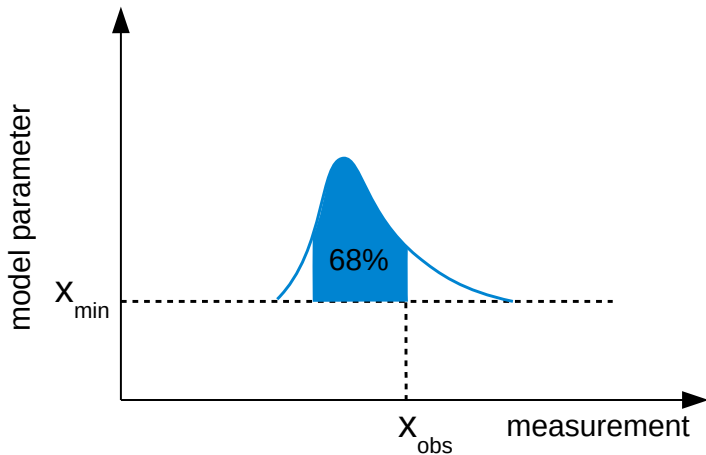
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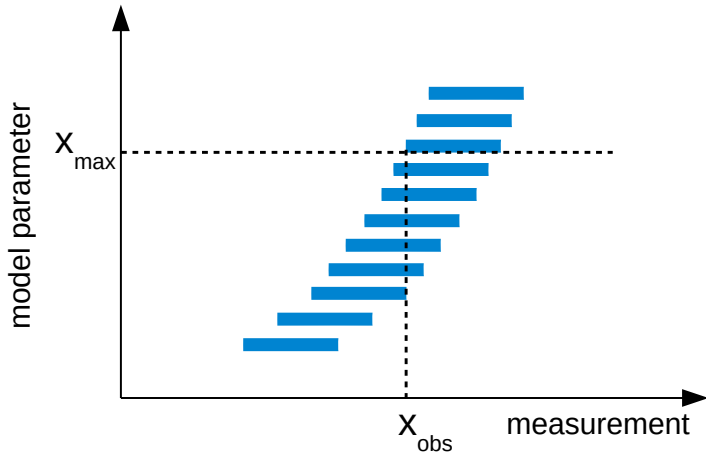
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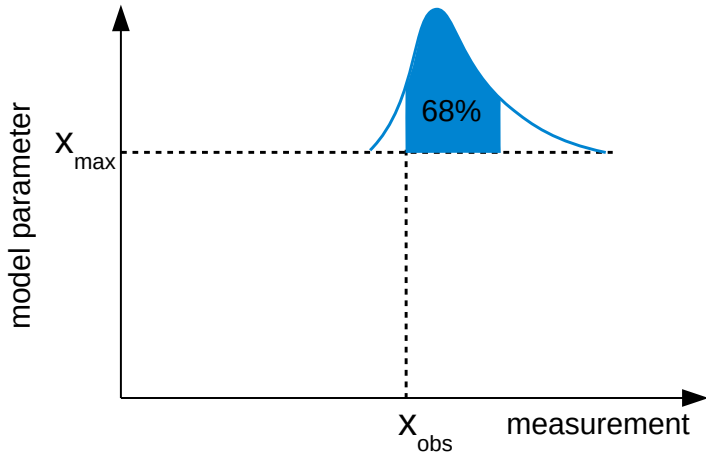
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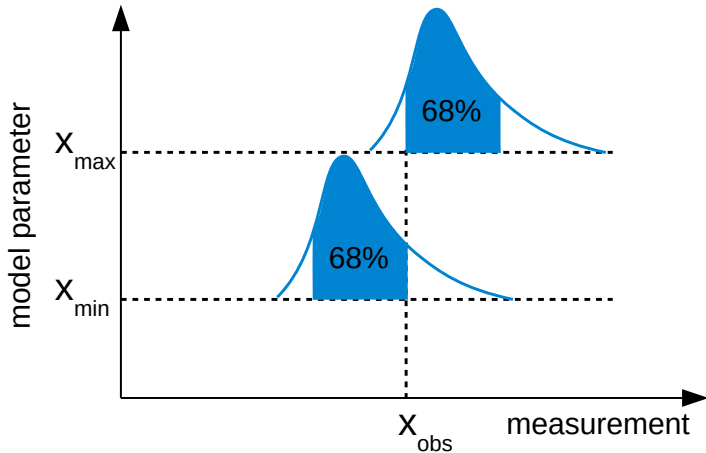
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Confidence Intervals

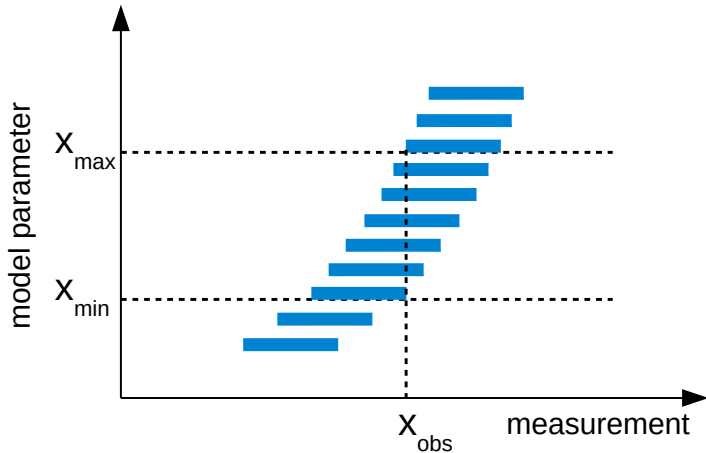
Neyman construction



Confidence interval is $[X_{\min}, X_{\max}]$

Confidence Intervals

Neyman construction



Confidence interval is $[X_{\min}, X_{\max}]$

Application: Uncertainty of Efficiency

(see Exercises No. 3)

- Efficiency: **numerator = subset of denominator**
→ fully correlated, don't use Gaussian error propagation!
- Good description: **binomial uncertainties**
 - Given true efficiency ϵ
 - Draw from a population of N events: expect $\langle n \rangle = \epsilon N$ events on average
 - Variance of n

$$V[n] \equiv \sigma_n^2 = N\epsilon(1 - \epsilon)$$

- But don't know true efficiency: replace with estimator $\hat{\epsilon}$ from measurement (random sample)

$$\epsilon \rightarrow \hat{\epsilon} = \frac{n}{N}, \quad \hat{\sigma}_n^2 = N\hat{\epsilon}(1 - \hat{\epsilon})$$

- Uncertainty on efficiency estimator

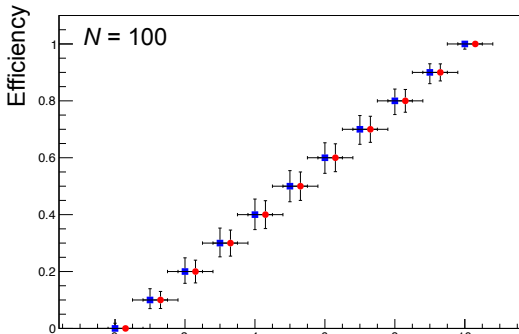
$$\sigma_{\hat{\epsilon}}^2 = \frac{1}{N^2} \hat{\sigma}_n^2 = \frac{\hat{\epsilon}(1 - \hat{\epsilon})}{N}$$

Application: Uncertainty of Efficiency

(see Exercises No. 3)

- Problem with **binomial uncertainty from measured efficiency**: variance $\rightarrow 0$ for $\hat{\epsilon} \rightarrow 0, 1$
- Solution: construction of proper **68% confidence level intervals**

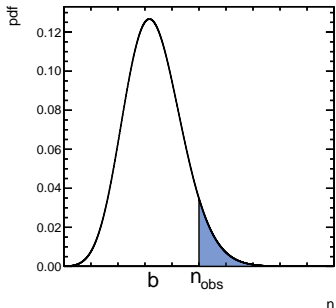
Comparison: Clopper-Pearson – Binomial



3.3.2. Hypothesis testing

- How well is the data described by my hypothesis (model)?
- For example, **how likely that observed peak in data just upward fluctuation** of the background, i. e. no Higgs boson present?
- Assume simple counting experiment
 - b : expected number of background events (=model)
 - n_{obs} : number of observed events
- **“p value”**: **probability of upward fluctuation as large as or larger than observed in data**

$$p \equiv P(n \geq n_{\text{obs}} | b) = \int_{n_{\text{obs}}}^{\infty} dn \mathcal{P}(n|b)$$



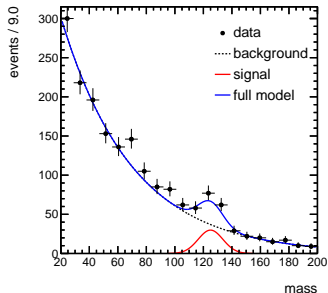
p value is **not the probability of a hypothesis**

p quantifies level of (dis-)agreement between model and data:

→ judgement call whether to keep model or reject it

- Typically problems not as simple as a single measured quantity
 - Multiple channels/measurements, e. g. binned distribution
 - Prediction depends on several parameters, e. g. POI, nuisance params.
- Formally: use appropriate **test statistic** $t : \mathbb{R}^D \rightarrow \mathbb{R}$
 - In general any function that **combines relevant information** from an experiment, e. g. number of observed events per bin, **into one single number** reflecting the agreement between data and hypothesis
 - Likelihood is example for test statistic
- t can be used to compute p value for complex models

$$p = \int_{t_{\text{obs}}}^{\infty} \mathcal{P}(t) dt$$

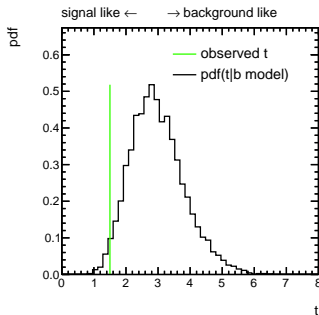


Test Statistic

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$$p = \int_{t_{\text{obs}}}^{\infty} \mathcal{P}(t) dt$$

- Requires pdf $\mathcal{P}(t)$ of the test statistic t : typically from simulation (*toy data*)

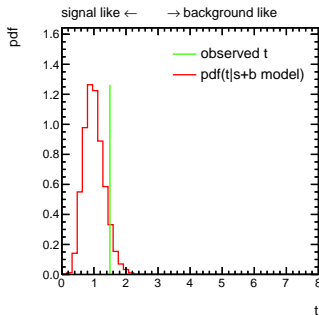


Test Statistic

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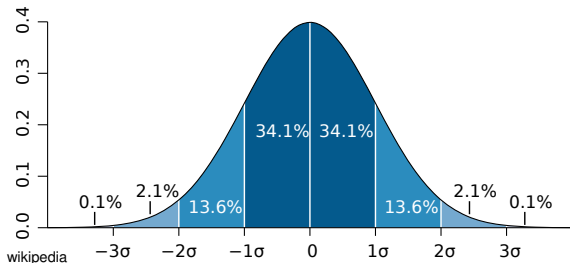


Significance

- Often, p value converted into equivalent **significance Z**:
upward fluctuation from 0 by Z of normal-distributed variable corresponding to same p value
 - Corresponds to Z standard deviations σ of the Gaussian distribution

$$Z = \Phi^{-1}(1 - p)$$

Φ : cumulative (=quantile) function of normal distribution

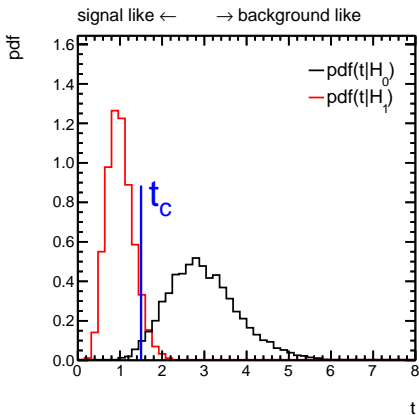


- Convention** to classify effects by significance
 - 3σ : *evidence* for signal (0.3% chance of background fluctuation)
 - 5σ : *discovery* of signal (0.00006% chance of background fluctuation)

- **Cannot determine whether a hypothesis is true** (frequentist)
- But can **define rules how to reject hypothesis** in favour of an alternative hypothesis
- Can determine **probability of wrong choice**: how often wrong choice is made would the experiment be repeated very often

Distinguishing Hypotheses

- Classification problem: how to interpret outcome of an experiment?
 - Want to distinguish between two alternative hypotheses, e. g.
 - H_0 : there are only background events, e. g. no Higgs boson
 - H_1 : there is a signal, e. g. a Higgs boson, contribution to the data
 - Define **test statistic t** , e. g. \mathcal{L}
 - Construct pdf $\mathcal{P}(t|H_i)$ of t under H_0 and H_1
 - In reality often from simulation
 - How compatible is t_{obs} with H_i ?
- Set a critical value t_c and reject H_0 in favour of H_1 if $t_{\text{obs}} < t_c$



Distinguishing Hypotheses

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 - H_0 : there are only background events, e. g. no Higgs boson
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- Types of errors

I: reject H_0 although it is true

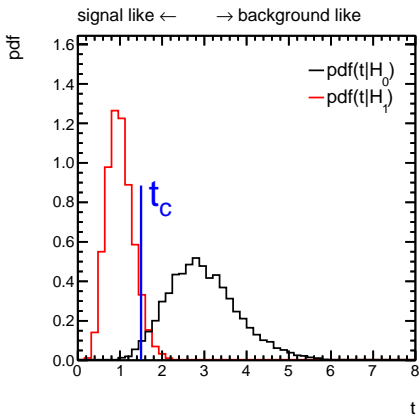
II: accept H_0 although H_1 is true

$$\text{I: } P(t < t_c | H_0) = \int_{-\infty}^{t_c} dt \mathcal{P}(t | H_0) \equiv \alpha$$

$$\text{II: } P(t > t_c | H_1) = \int_{t_c}^{\infty} dt \mathcal{P}(t | H_1) \equiv \beta$$

α : significance of test

$1 - \beta$: power of test



Distinguishing Hypotheses

- Classification problem: how to interpret outcome of an experiment?
- Want to distinguish between two alternative hypotheses, e. g.
 - H_0 : there are only background events, e. g. no Higgs boson
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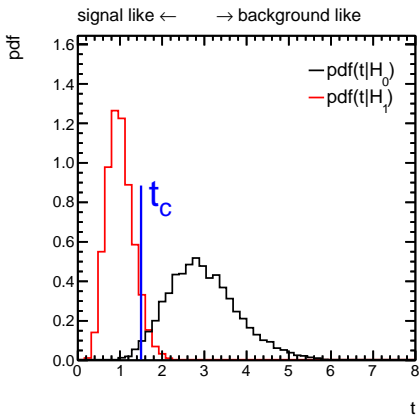
I: reject H_0 although it is true

II: accept H_0 although H_1 is true

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$$\text{II: } P(t > t_c | H_1) = \int_{t_c}^{\infty} dt \mathcal{P}(t | H_1) \equiv \beta$$

By choice of t_c (choice of α): association to one or the other hypothesis performed
“at the confidence level α ”



Difference Between α and p Value?

- In one sense, no difference
 - Both are $\int_x^\infty dx' \mathcal{P}(x')$
- But **conceptually very different**
 - α is computed before one sees the data: predefined property of the test
 - p depends on the data (property of the data) and is a random variable

Neyman-Pearson Lemma

- When performing a test between two hypotheses (models) H_0 and H_1 , the **likelihood ratio test**, which rejects H_0 in favour of H_1 if

$$Q = \frac{\mathcal{L}_{H_1}}{\mathcal{L}_{H_0}} > Q_c \quad \text{with } P(Q > Q_c | H_0) = \alpha$$

is the **most powerful test** at a significance level α

- The test statistic Q is called **likelihood ratio**
- For a number of reasons, usually

$$q = -2 \ln Q$$

- Used test statistic in particle-physics experiments evolved with time
- Hypothesis test example
 - H_0 : background-only hypothesis, e. g. SM without Higgs boson
 - H_1 : Higgs boson with fixed mass and signal strength μ present in data
- At the LHC commonly: profile likelihood ratio test statistics

$$q_\mu = -2 \ln \frac{\mathcal{L}(n_{\text{obs}}; \mu, \hat{\theta}(\mu))}{\mathcal{L}(n_{\text{obs}}; \hat{\mu}, \hat{\theta})}$$

- Nuisance parameters are profiled in nominator
- Global maximum of \mathcal{L} under (μ, θ) as denominator with $0 \leq \hat{\mu} \leq \mu$
- Allows usage of certain approximations when computing $\mathcal{P}(q_\mu | H_i)$ (“asymptotic formulae” based on Wilks and Wald theorem¹)

¹G. Cowan “Asymptotic formulae for likelihood-based tests of new physics”, Eur. Phys. J. C71:1554 (2011)

3.3.3. Search for new physics (exclusion limits)

- Assume measurement with a given sensitivity: no signal observed
- How much signal can “hide” in the bkg. fluctuations (+uncertainty)?
- **How large could a signal be at most?**

Excluding Parameters

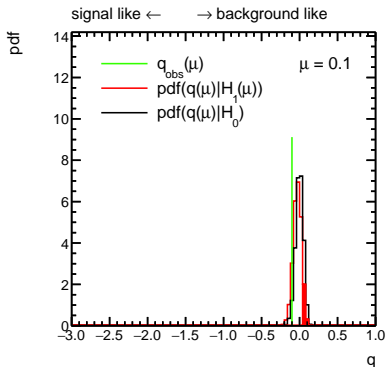
- No signal observed: **how large could a signal be at most?**
- Formally: test statistic q , depending on signal-strength modifier μ

$$q(\mu) = -2 \ln \frac{\mathcal{L}_{H_1}(\mu)}{\mathcal{L}_{H_0}}$$

What is largest signal $\mu \equiv \mu_{1-\alpha}$
for which H_1 would be rejected at
significance level α ?

$$\alpha = \int_{q_{\text{obs}}}^{\infty} dq \mathcal{P}(q(\mu_{1-\alpha}) | H_1) \equiv \text{CL}_{\text{s+b}}$$

i. e. for $\mu_{1-\alpha}$: $q_{\text{obs}} = q_c$



Excluding Parameters

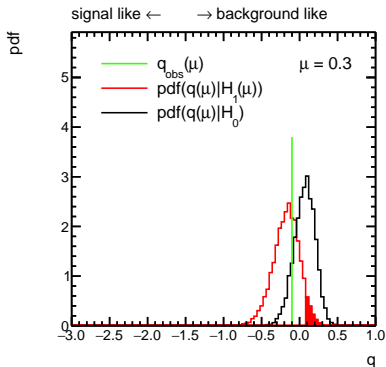
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Excluding Parameters

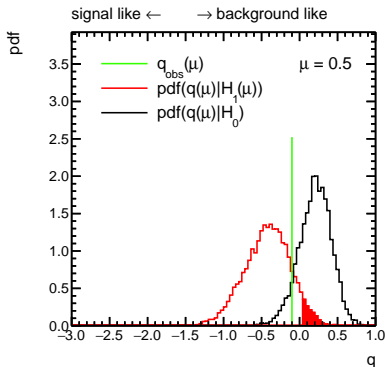
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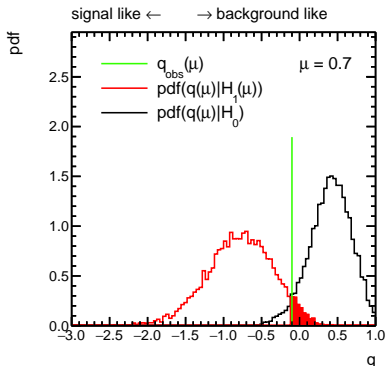
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i. e. for $\mu_{1-\alpha}$: $q_{\text{obs}} = q_c$



Excluding Parameters

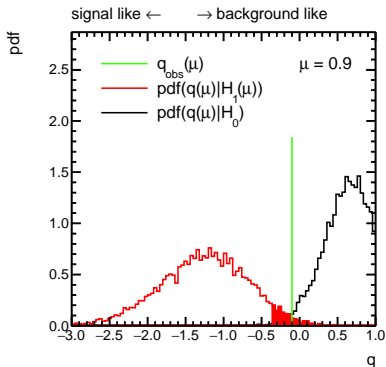
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Excluding Parameters

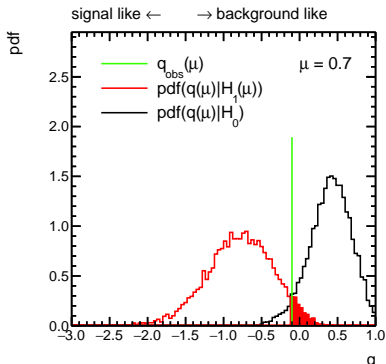
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i. e. for $\mu_{1-\alpha}$: $q_{\text{obs}} = q_c$



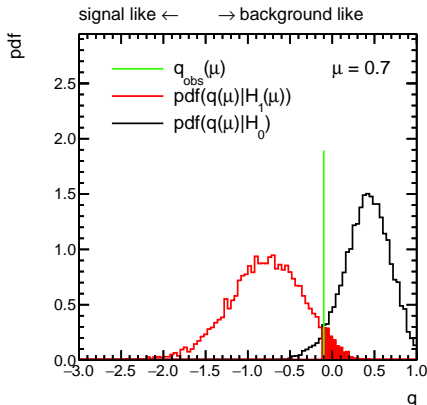
$\mu_{1-\alpha}$ is called **upper limit** on μ at confidence level C.L. = $1 - \alpha$
In particle physics usually 95 % C.L. limit

(Observed) Upper Limit: Interpretation

- “Maximal signal that we would still reject”
- 95 % C.L. upper limit on μ : largest value of μ that would still be rejected in a test with significance 5% given the data
 - NB: limit is a **function of the data** (depends on q_{obs})!
- μ for which $\text{CL}_{s+b} = 0.05$:

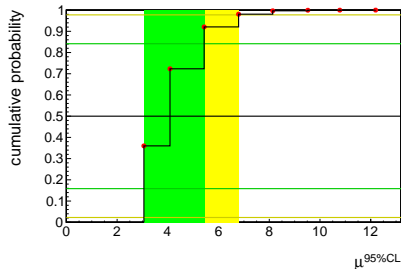
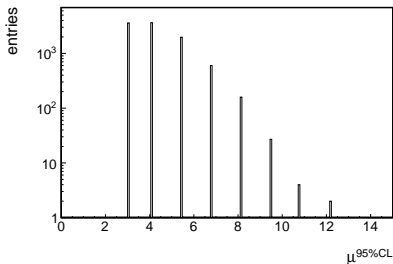
$$0.05 = \int_{q_{\text{obs}}}^{\infty} dq \mathcal{P}(q(\mu_{95}) | H_1)$$

- Upper **limit covers true value** ($\mu_{\text{true}} < \mu_{95}$) **with probability C.L. = 95 %**
 - If the experiment is repeated many times, μ_{95} would be larger than μ_{true} in 95 % of the cases
- Still **5 % chance of wrong exclusion**, i. e. that $\mu_{\text{true}} > \mu_{95}$



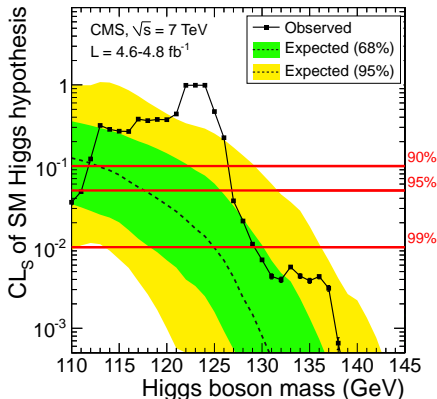
Expected Limit

- **Estimate what observed limit would look like in case of no signal**
- Obtained e. g. from *toy dataset*
 - Sample toy data for q under background-only hypothesis from $\mathcal{P}(q|H_0)$
 - Treat each as observation and compute μ_{95} limit
 - Obtain quantiles from distribution of all μ_{95}
- **Expected limit = median of μ_{95} distribution**
 - 16 and 84% quantiles: **68% confidence interval**
 - 2.5 and 97.5% quantiles: **95% confidence interval**
→ “Brazilian band” plots



Before the Higgs-Boson Discovery

- Combination of Higgs-boson search results by CMS [Phys.Lett. B710 (2012) 26]

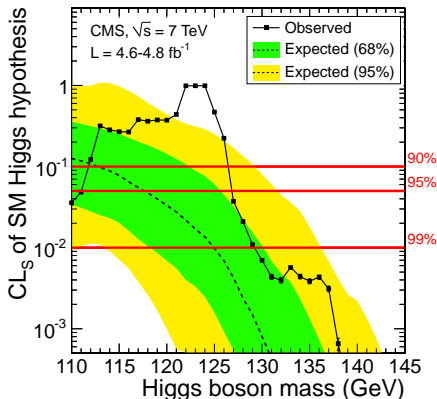


- Tested hypotheses
 - H_0 : no Higgs boson
 - H_1 : SM Higgs boson
- Test statistic q evaluated for SM Higgs boson of different mass ($\mu = 1$ in each case)

- Excluding a SM Higgs boson at 95% CL with masses
 - $m_H > 118$ GeV expected (from toy data under H_0)
 - $m_H > 127$ GeV observed (from real data)

Before the Higgs-Boson Discovery

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- Tested hypotheses
 - H_0 : no Higgs boson
 - H_1 : SM Higgs boson
- Test statistic q evaluated for SM Higgs boson of different mass ($\mu = 1$ in each case)

- Observed exclusion weaker than expected (smaller mass range)
- Around $m_H = 125$ GeV, q_{obs} differs significantly from expectation:
indication that H_0 is wrong!