

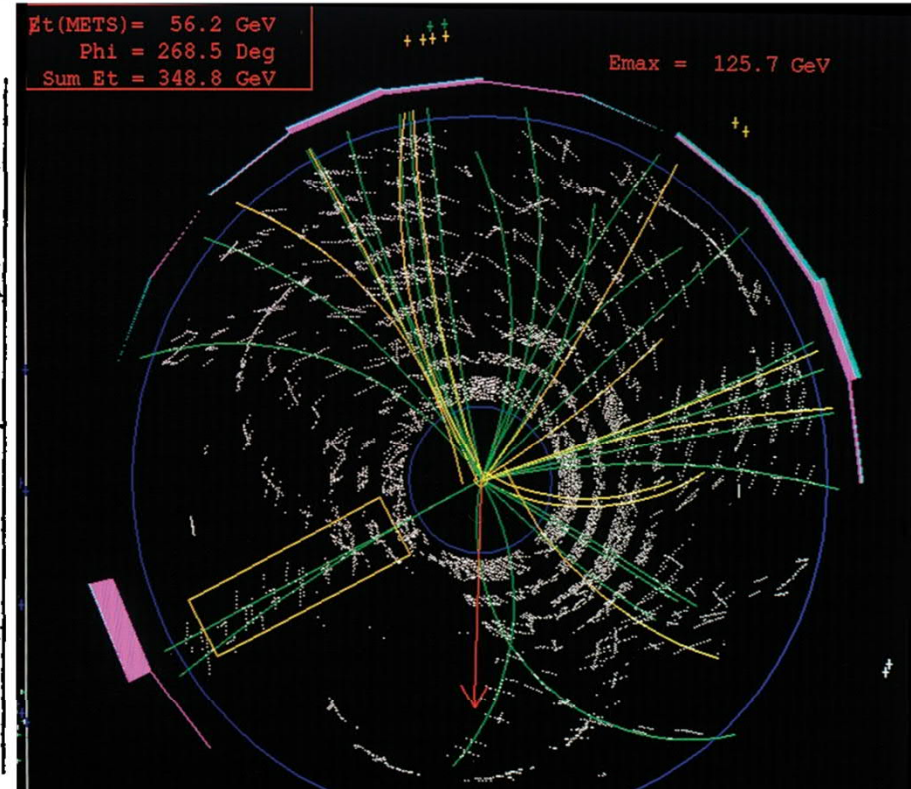
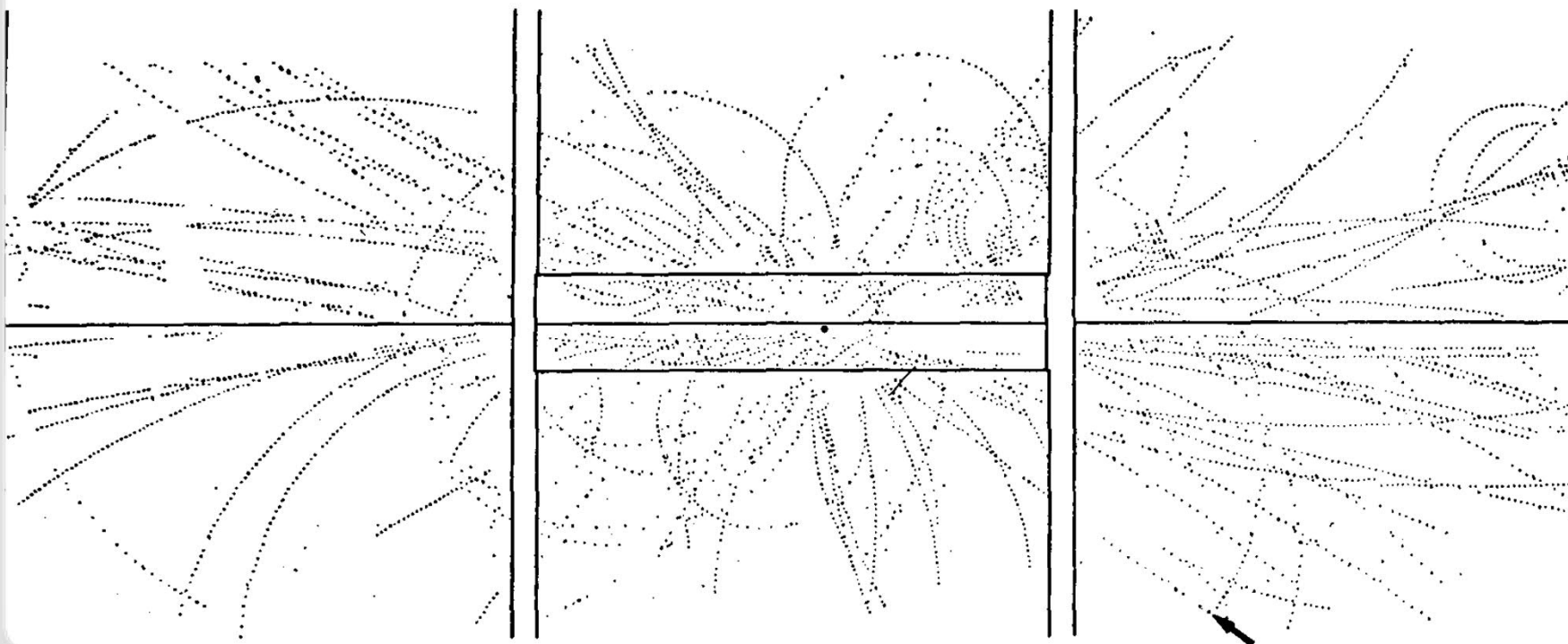
KSETA-Course: Accelerator-Based Particle Physics

Electroweak Physics

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EVENT 2958. 1279.



Elektromagnetic Interaction & CC

- Electromagnetic Interaction (Fermion with charge e)

$$\mathcal{L}_{em} = \underbrace{e}_{\text{elem. charge}} \cdot \underbrace{j_{\mu}^{em}}_{\text{EM current}} \cdot \underbrace{A^{\mu}}_{\text{photon}} = \frac{gg'}{\sqrt{g^2 + g'^2}} \bar{f} \gamma_{\mu} f A^{\mu}$$



with $e = g \sin\vartheta_w = g' \cos\vartheta_w$ weak mixing angle

- Charged Current: V-A structure

$$\mathcal{L}_{cc} = \frac{g}{\sqrt{2}} [J_{\mu}^{+CC} W^{\mu-} + J_{\mu}^{-CC} W^{\mu+}]$$

CC for quarks

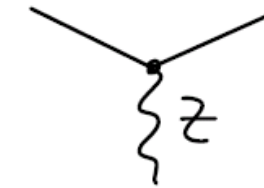
$$J_{\mu}^{+CC} = (\bar{u}, \bar{c}, \bar{t}) \underbrace{\gamma_{\mu} \frac{1}{2} (1 - \gamma_5)}_{\text{V-A}} V_{CKM} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

quark mixing



NC & Selfcoupling

- Neutral Currents $\mathcal{L}_{NC} = -\frac{\sqrt{g^2 + g'^2}}{2} J_\mu^{NC} Z^\mu$



with:

$$J_\mu^{NC} = \bar{f} \gamma_\mu \frac{1}{2} (g_V^f - g_A^f \gamma_5) f$$

vector-coupling

axial-coupling

$$g_V^f = (I_3^f - 2q^f \sin^2 \vartheta_W) \quad g_A^f = I_3^f$$

(I_3^f third component of isospin, q^f fermion charge)

- Selfcouplings of gauge bosons: only $WW\gamma$, WWZ

$$\mathcal{L}_{WW\gamma} = -ie [A_\mu (W^{-\mu\nu} W_\nu^+ - W^{+\mu\nu} W_\nu^-) + F_{\mu\nu} W^{+\mu} W^{-\nu}]$$

$$\mathcal{L}_{WWZ} = -ie \cot \vartheta_W [Z_\mu (W^{-\mu\nu} W_\nu^+ - W^{+\mu\nu} W_\nu^-) + Z_{\mu\nu} W^{+\mu} W^{-\nu}]$$

- +quartic couplings $WWWW$, $WWZZ$, $WWZ\gamma$, $WW\gamma\gamma$

Cross Section

- Resonant (s-channel-) production von Z-bosons in e^+e^- -scattering
 - Photon and Z-boson: identical quantum numbers ($J^P = 1^-$) \rightarrow interference
 - Matrix-element:

$$|M|^2 = \left| \begin{array}{c} e^- \quad \quad \quad \bar{f} \\ \searrow \quad \quad \nearrow \\ \quad \quad \gamma^* \quad \quad \quad \\ \nearrow \quad \quad \searrow \\ e^+ \quad \quad \quad f \end{array} + \begin{array}{c} e^- \quad \quad \quad \bar{f} \\ \searrow \quad \quad \nearrow \\ \quad \quad Z \quad \quad \quad \\ \nearrow \quad \quad \searrow \\ e^+ \quad \quad \quad f \end{array} \right|^2$$

- cross section: $\sigma(e^+e^- \rightarrow f\bar{f}) = \sigma_{\gamma^*} + \sigma_{\gamma^*/Z} + \sigma_Z$
- $\sqrt{s} \ll m_Z$: photon exchange dominant \rightarrow simple QED
- $\sqrt{s} \simeq m_Z$: Z-boson-exchange dominant, photon- and interference-term negligible

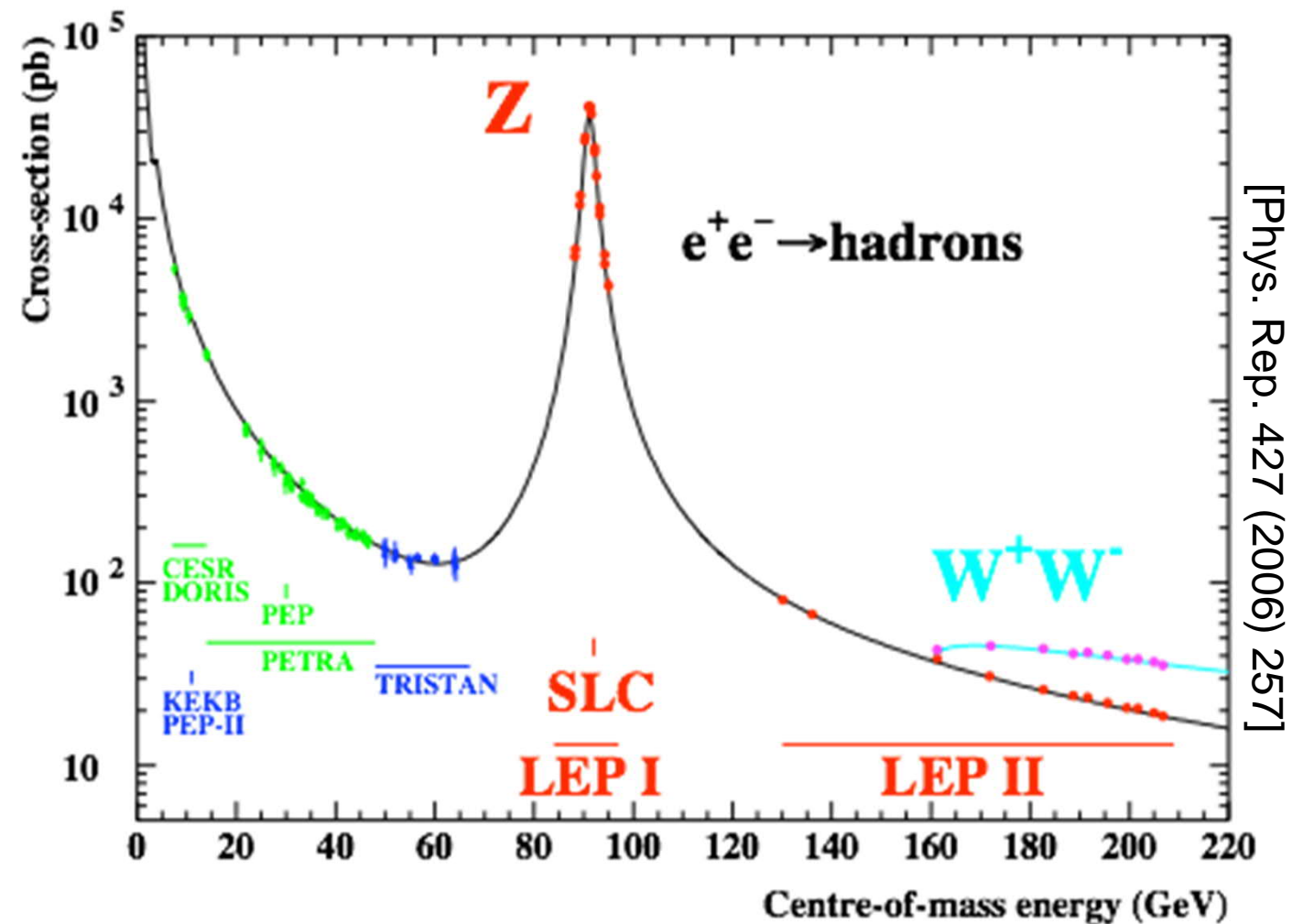
Width of the Z-Resonance

- Z = unstable particle
 - finite lifetime τ_Z
 - decaywidth $\Gamma_Z = 1/\tau_Z$ („smeared” mass)
 - modified propagator
- Decay width
 - sum of partial decay widths („partial widths”)
 - for $\sqrt{s} \simeq m_Z$:

$$\Gamma_Z = \sum_f \Gamma_f = \sum_{q=u,d,s,c,b} \Gamma_q + \sum_{\ell=e,\mu,\tau} \Gamma_\ell + \sum_{\nu=\nu_e,\nu_\mu,\nu_\tau} \Gamma_\nu$$

- in leading order:

$$\Gamma_f = N_C^f \frac{G_F m_Z^3}{6\sqrt{2}\pi} \left[(g_V^f)^2 + (g_A^f)^2 \right]$$



Number of light neutrinos

- Cross section for a fermion f:

$$\sigma_f = \underbrace{\frac{12\pi}{m_Z^2} \cdot \frac{\Gamma_e \cdot \Gamma_f}{\Gamma_Z^2}}_{\sigma_f^0} \cdot \frac{s \cdot \Gamma_Z^2}{(s - m_Z^2)^2 + s^2 \frac{\Gamma_Z^2}{m_Z^2}}$$

Breit-Wigner
 $\Rightarrow 1$ for $s = m_Z^2$

- Compare different cross sections to find Γ_{inv} from Γ_Z and other f

$$\frac{\Gamma_{inv}}{\Gamma_e} \equiv R_{inv}^0 = \frac{\Gamma_Z}{\Gamma_e} - 3 - R_e^0 = \sqrt{\frac{12\pi}{m_Z^2} \cdot \frac{R_e^0}{\sigma_{had}^0}} - 3 - R_e^0$$

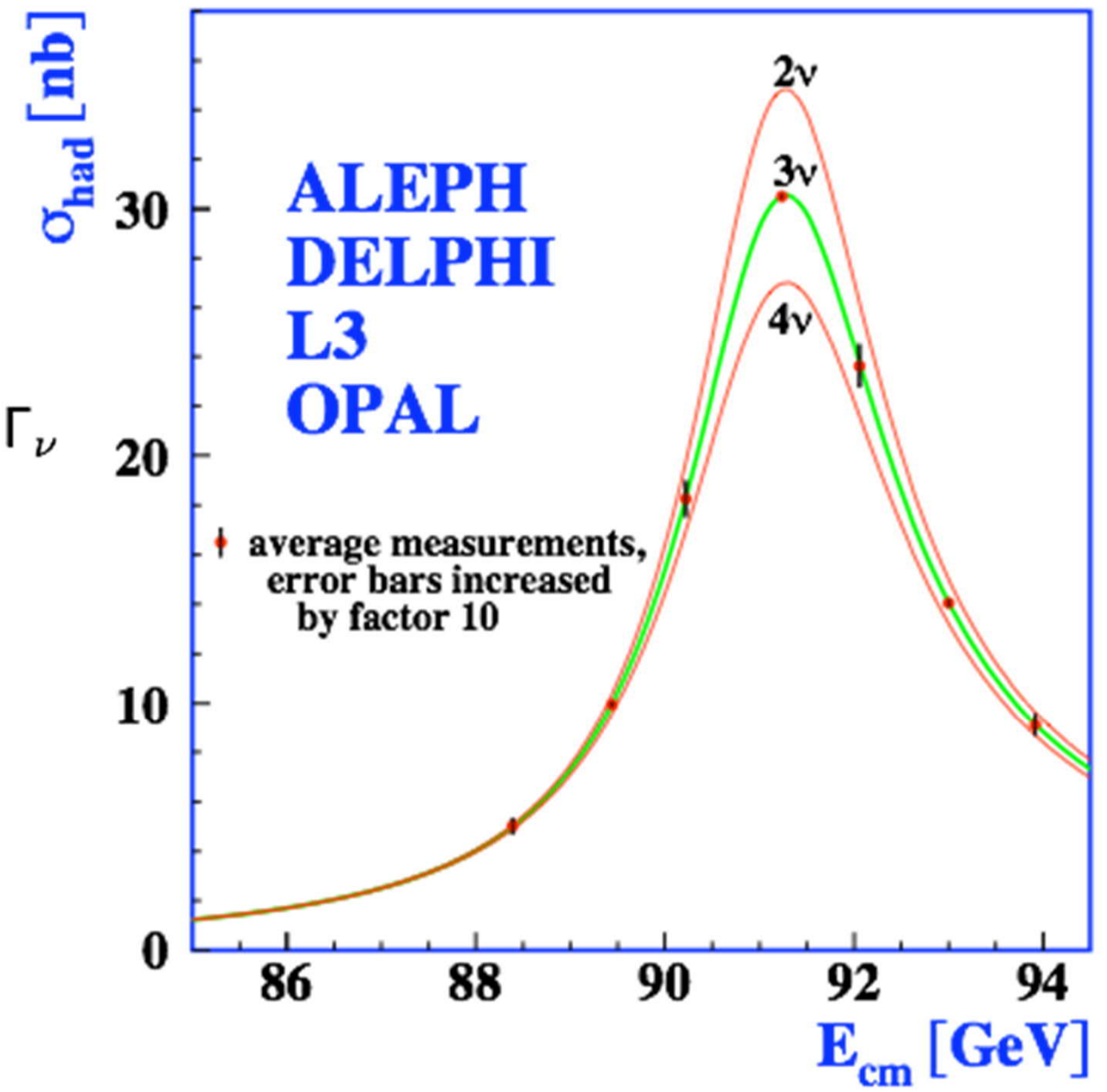
- Divide Γ_{inv} by Γ_ν derived from theory

$$\Rightarrow N_\nu = 2.9840(82)$$

Number of neutrinos in pictures

$$\sigma_f = \frac{12\pi}{m_Z^2} \cdot \frac{\Gamma_e \cdot \Gamma_f}{\Gamma_Z^2} \cdot \frac{s \cdot \Gamma_Z^2}{(s - m_Z^2)^2 + s^2 \frac{\Gamma_Z^2}{m_Z^2}}$$

$$\Gamma_Z = \sum_f \Gamma_f = \sum_{q=u,d,s,c,b} \Gamma_q + \sum_{\ell=e,\mu,\tau} \Gamma_\ell + \sum_{\nu=\nu_e,\nu_\mu,\nu_\tau} \Gamma_\nu$$

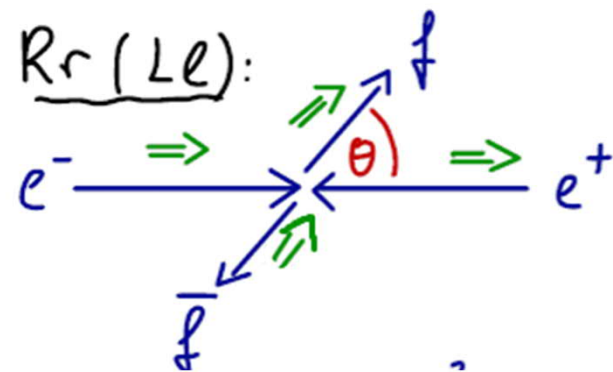


[Phys. Rep. 427 (2006) 257]

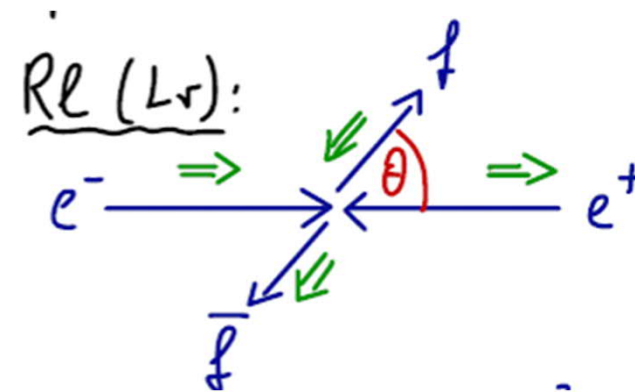
$\sqrt{s} \ll m_Z$: differential Xsec

$\sqrt{s} \ll m_Z$: photon exchange (pure QED)

helicities: photon \Rightarrow spin 1



$$\sigma \sim (1 + \cos \vartheta)^2$$



$$\sigma \sim (1 - \cos \vartheta)^2$$

$$\Rightarrow \frac{d\sigma_\gamma}{d \cos \vartheta} = N_c^f \cdot q_f^2 \cdot \frac{\pi \alpha^2}{2s} (1 + \cos^2 \vartheta)$$

colors

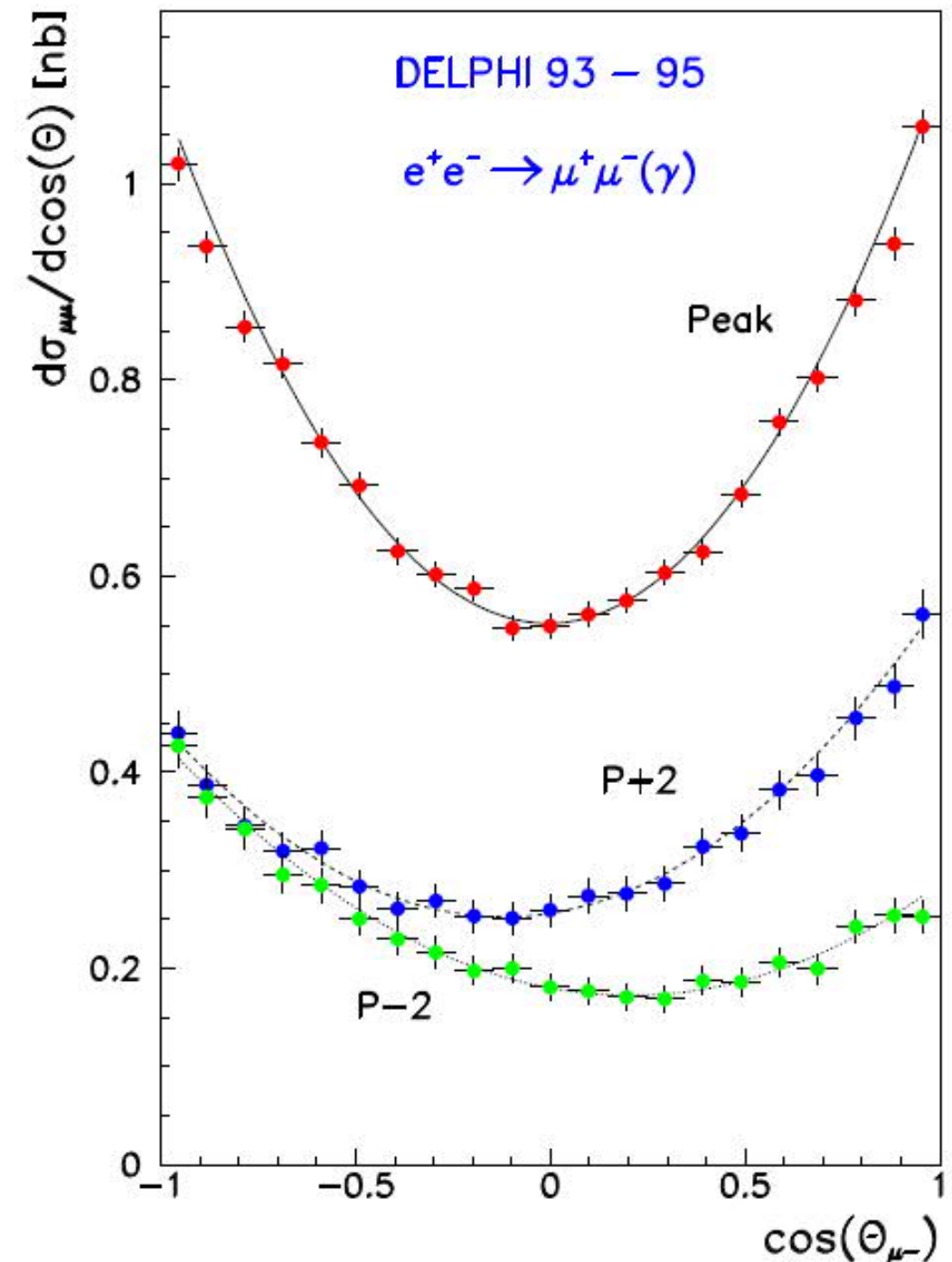
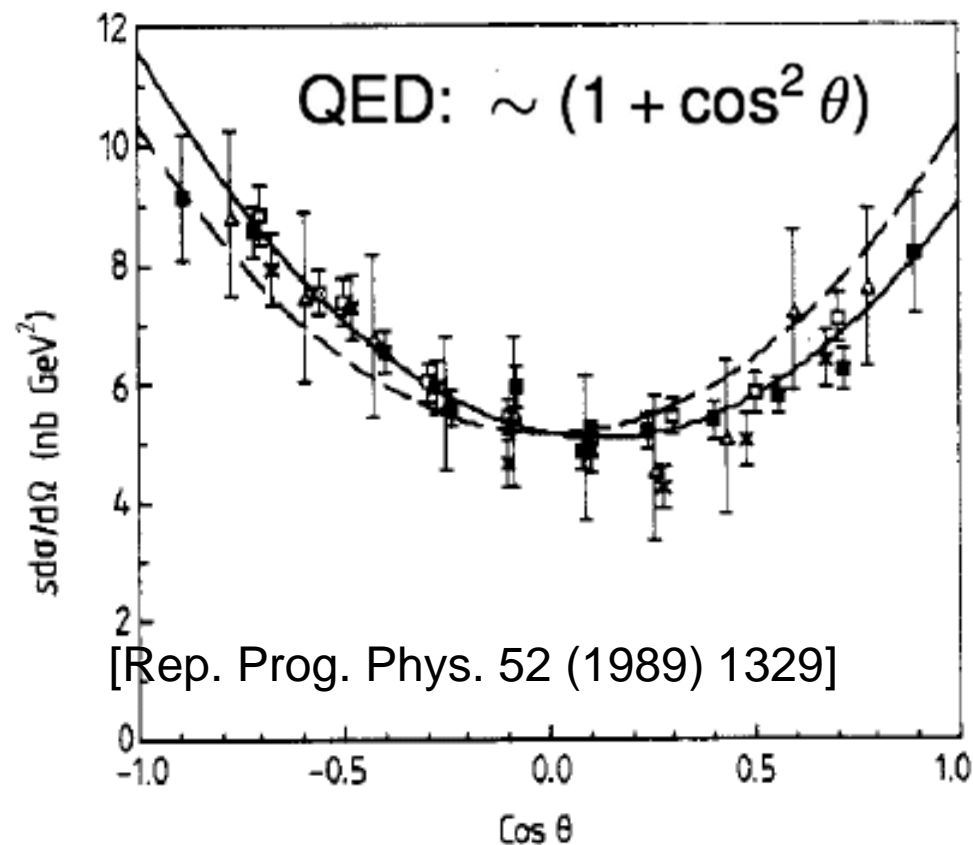
charge

$\sqrt{s} \ll m_Z$: Angular Distribution

LEP: $\sqrt{s} = m_Z \pm 2 \text{ GeV}$

- Before LEP: measurements at PETRA (DESY) \rightarrow first deviations from pure QED
- LEP: γ^*/Z -interference und Z central physics topic

PETRA: $\sqrt{s} = 29\text{-}35.5 \text{ GeV}$



[Phys. Rep. 427 (2006) 257]

Asymmetries

- Generic definition of an **asymmetry**:
Partition a dataset into two parts $X, Y \rightarrow A = \frac{X - Y}{X + Y}$
- Why asymmetries?
 - Asymmetries = **Ratios**, not absolute rates
 - Backgrounds and systematic effects on numerator/denominator equal or similar
→ **Reduction of uncertainties** due to cancellations
 - Increased sensitivity to **small differences**

Differential Cross Section

Angular distribution for Z exchange:

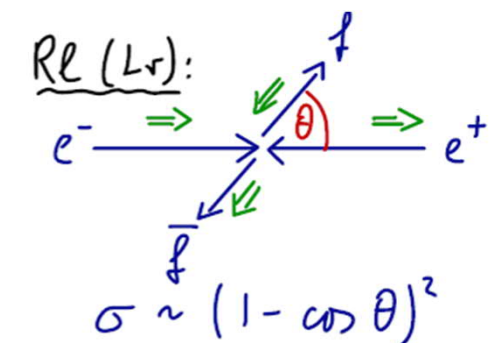
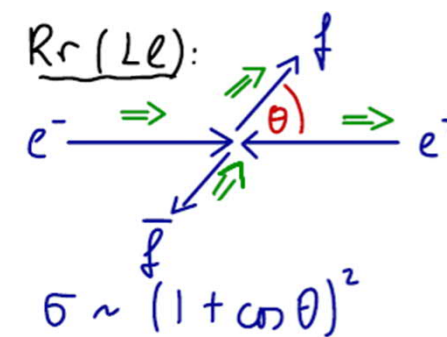
→ 4 helicities (using $g_R = g_V + g_A, g_L = g_V - g_A$)

$$\frac{d\sigma_f}{d\cos\theta} \sim (g_L^e)^2 (g_L^f) (1 + \cos\vartheta)^2$$

$$\frac{d\sigma_f}{d\cos\theta} \sim (g_L^e)^2 (g_R^f) (1 - \cos\vartheta)^2$$

$$\frac{d\sigma_f}{d\cos\theta} \sim (g_R^e)^2 (g_R^f) (1 + \cos\vartheta)^2$$

$$\frac{d\sigma_f}{d\cos\theta} \sim (g_R^e)^2 (g_L^f) (1 - \cos\vartheta)^2$$



$$\rightarrow \frac{d\sigma_f}{d\cos\theta} = \frac{3}{8} \sigma_f [(1 + \cos^2\vartheta) + 2 A_e A_f \cos\vartheta]$$

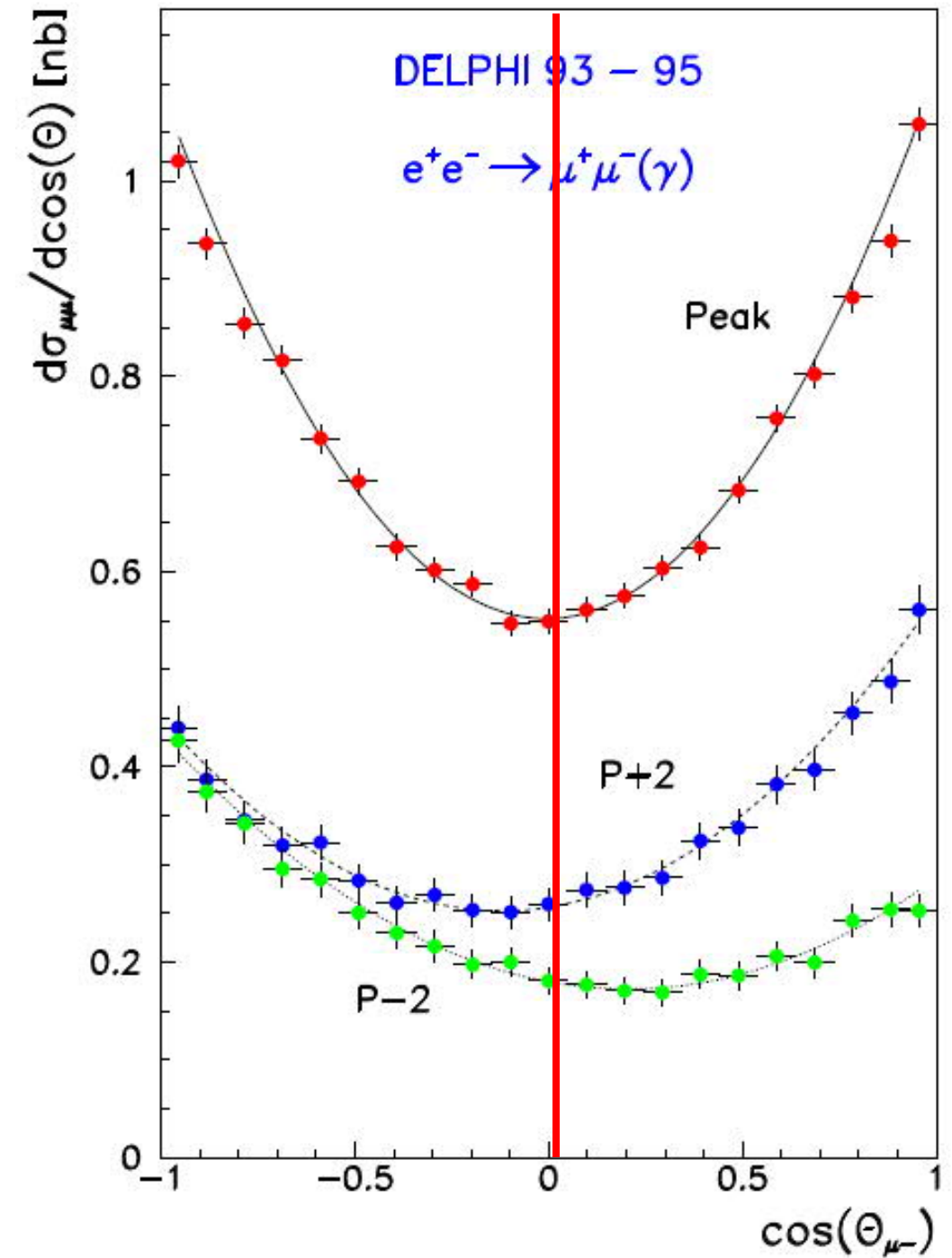
$$A_f = \frac{(g_L^f)^2 - (g_R^f)^2}{(g_L^f)^2 + (g_R^f)^2} = 2 \frac{g_V^f / g_A^f}{1 + \left(\frac{g_V^f}{g_A^f}\right)^2}$$

access ratio $\left(\frac{g_V^f}{g_A^f}\right)$

Forward-Backward-Asymmetry

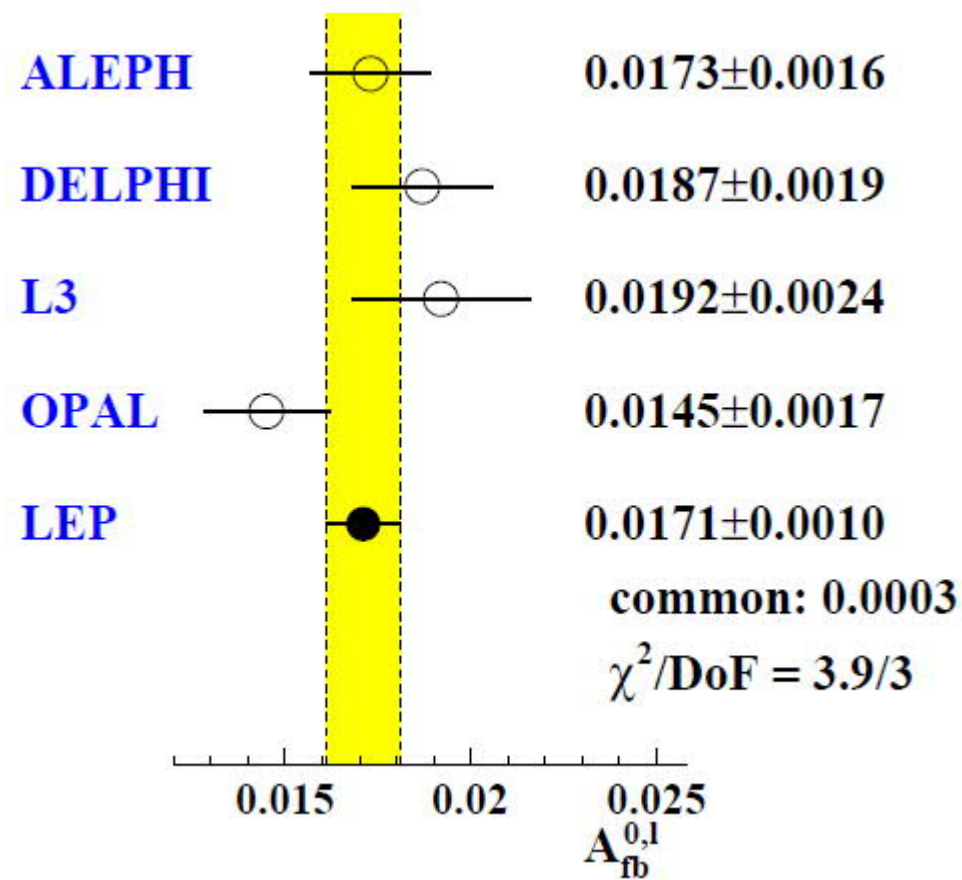
$$\sigma_F \equiv \int_0^{\pi/2} \frac{d\sigma}{d \cos \vartheta} d\vartheta \quad \sigma_B \equiv \int_{\pi/2}^{\pi} \frac{d\sigma}{d \cos \vartheta} d\vartheta$$

$$A_{FB} \equiv \frac{\sigma_F - \sigma_B}{\sigma_F + \sigma_B} = \dots = \frac{3}{4} A_e A_f$$

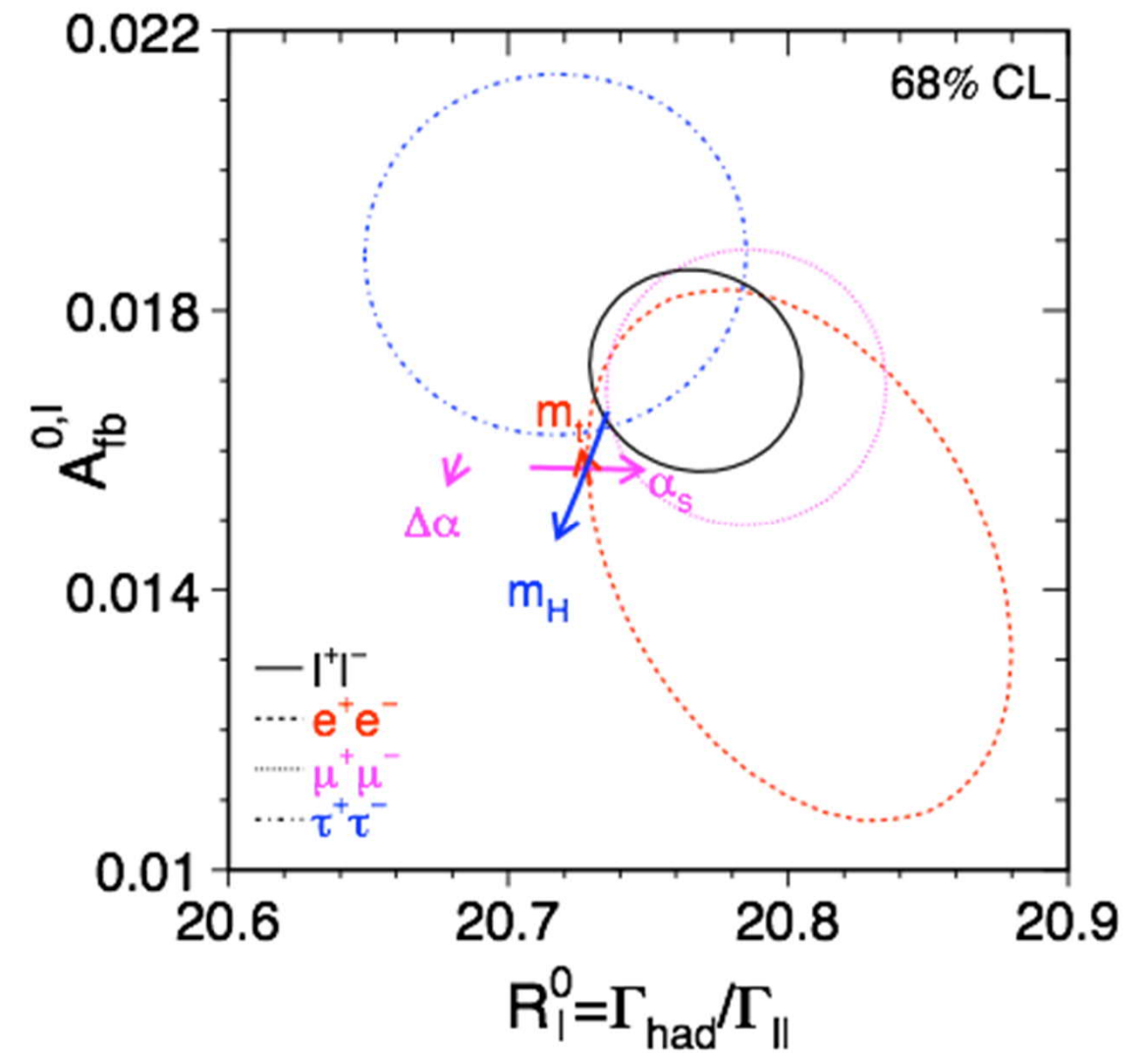


A_{FB} : Results from LEP

LEP-average: A_{FB} for leptons

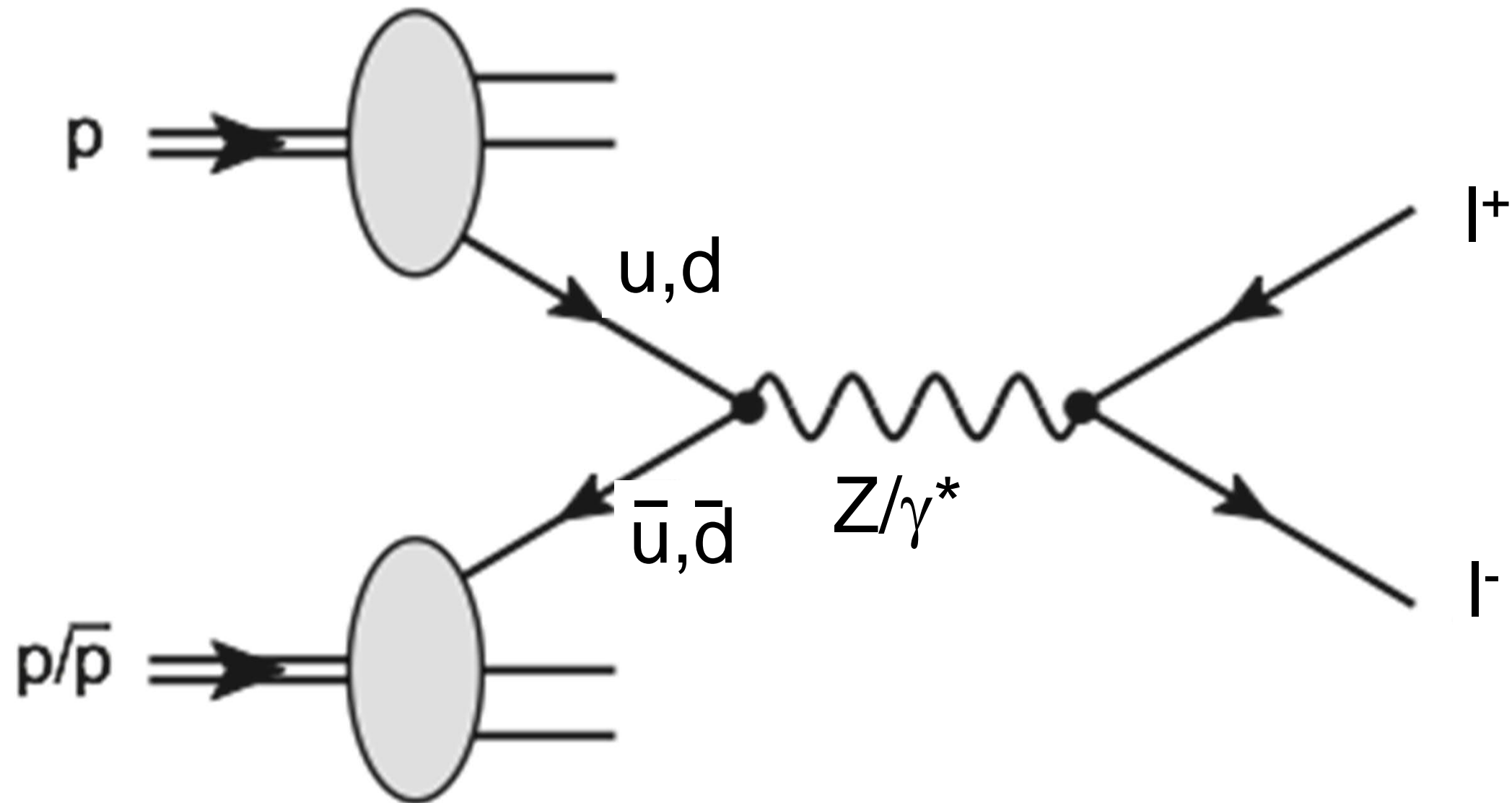


A_{FB} separately for e, μ , τ vs. R^0



[Phys. Rep. 427 (2006) 257]

Z at hadron colliders



- Lepton final states preferred for low backgrounds, precise reco
- Learn about proton structure (See other lecture)
- Unknown initial state complicates studies of EWK physics

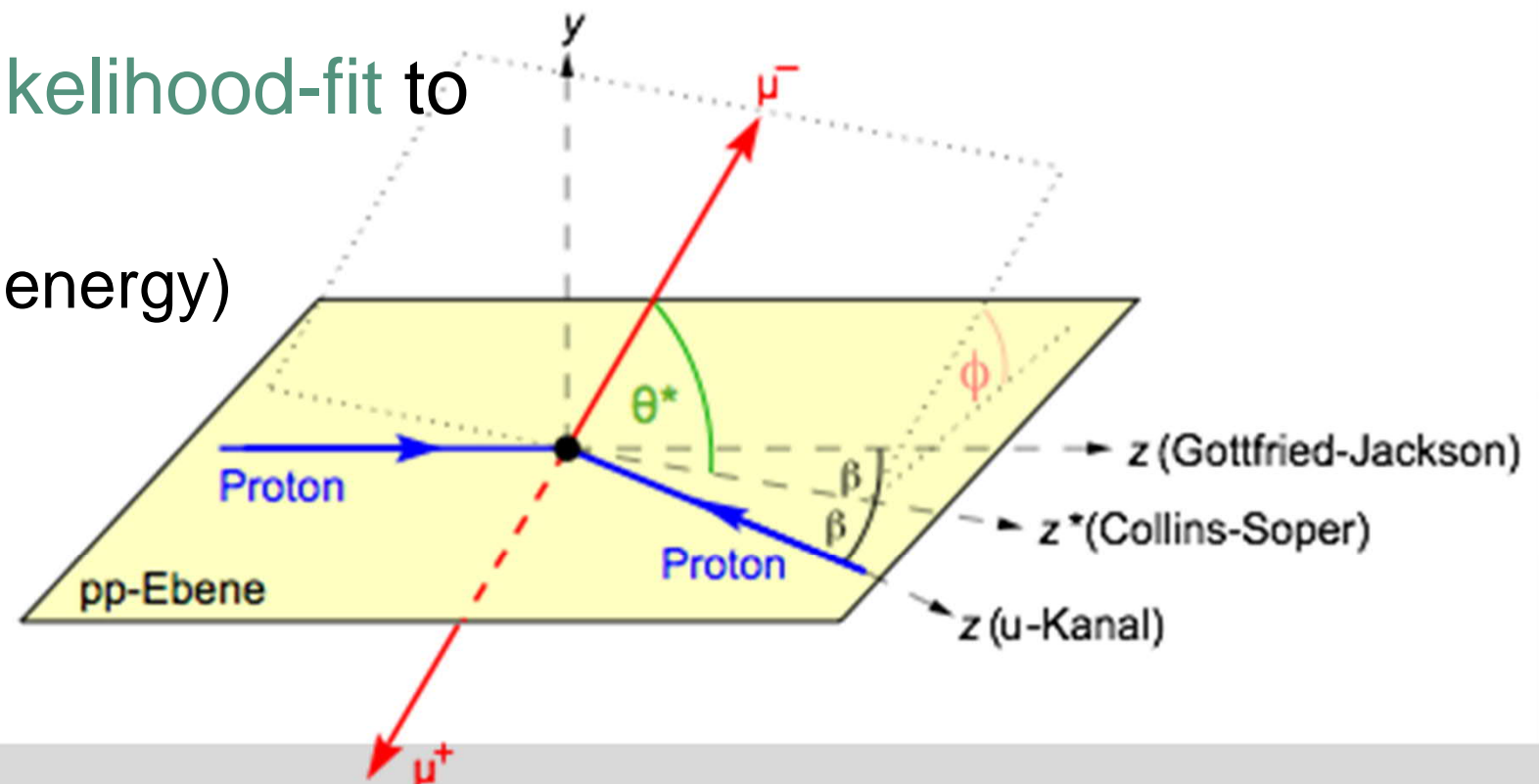
Weak mixing angle

- Drell-Yan-Prozess: $qq \rightarrow \gamma^*/Z \rightarrow ff$
 - Standard model: relative couplings of fermions to γ/Z given by $\sin^2 \theta_{W,\text{eff}}^f$

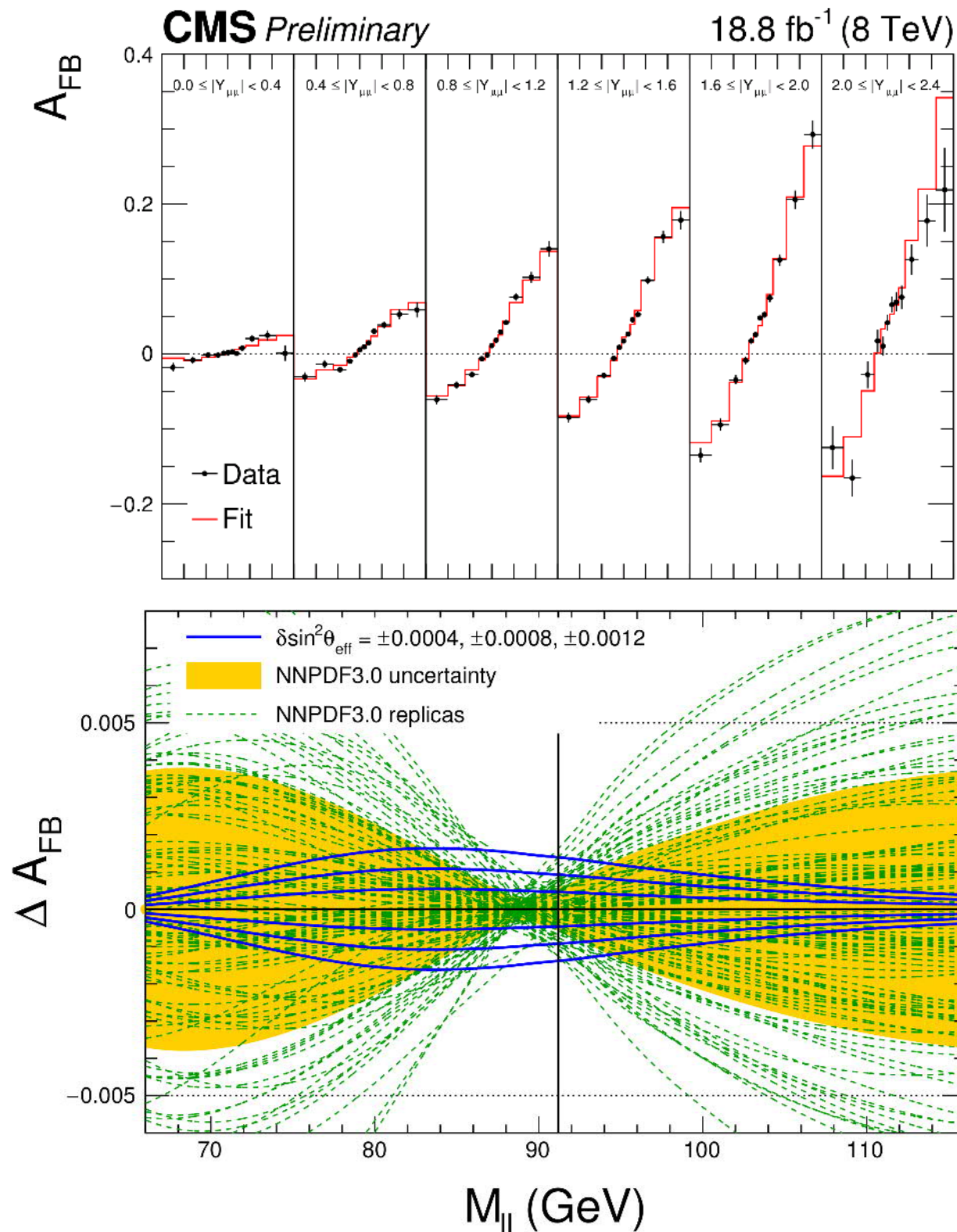
$$\sin^2 \theta_{W,\text{eff}}^f = \frac{I_3^f}{2q^f} \left(1 - \frac{g_V^f}{g_A^f} \right)$$
 - Assumption: differential cross section and PDFs known
 \rightarrow Extraktion von $\sin^2 \theta_{W,\text{eff}}^f$
 - Tricky: which direction did the quark/antiquark come from?

- Simultaneous maximum-likelihood-fit to

- Lepton pair mass $M(\mu\mu)$
 (= partonic center of mass energy)
- A_{FB} (Collins-Soper-frame)



Weak Mixing Angle



CMS-results:

$$\sin^2 \theta_{W,\text{eff}}^f = 0.23101 \pm 0.00052$$

- consistent with LEP- resultat ($\sin^2 \theta_{W,\text{eff}}^f = 0.23153 \pm 0.00016$)
- uncertainty: 0.2% (LEP: 0.07%)

Systematic uncertainties:

=> reduce effect of PDF uncertainty by simultaneous fit

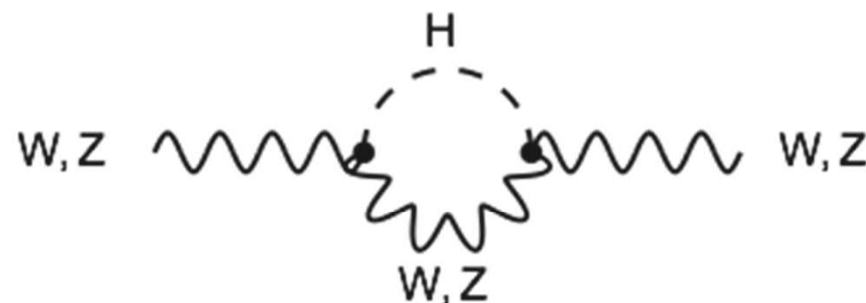
[CMS-SMP-16-007]

Measuring the W-mass

- electroweak theory predicts
 - Connection of W- and Z-mass by the weak mixing angle

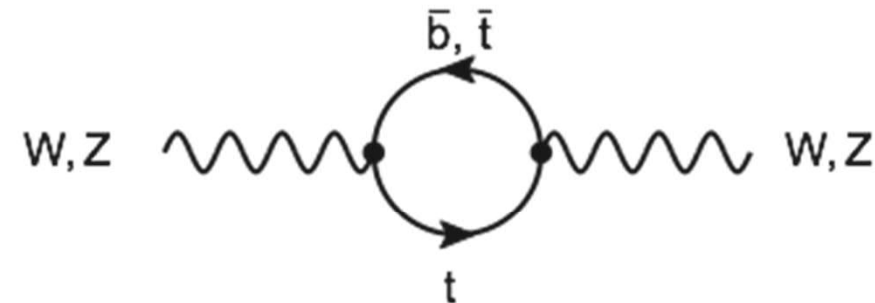
$$m_W^2 = \frac{g^2 v^2}{4}, \quad m_Z^2 = \frac{v^2}{4} (g^2 + g'^2) \quad \rightarrow \quad \rho_0 = \frac{m_W^2}{m_Z^2 \cos^2 \theta_W} = 1$$

- Connection with Top-quarks and Higgs-Boson masses through loop diagram, i.e. „self-energy” of W and Z



$$\sim g' \ln(m_H / m_W)$$

weak dependence
(logarithmic)



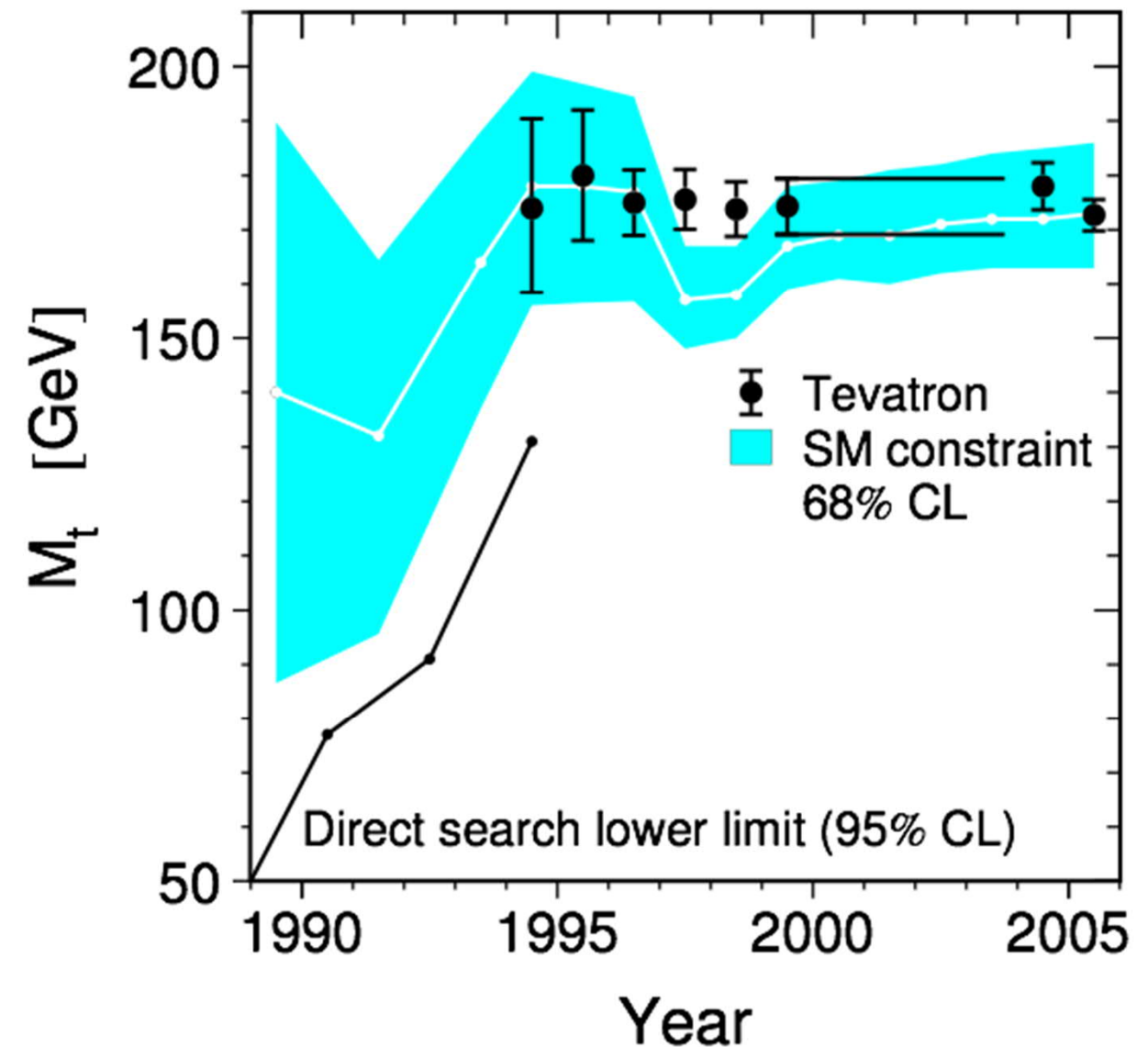
$$\sim G_F m_t^2$$

strong dependence
(quadratic)

Measuring the W-mass

- Looking for indirect effect needs **highest precision**
- Z-mass: **extremely precise** measurements at LEP (uncertainty: $2 \cdot 10^{-5}$)
- W-mass: LEP + Tevatron
- **Prediction** of Top-quark mass before discovery
- bounds on **allowed Higgs-boson masses** before discovery

vgl. Tevatron (2012):
 $m_t = 173.2 \pm 0.9 \text{ GeV}$

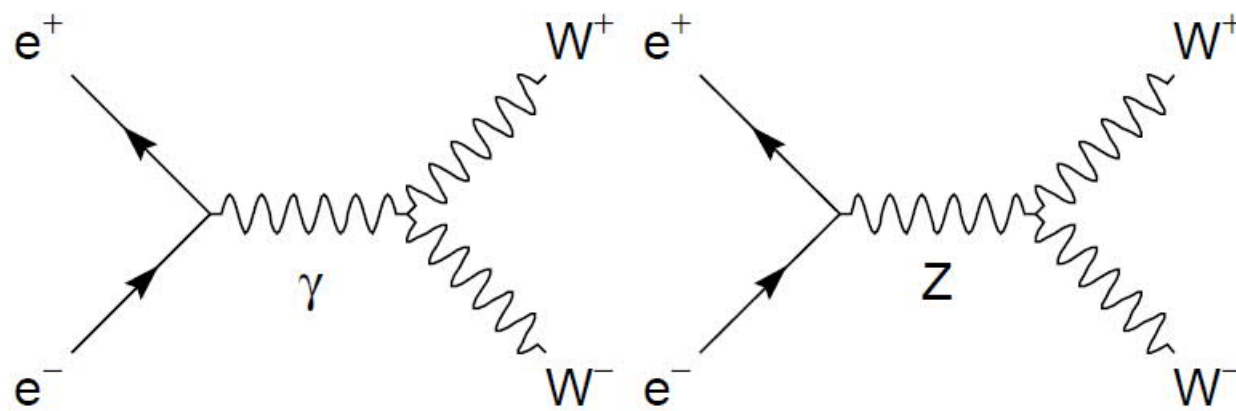


[Phys. Rept. 427 (2006) 257]

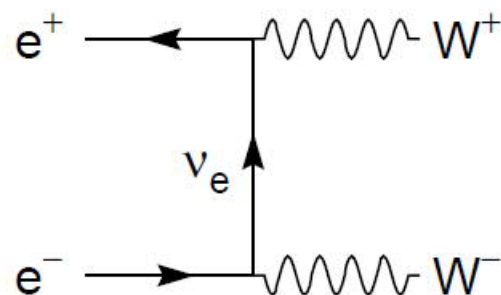
W-pair creation at LEP-II

- LEP-II: passes **kinematic Threshold** for W-boson-pair creation

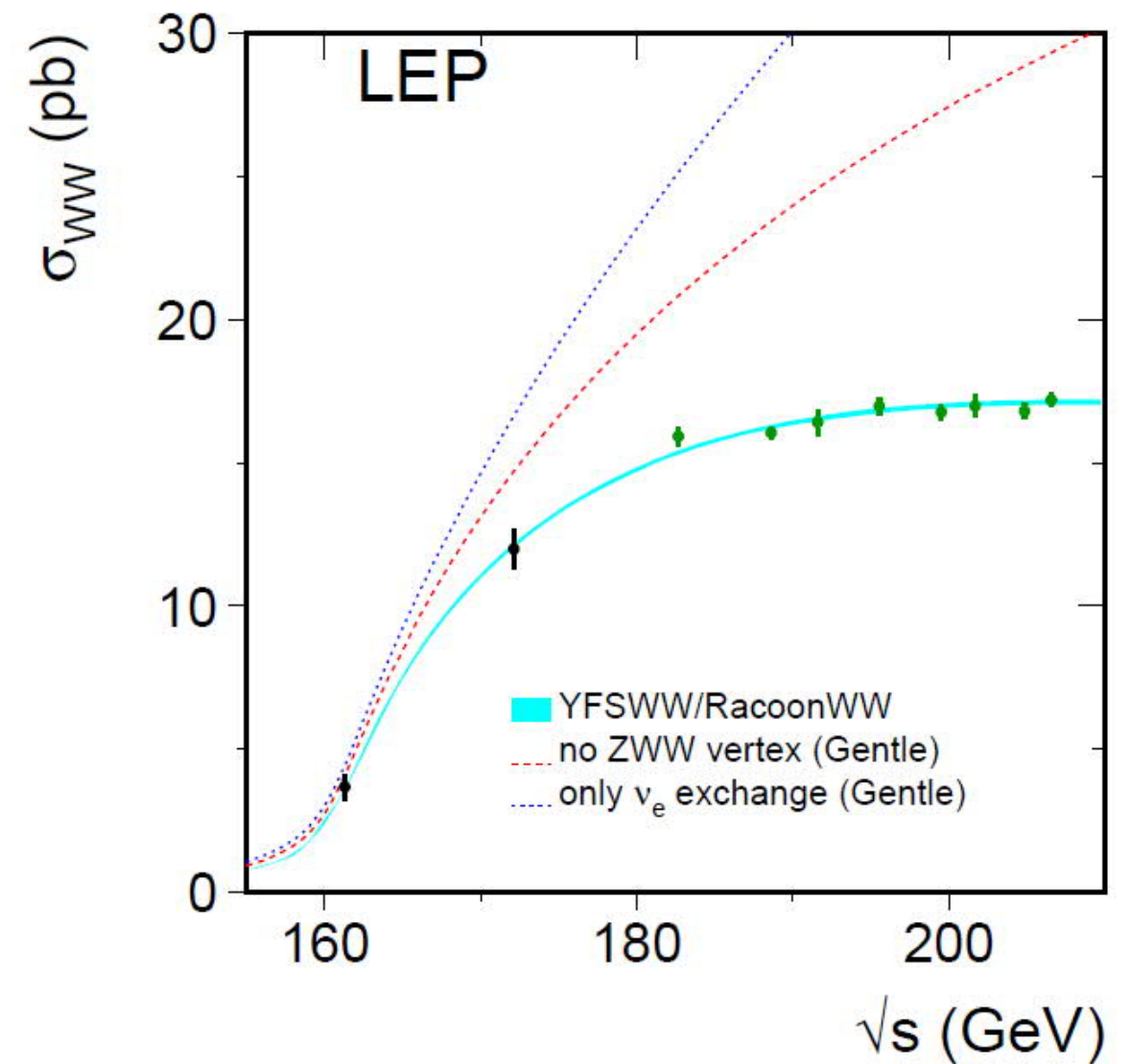
- Feynman-diagrams (Born)
 - TGC: γWW und ZWW



- Neutrino-exchange



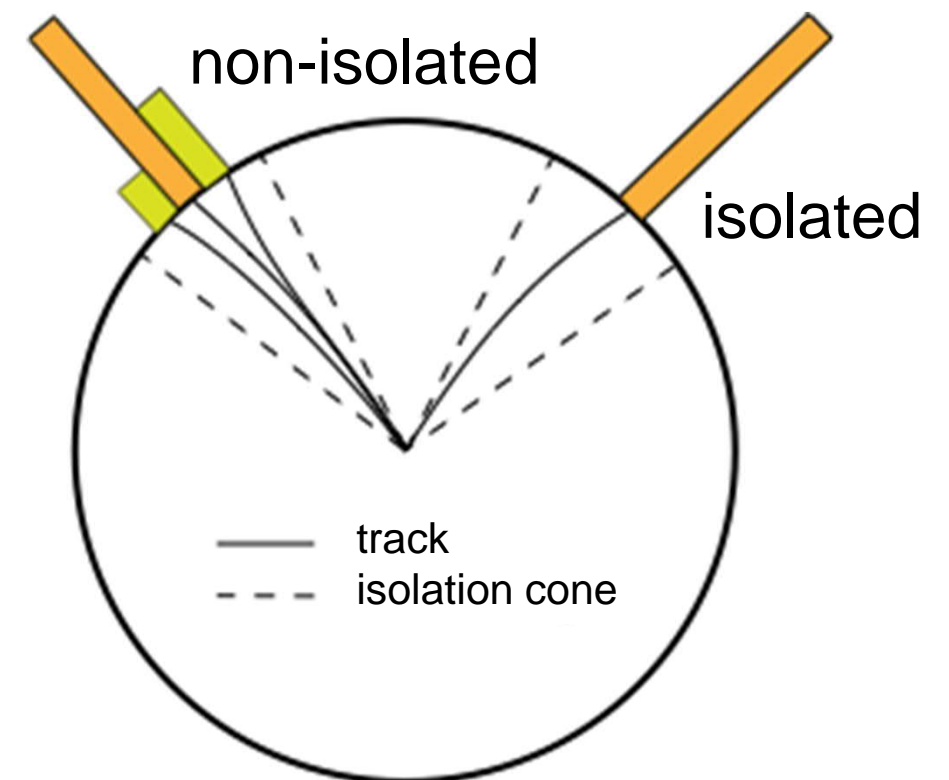
Cross section: $e^+e^- \rightarrow W^+W^-$



[arXiv:1302.3415]

W-Mass at Hadron-Colliders

- Start the reconstruction with **lepton**
 - **Isolated** leptons with high transverse momentum → suppress multijet background
 - **Prototype** for many high- p_T -analyses at hadron-colliders
 - Additional hadronic activity → recoil against W

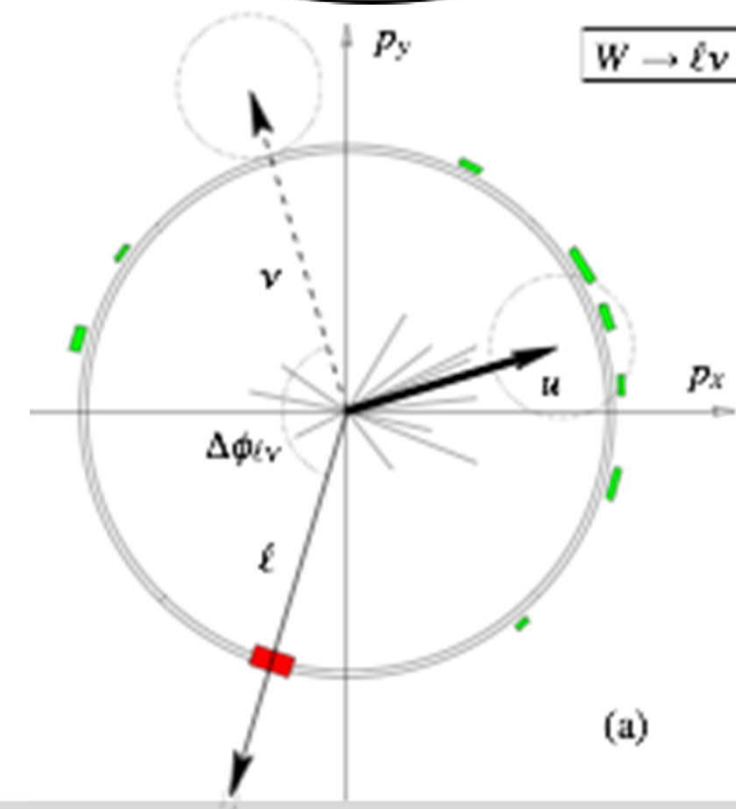


- Observable: **transverse mass**

$$m_T^2 = (E_T^\ell + E_T^\nu)^2 - (\vec{p}_T^\ell + \vec{p}_T^\nu)^2$$

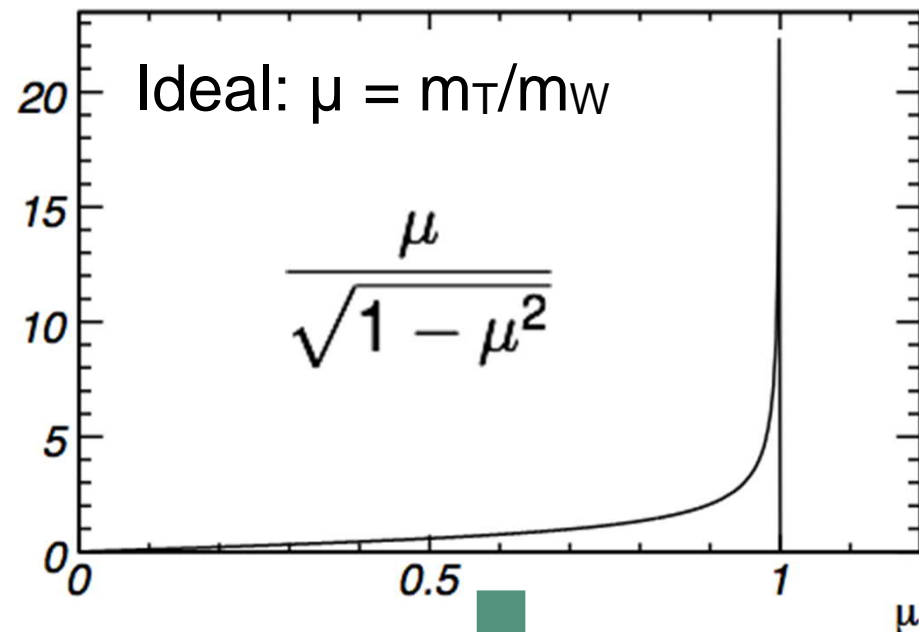
$$\approx 2 |\vec{p}_T^\ell| |\vec{p}_T^\nu| (1 - \cos \Delta\phi_{\ell\nu})$$

- p_T^ν missing transverse energy in the event → assumed to represent neutrino

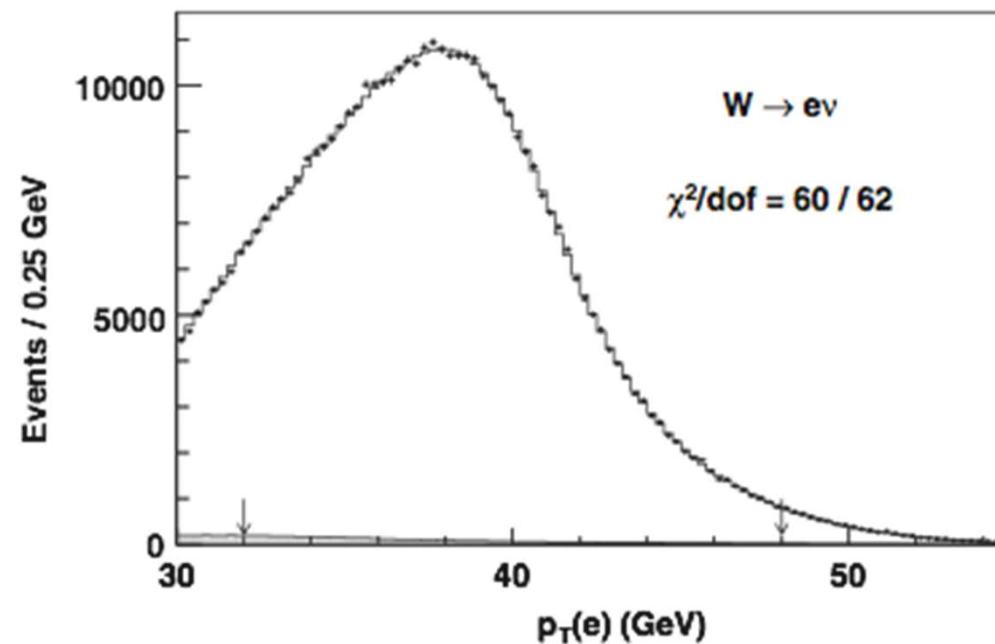
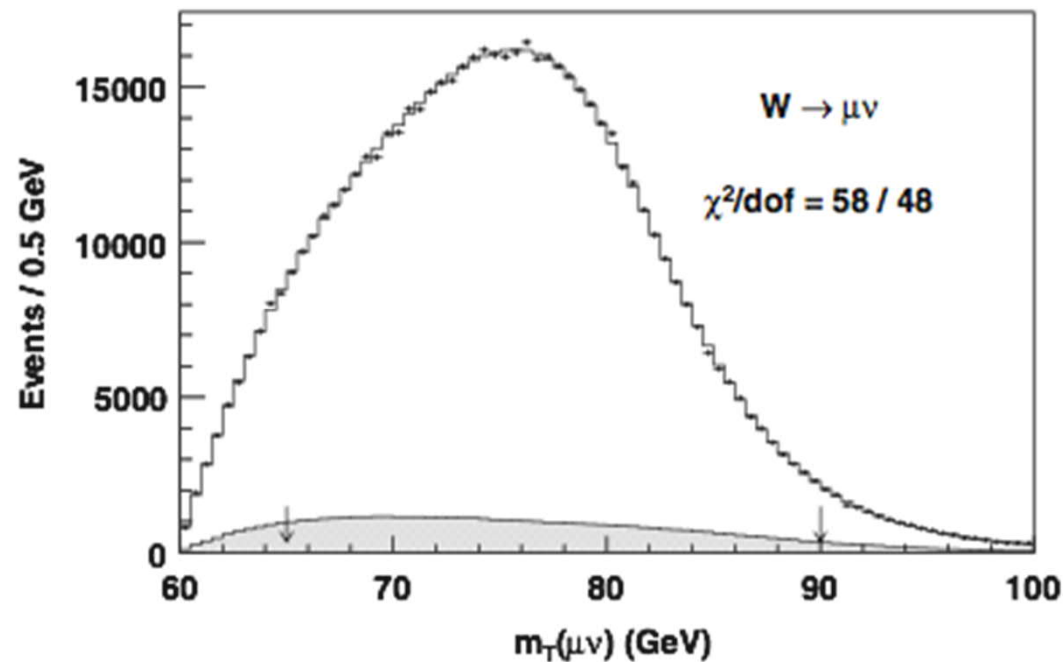


[CERN-OPEN-2008-020]

W-Mass at Hadron-Colliders

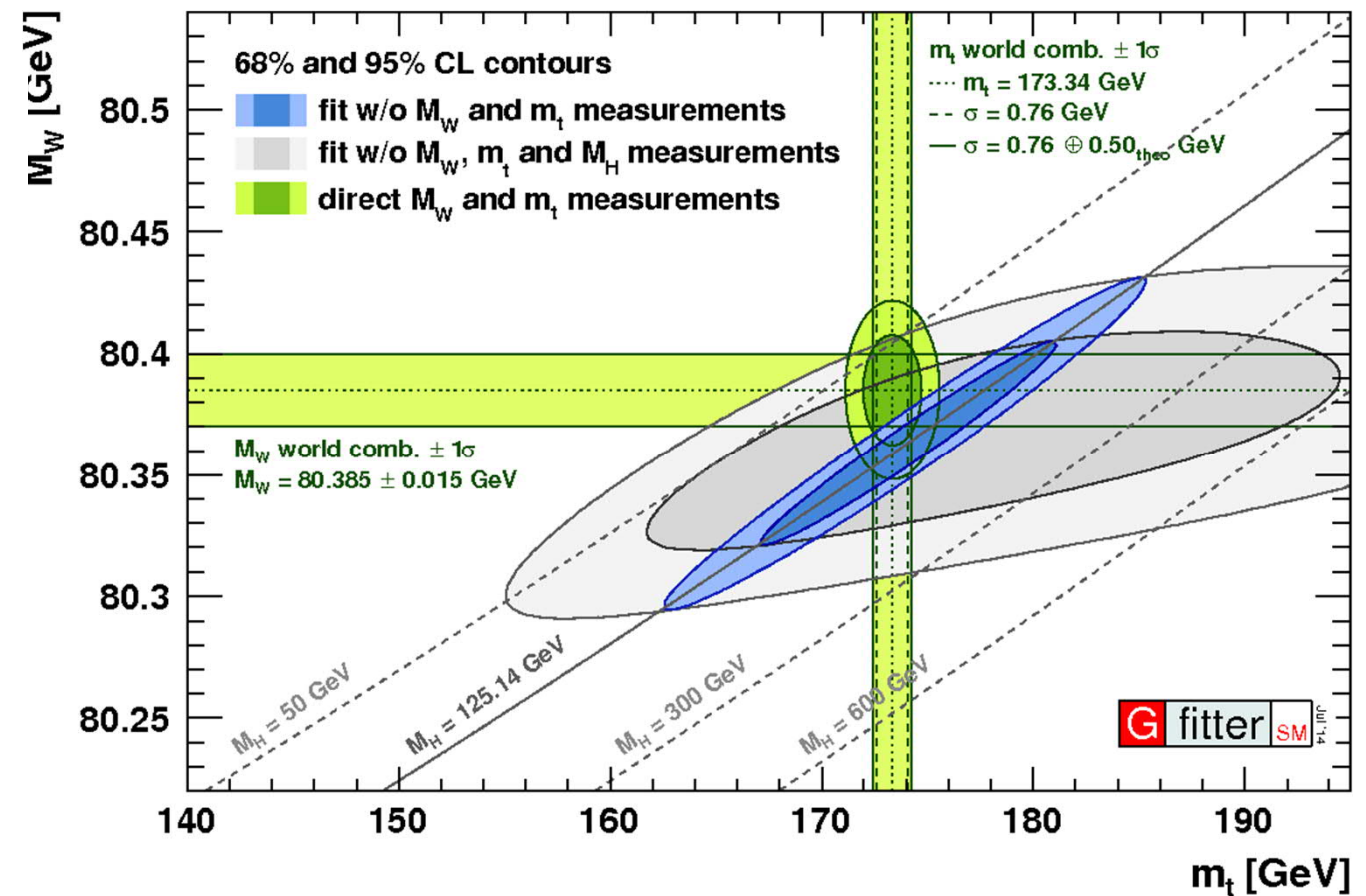
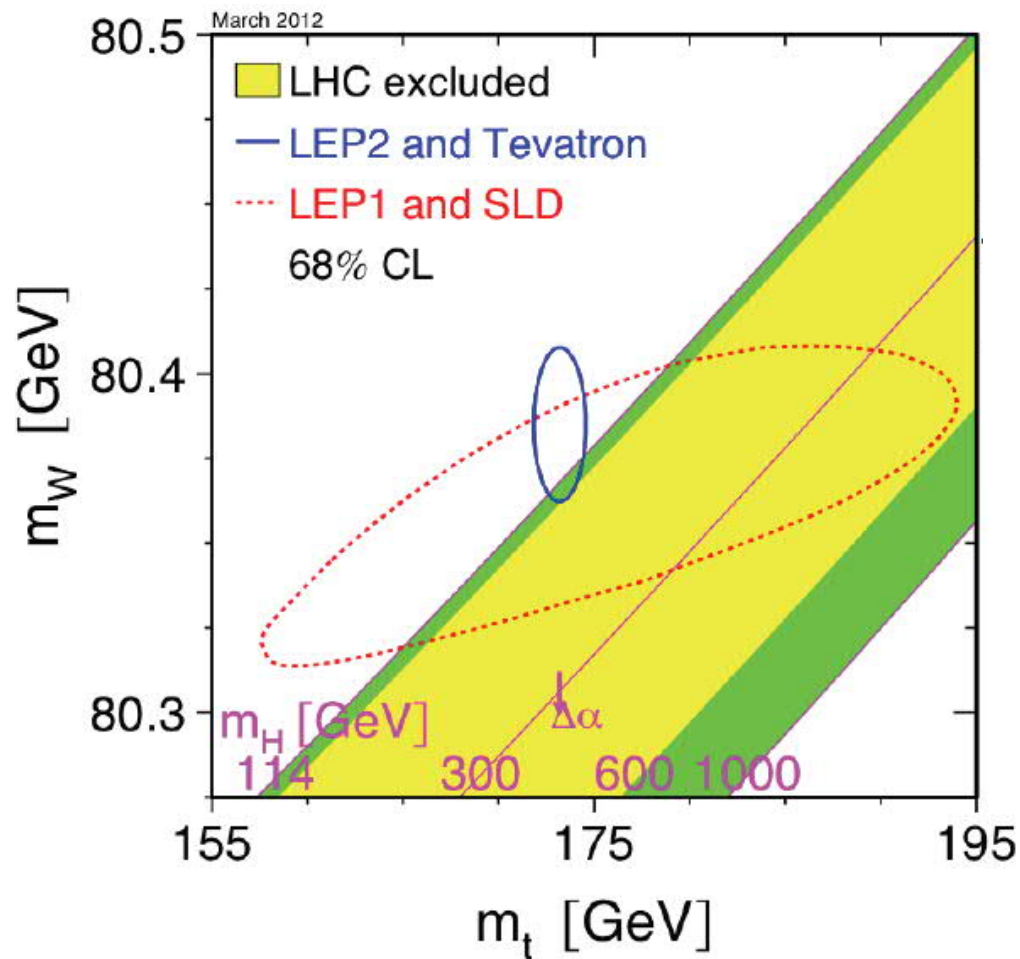


- Analysis: Likelihood-fit with **templates** from simulations with different W-masses
- Precision limited by **systematic uncertainties**: parton-densities, lepton energies, ...



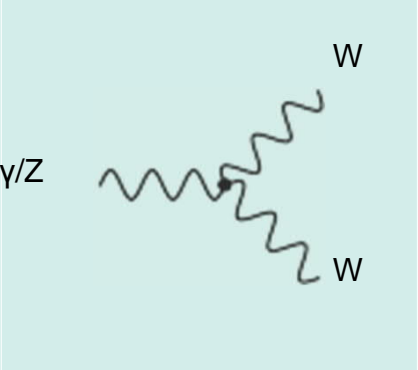
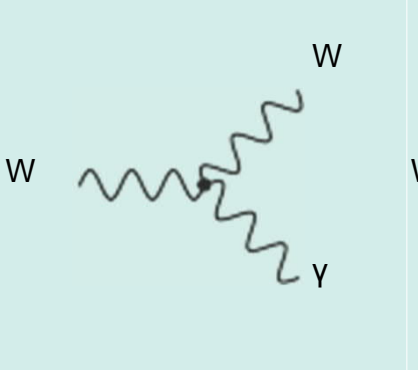
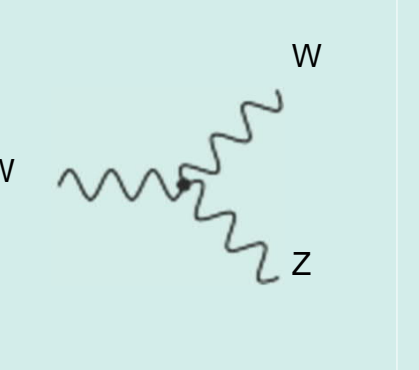
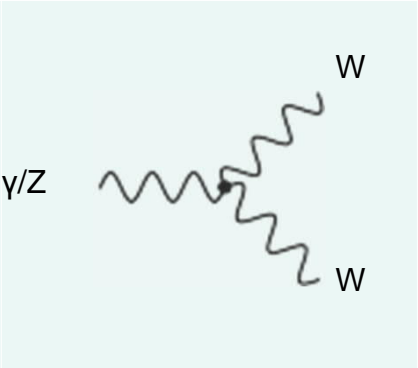
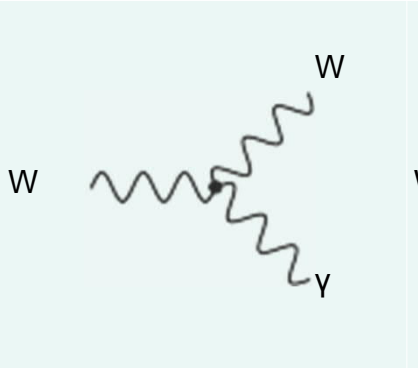
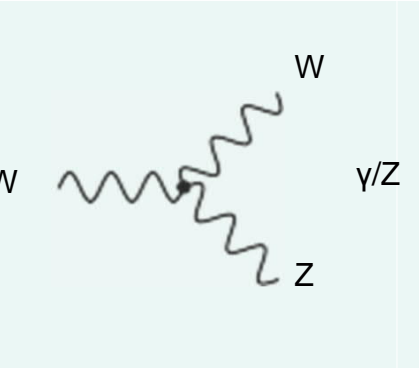
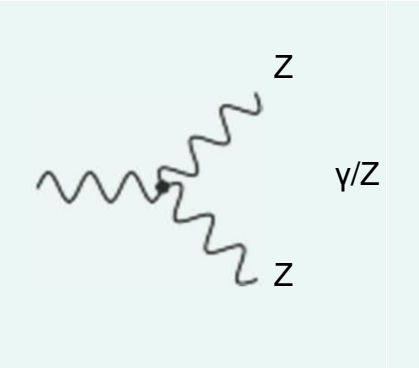
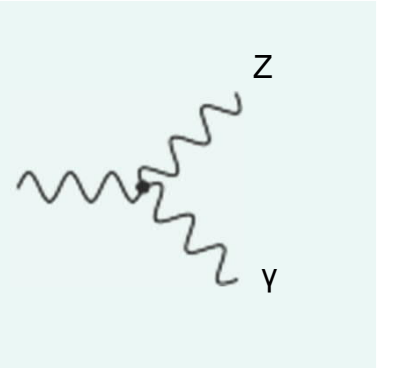
[Phys. Rev. Lett. 108 (2012) 151803]

Then and Now



- Pre-Higgs discovery: use $M_W + M_{\text{top}}$ to constrain Higgs mass
- Post Higgs discovery: use $M_W + M_{\text{top}} + M_H$ to constrain exotic theories that could add more particles to the loops
 → very strong limits on supersymmetry

Anomalous TGC

	WW	γ W	WZ	ZZ	γ Z
SM					
aTGC					

- Modified **couplings** caused by physics beyond the SM:
higher cross sections, especially at high **V-transverse-momentum**

Anomalous TGC: Lagrange density

- Remember: SM only allows $WW\gamma$ and WWZ triple boson vertices

$$\mathcal{L}_{WW\gamma} = -ie[A_\mu(W^{-\mu\nu}W_\nu^+ - W^{+\mu\nu}W_\nu^-) + F_{\mu\nu}W^{+\mu}W^{-\nu}]$$

$$\mathcal{L}_{WWZ} = -ie \cot\vartheta_w[Z_\mu(W^{-\mu\nu}W_\nu^+ - W^{+\mu\nu}W_\nu^-) + Z_{\mu\nu}W^{+\mu}W^{-\nu}]$$

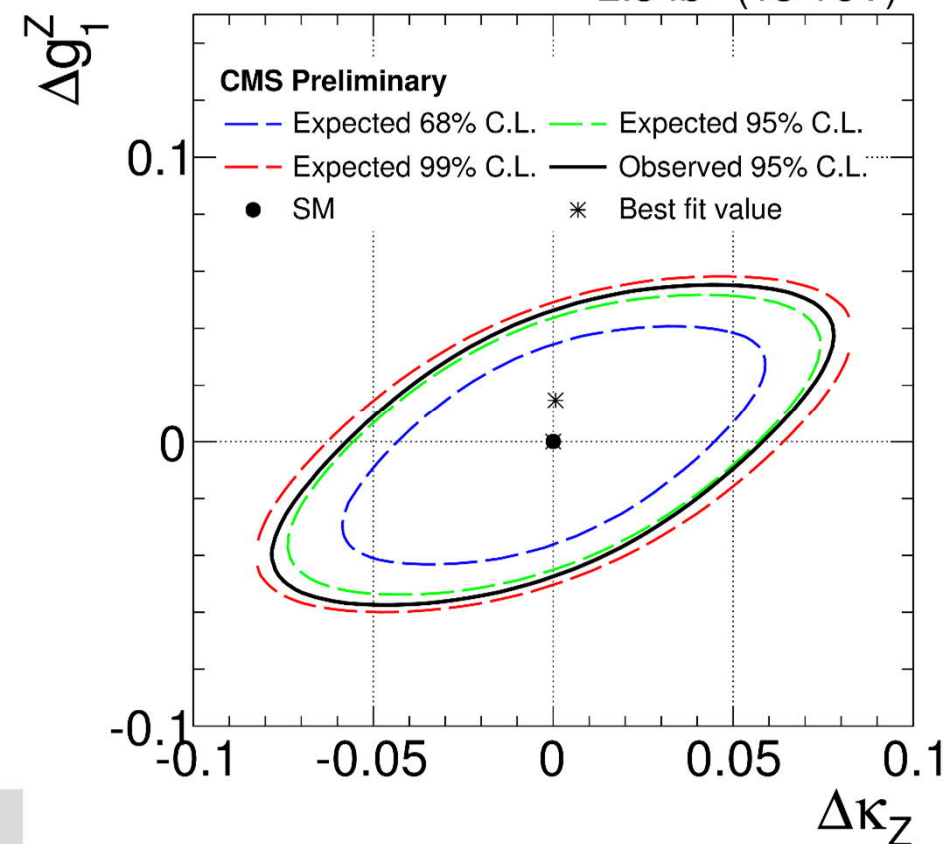
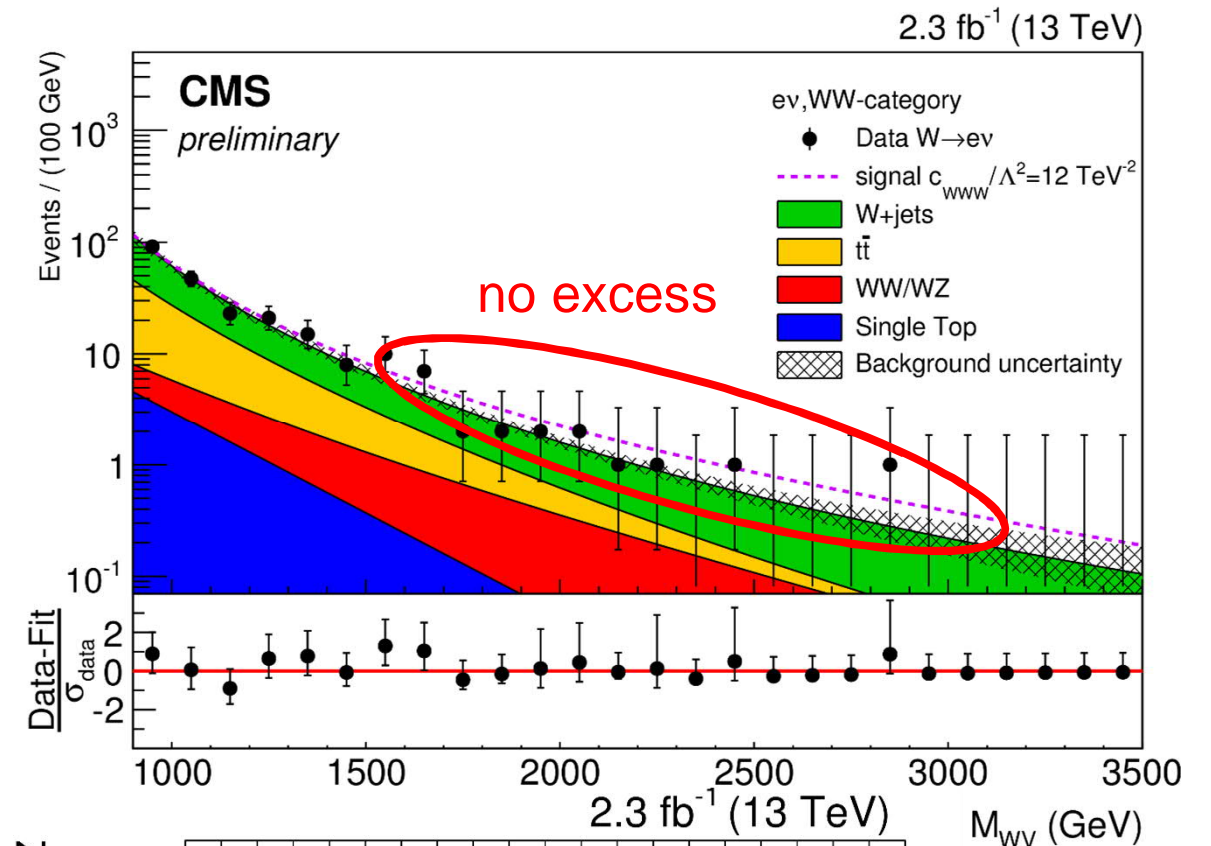
- Most general possible effective Lagrange density ($V = \gamma, Z$)

$$\begin{aligned} \mathcal{L}_{WWV}^{\text{eff}} = & -ig_{WWV} \left[g_1^V V_\mu (W^{-\mu\nu}W_\nu^+ - W_{\mu\nu}^+ W_\nu^-) + \kappa_V V_{\mu\nu} W^{+\mu}W^{-\nu} \right. \\ & + \frac{\lambda_V}{m_W^2} V_{\mu\nu} W^{+\nu\rho}W_\rho^{-\mu} + ig_5^V \epsilon_{\mu\nu\rho\sigma} ((\partial^\rho W^{-\mu})W^{+\nu} - W^{-\mu}(\partial^\rho W^{+\nu})) V^\sigma \\ & \left. + ig_4^V W_\mu^- W_\nu^+ (\partial^\mu V^\nu + \partial^\nu V^\mu) - \frac{\tilde{\kappa}_V}{2} W_\mu^- W_\nu^+ \epsilon^{\mu\nu\rho\sigma} V_{\rho\sigma} - \frac{\tilde{\lambda}_V}{2m_W^2} W_{\rho\mu}^- W_\nu^{+\mu} \epsilon^{\nu\rho\alpha\beta} V_{\alpha\beta} \right] \end{aligned}$$

- SM: $g_1^V = \kappa_V = 1$, all other couplings vanish
- C- und P-Erhaltung: $g_1^V, \kappa_V \neq 1, \lambda_V \neq 0, g_4^V = g_5^V = \tilde{\kappa}_V = \tilde{\lambda}_V = 0$

Limits on aTGCs

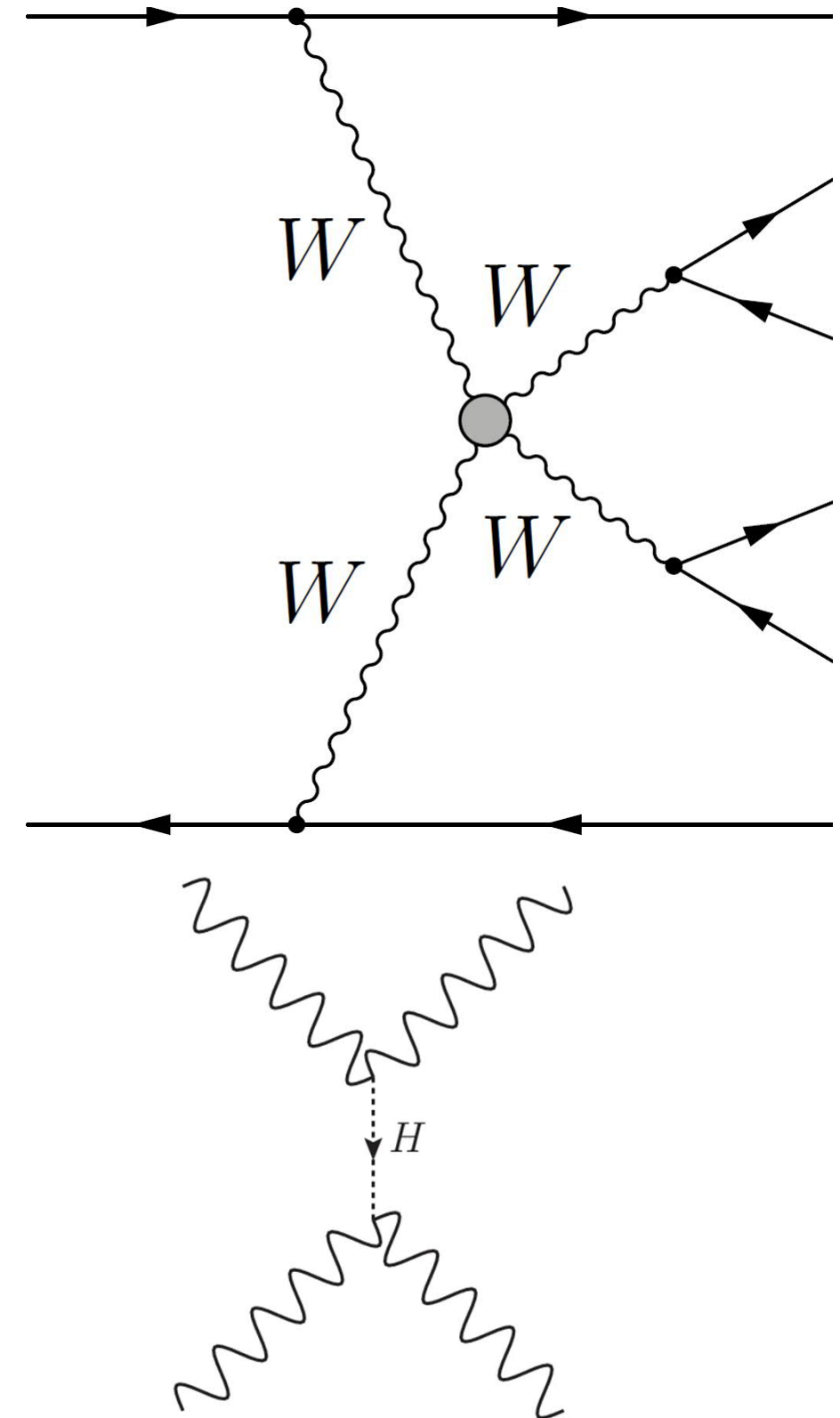
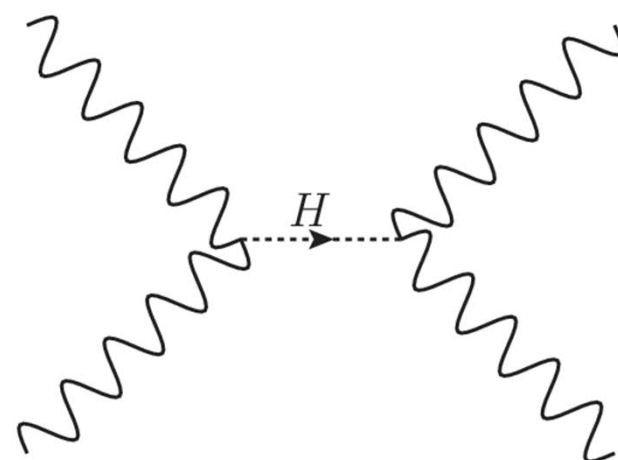
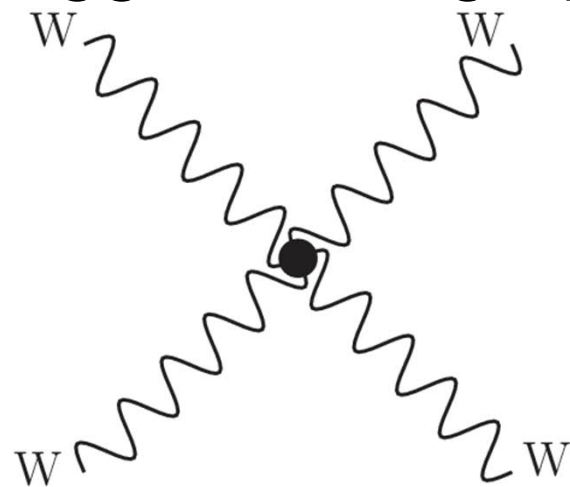
- Example:
CMS W + W/Z semi-leptonic
- Isolated lepton + MET
→ leptonic W candidate
- Fat jet with substructure
→ hadronic W candidate
- Search for excess at high diboson invariant masses
→ high aTGC contributions
- Extract limits from likelihood contours in signal+background fit



[CMS-PAS-SMP-16-012]

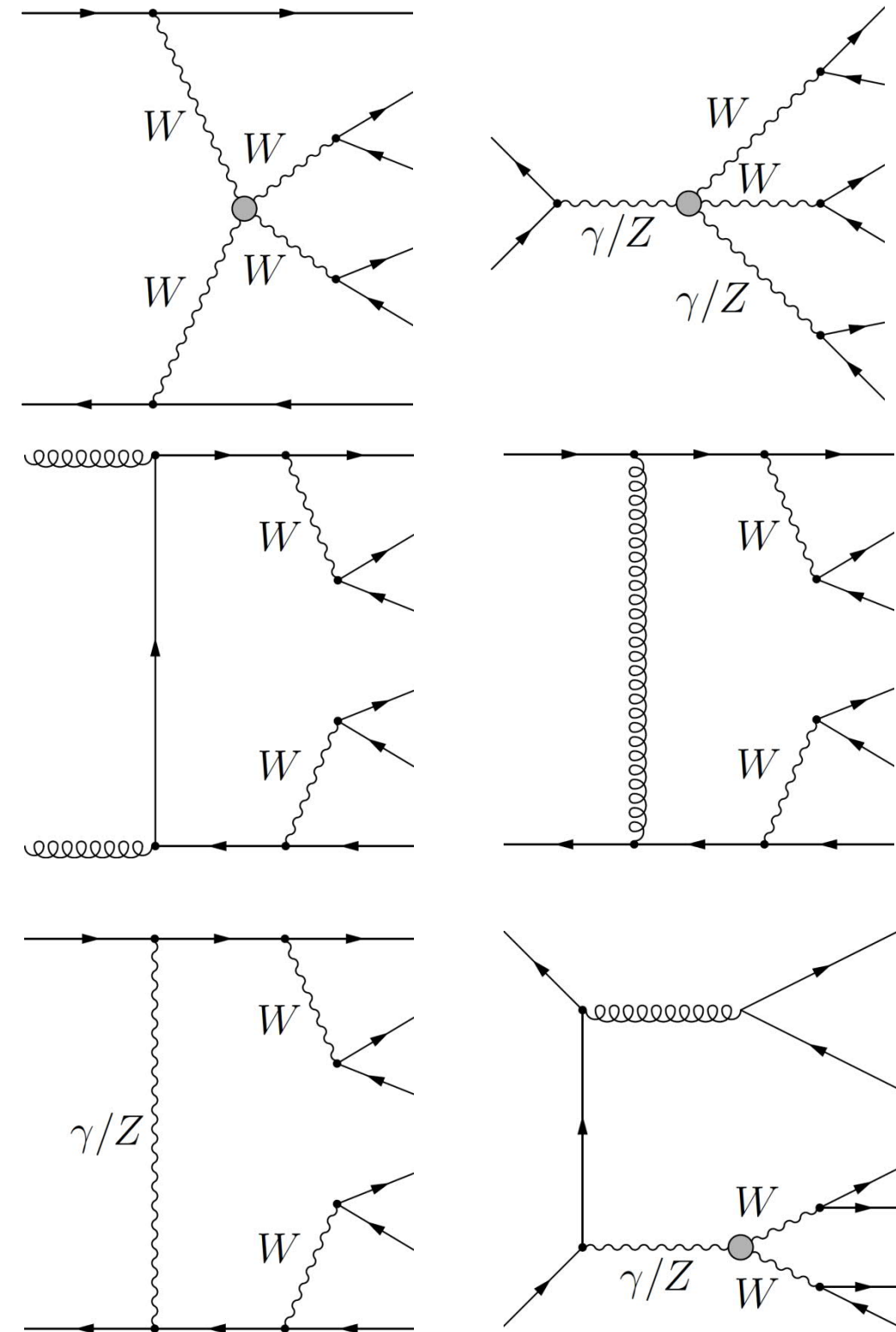
Vector-Boson Scattering

- Study quartic vertex in vector boson scattering
- Similar to VBF process in Higgs boson physics
- Scattering cross section with longitudinally polarized bosons not unitary at high energies
- Regularised in SM by interference with Higgs boson graphs



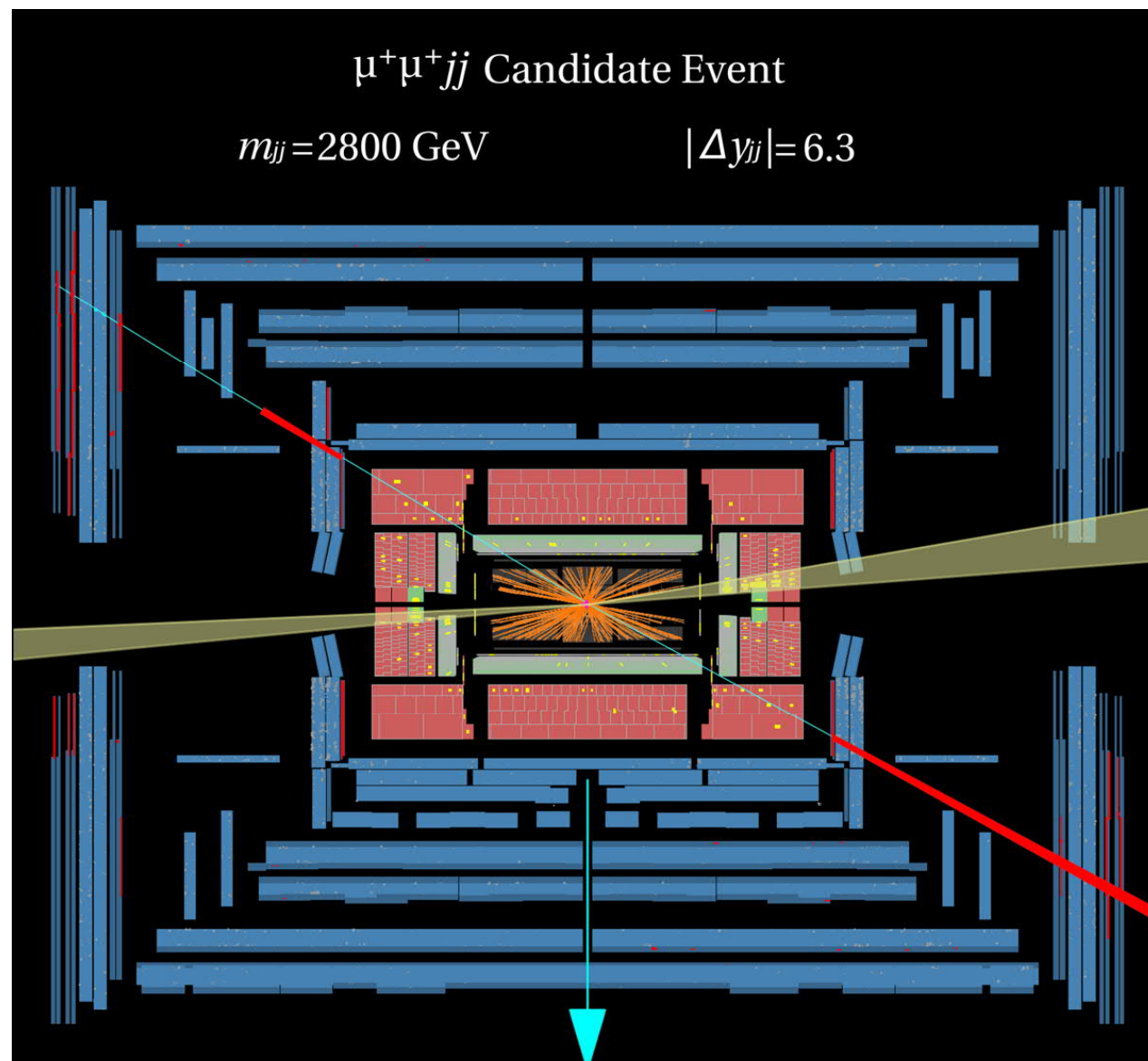
Vector-Boson Scattering

- Detailed test of the Higgs mechanism
- $2W + 2\text{jet}$ processes „common“ even without quartic vertex
- Enhance VBS with suitable selections:
 - jets with high dijet-mass
 - large difference in rapidity
- Only look at $W^{+/-}W^{+/-}$
=> no gluons in the initial state

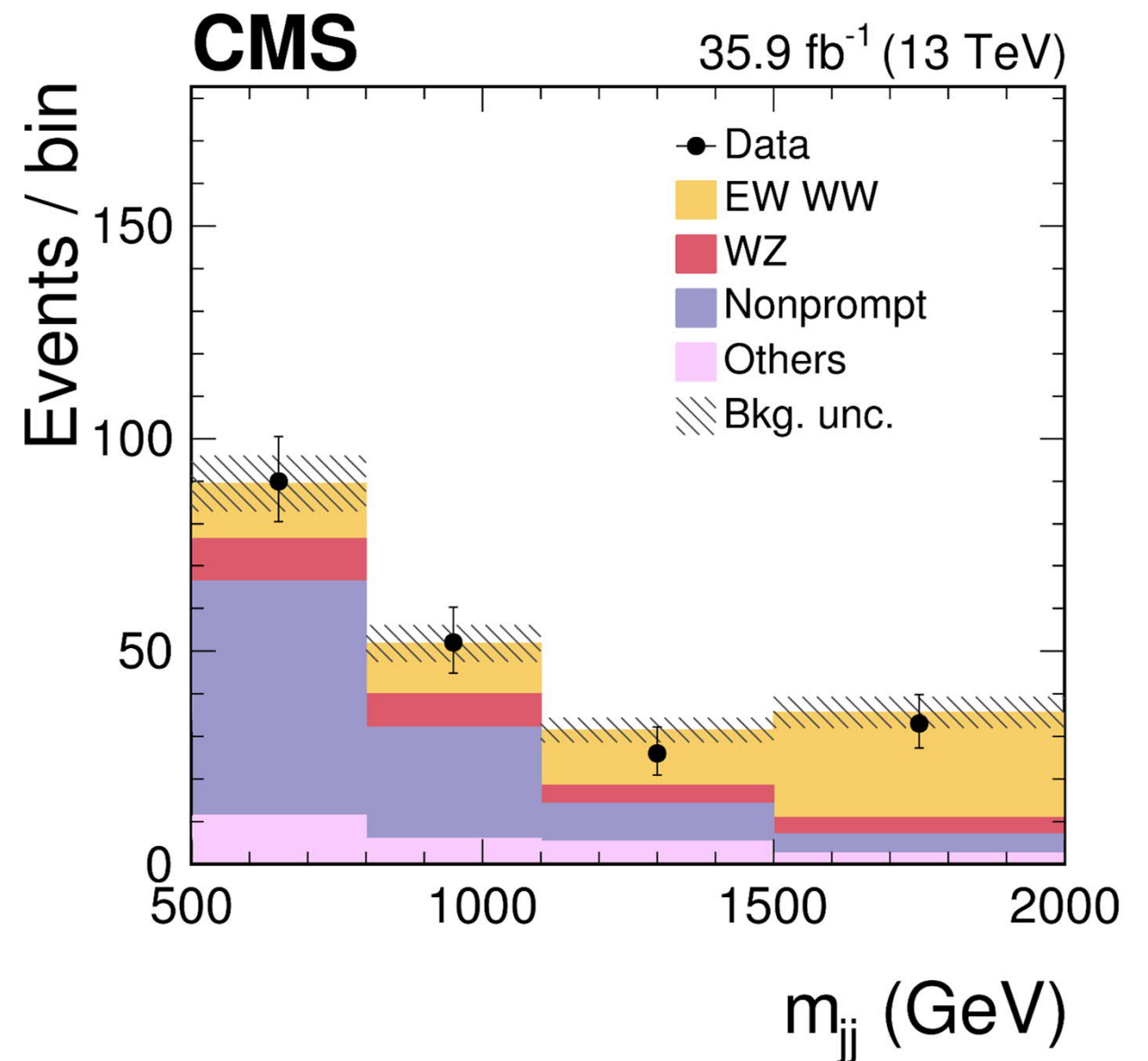


Quartic Couplings

- Measurement in $pp \rightarrow W^+W^+jj$
- Limits on anomalous couplings



[Atlas-CONF-2014-013]



[CMS-SMP-17-004]